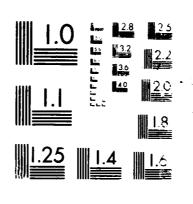
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A DETAILED NUMERICAL, GRAPHICAL, AND EXPERIMENTAL STUDY OF OBLIQUE SHOCK WAVE REFLECTIONS

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M3H 5T6

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Oblique Shock-Wave Reflections Numerical Simulation 10 11 Interferometry - Isopycnics Shock-Tube Flows 20 6					
An extensive series of numerical calculations of oblique-shock-wave reflections in air and argon have been performed using a version of the second-order Eulerian Godunov scheme for inviscid compressible flow. This scheme is among the best of the upwind schemes developed in recent years. The results have been compared with the best available interferometric data from the UTIAS 10 cm x 18 cm shock tube, for fifteen different cases. The objective of this portion of the study was to assess the accuracy of the computer code in computing two-dimensional shocked flow of an inviscid perfect gas. A significant portion of our analysis is devoted to the question of the extent of influence of viscous and vibrational nonequilibrium effects on the experimental flow fields.					
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19. ABSTRACT (Continued)

Further parametrized series of calculations were performed in an effort to study the feasibility of numerically constructing inviscid transition lines in the $(M_{\mbox{\scriptsize g}},\, \theta_{\mbox{\scriptsize W}})$ -plane. Good agreement with analytic predictions was found for low values of $M_{\mbox{\scriptsize g}}$ and, as might be expected, there are substantial discrepancies for $M_{\mbox{\scriptsize g}}=8.75$. The possibility of using such numerical results in the formulation of accurate transition criteria is discussed.

Overall, the computer code has been found to represent a significant predictive capability. The future extension of the code to permit the detailed modelling of nonequilibrium and viscous effects is, however, an important objective.

SUMMARY

An extensive series of numerical calculations of oblique-shock-wave reflections in air and argon have been performed using a version of the second-order Eulerian Godunov scheme for inviscid compressible flow. This scheme is among the best of the upwind schemes developed in recent years.

. The results have been compared with the best available interferometric data from the UTIAS 10 cm x 18 cm shock tube, for fifteen different cases. The objective of this portion of the study was to assess the accuracy of the computer code in computing two-dimensional shocked flow of an inviscid perfect gas. A significant portion of our analysis is devoted to the question of the extent of influence of viscous and vibrational nonequilibrium effects on the experimental flow fields.

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Overall, the computer code has been found to represent a significant predictive capability. The future extension of the code to permit the detailed modelling of nonequilibrium and viscous effects is, however, an important objective.

PREFACE

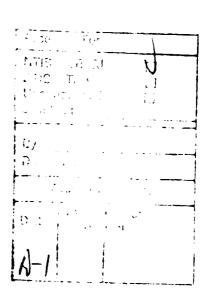
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TABLE OF CONTENTS

Section	n Pa	ge
	SUMMARY	ii
	PREFACE	iv
	LIST OF ILLUSTRATIONS	vi
1	INTRODUCTION	1
2	OBLIQUE SHOCK-WAVE REFLECTIONS	3
3	EXPERIMENTAL TECHNIQUES	6
	3.1 Experimental Facility	6
	3.2 Data Reduction Techniques	7
4	NUMERICAL METHOD	8
5	COMPUTATIONAL RESULTS	11
	5.1 Comparison of Experiment with Calculation	11
	5.2 Transition Sequences	23
6	CONCLUSIONS	28
7	LIST OF REFERENCES	30





LIST OF ILLUSTRATIONS

Figure		Page
	<pre>1 - Schematic diagrams of types of oblique shock-wave reflections: (a) RR; (b) SMR; (c) CMR; (d) DMR; also definitions of L and s.</pre>	34
	2 - Regions of RR, SMR, CMR, and DMR and their transition boundaries i the ($^{\rm M}_{\rm S}$, $^{\rm G}_{\rm W}$)-plane for perfect (frozen) air solid lines and imperfect (equilibrium) air broken lines, $^{\rm C}_{\rm O}$ = 2.00 kPa, $^{\rm C}_{\rm O}$ = 300 K, $^{\rm C}_{\rm Y}$ = 1.40.	n 35
	3 - Numerical scheme for flow initialization; (a) starting procedure;(b) shock reaching corner; (c) elimination of small disturbances.	36
	4 - Case 1, M_S = 2.05, θ_W = 600, Argon, γ = 5/3, RR.	37
	4a - Interferogram	37
	. 4b - Calculated isopycnics using the experimental fringes	37
	4c - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \bar{u}	38
	4d - Whole-flowfield contour-plots	39
	4e - Blowup-frame plots	42
	5 - Case 2, M_S = 1.26, θ_W = 450, Air, γ = 1.4 and Hansen EOS, RR.	44
	5a - Interferogram	44
	$5b_p$ - Calculated isopycnics (γ =1.4) using the experimental fringes	44
	$5b_{H}$ - Calculated isopycnics (Hansen) using the experimental fringes $5c$ - Wall plot for ρ/ρ_{0} , γ = 1.4 and Hansen calculations, with	45
	experimental data	45
	$5c_p$ - Wall plot for p/p_0 , ρ/p_0 with experimental data included, e, \tilde{u} ; γ = 1.4.	46
	$5c_{H}$ - Wall plot for p/p ₀ , ρ/ρ_{0} with experimental data included, e,	
	u; Hansen.	47
	$5dp$ - Whole-flowfield contour-plots; $\gamma = 1.4$.	48
	$5e_0$ - Blowup-frame plots: $y = 1.4$	51

Fig	ure	Page
	5d _H - Whole-flowfield contour-plots; Hansen	53
	5e _H - Blowup-frame plots; Hansen	56
6	- Case 3, M_S = 1.50, θ_W = 450, Air, γ = 1.4 and Hansen EOS, SMR.	58
	6a - Interferogram	58
	$6b_p$ - Calculated isopycnics (γ = 1.4) using the experimental fringes	58
	$6b_{\rm H}$ - Calculated isopycnics (Hansen) using the experimental fringes $6c$ - Wall plot for ρ/ρ_0 , γ = 1.4 and Hansen calculations, with	59
	experimental data	59
	$6c_p$ - Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e,	
	\tilde{u} ; $\gamma = 1.4$	60
	$6c_H$ - Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e,	٠.
	u; Hansen	61
	$6dp$ - Whole-flowfield contour-plots; $\gamma = 1.4$	62
	<pre>6ep - Blowup-frame plots; γ = 1.4</pre>	65
	6d _H - Whole-flowfield contour-plots; Hansen	67
	6e _H - Blowup-frame plots; Hansen	70
7	- Case 4, M_S = 3.03, θ_W = 470, Air, γ = 1.4 and Hansen EOS, DMR.	72
	7a - Interferogram	72
	$7b_p$ - Calculated isopycnics ($\gamma = 1.4$) using the experimental fringes	72
	$7b_{\mbox{\scriptsize H}}$ - Calculated isopycnics (Hansen) using the experimental fringes	72
	7c - Wall plot for p/p_0 , $\gamma = 1.4$ and Hansen calculations, with	70
	experimental data	73
	7 cp - Wall plot for p/p ₀ , ρ/ρ_0 with experimental data included, e, \tilde{u} ; $\gamma = 1.4$	74
	$7c_{\rm H}$ - Wall plots for p/p _o , $\rho/\rho_{\rm O}$ with experimental data included, e,	
	u; Hansen.	75
	$7d_p$ - Whole-flowfield contour-plots; $\gamma = 1.4$	76
	$7e_p$ - Blowup-frame plots; $\gamma = 1.4$	79
	7d _H - Whole-flowfield contour-plots; Hansen	81
	7. Diama farma platas Hancon	0.4

igure		Page
8 - Case 5, M _s =	2.65, $\theta_{\rm W}$ = 300, Air, γ = 1.4 and Hansen EOS, CMR.	86
3a - Interferog	ram	36
8bp - Calculate	d isopycnics (γ = 1.4) using the experimental fringe	s 36
3h _H - Calculate	d isopycnics (Hansen) using the experimental fringes	87
3c - Wall plot	for p/p_0 , $\gamma = 1.4$ and Hansen calculations, with	
experiment	al data	87
8cp - Wall plot u; γ =	for p/p_0 , ρ/ρ_0 with experimental data included, e, 1.4.	88
8c _H - Wall plot	for p/p_0 , ρ/ρ_0 with experimental data included, e,	
u; Hans	·	89
8dp - Whole-flo	wfield contour-plots; γ = 1.4	90
8ep - Blowup-fr	ame plots; $\gamma = 1.4$	93
8d _H - Whole-flow	wfield contour-plots; Hansen	95
8e _H - Blowup-fra	ame plots; Hansen	98
9 - Case 6, M _s =	5.07, $\theta_{W} = 30^{\circ}$, Argon, $\gamma = 5/3$, CMR.	100
9a - Interferog	ram	100
9b - Calculated	isopycnics using the experimental fringes	100
9c - Wall plots	for p/p_0 , ρ/ρ_0 with experimental data included, e, \hat{u}	101
	field contour-plots	102
9e - Blowup-fram	ne plots	105
10 - Case 7, M _s =	10.37, $\theta_{\rm W}$ = 100, Air, Hansen EOS, CMR.	107
10a - Interfero	gram	107
10b - Calculated	d isopycnics using the experimental fringes	107
10c - Wall plots	s for p/p_0 , ρ/ρ_0 with experimental data included, e,	ũ 108
	wfield contour-plots	109
10e - Blowup-fra	ame plots	112

Figure

11 - Case 3, M_S = 1.56, θ_W = 400, Air, γ = 1.4 and Hansen EOS, SMR.	114
lla - Interferogram	114
$11b_p$ - Calculated isopycnics ($\gamma = 1.4$) using the experimental fringes	: 114
11b4 - Calculated isopycnics (Hansen) using the experimental fringes	115
11c - Wall plot for p/p_0 , $\gamma = 1.4$ and Hansen calculations, with	
experimental data	115
11cp - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \tilde{u} ; $\gamma = 1.4$	116
$11c_{H}$ - Wall plots for p/p_{O} , p/p_{O} with experimental data included, e,	110
u; Hansen.	117
11dp - Whole-flowfield contour-plots; γ = 1.4	118
11ep - Blowup-frame plots; $\gamma = 1.4$	121
11d _H - Whole-flowfield contour-plots; Hansen	123
lle _H - Blowup-frame plots; Hansen	126
12 - Case 9, M_S = 2.87, θ_W = 400, Air, γ = 1.4 and Hansen EOS, DMR.	128
12a - Interferogram	128
$12b_p$ - Calculated isopycnics (γ = 1.4) using the experimental fringes	; 128
12b _H - Calculated isopycnics (Hansen) using the experimental fringes	129
12c - Wall plots for p/p_0 , ρ/ρ_0 , γ = 1.4 and Hansen calculations, wit	
experimental data	129
12cp - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \tilde{u} ; $\gamma = 1.4$.	130
$12c_{H}$ - Wall plots for p/p_{o} , ρ/p_{o} with experimental data included, e,	
u; Hansen.	131
12dp - Whole-flowfield contour-plots; y = 1.4	132
12ep - Blowup-frame plots; y = 1.4	135
12d _H - Whole-flowfield contour-plots; Hansen	137
12e _H - Blowup-frame plots; Hansen	140

CIST OF ILLUSTRATIONS (Continued)	
	Page
13 - Case 10, M_S = 3.72, g_W = 400, Air, γ = 1.4 and Hansen EOS, DMR.	142
13a - Interferogram	142
$13b_p$ - Calculated isopycnics (γ = 1.4) using the experimental fringe	s 142
$13b_{\rm H}$ - Calculated isopycnics (Hansen) using the experimental fringes $13c$ - Wall plot for ρ/ρ_0 , γ = 1.4 and Hansen calculations, with	; 143
experimental data	143
$13c_p$ - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \tilde{u} ; $\gamma = 1.4$	144
$13c_{\rm H}$ - Wall plots for p/p ₀ , ρ/ρ_0 with experimental data included, e,	
u; Hansen	145
$13d_p$ - Whole-flowfield contour-plots; $\gamma = 1.4$	146
$13e_p$ - Blowup-frame plots; $\gamma = 1.4$	149
13d _H - Whole-flowfield contour-plots; Hansen	151
13e _H - Blowup-frame plots; Hansen	154
14 - Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR.	156
14a - Interferogram	156
$14b_p$ - Calculated isopycnics (γ = 1.4) using the experimental fringe	s 156
$14b_{\rm H}$ - Calculated isopycnics (Hansen) using the experimental fringes 14c - Wall plot for ρ/ρ_0 , γ = 1.4 and Hansen calculations, with	
experiment data	157
14c _p - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \tilde{u} ; $\gamma = 1.4$	158
$14c_{H}$ - Wall plots for p/p_{0} , ρ/ρ_{0} with experimental data included, e, \tilde{u} ; Hansen	159
$14dp$ - Whole-flowfield contour-plots; $\gamma = 1.4$	160
$14e_p$ - Blowup-frame plots; $\gamma = 1.4$	163
14d _H - Whole-flowfield contour-plots; Hansen	165
14e _H - Blowup-frame plots; Hansen	168

Figure		^o a ge
15 - Case 12, $M_S = 2.03$, $\theta_W = 27^{\circ}$, Air, $\gamma = 1.4$, SMR.		170
15a - Interferogram		170
15b - Calculated isopycnics using the experimental fringes		170
15c - Wall olots for p/p_0 , p/p_0 with experimental data included,	ē,	7171
15d - Whole-flowfield contour-plots		. 172
15e - Blowup-frame piots		175
16 - Case 13, M_S = 8.70, θ_W = 270, Air, Hansen EOS, DMR.		177
16a - Interferogram		177
16b - Calculated isopycnics using the experimental fringes		177
16c - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included,	e,	ũ 178
16d - Whole-flowfield contour-plots		179
16e - Blowup-frame plots		181
17 - Case 14, M_S = 7.19, θ_W = 200, Air, Hansen EOS, C/DMR.		183
17a - Interferogram		183
17b - Calculated isopycnics using the experimental fringes		183
17c - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included,	e,	ũ 184
17d - Whole-flowfield contour-plots		185
17e - Blowup-frame plots		188
18 - Case 15, $M_S = 8.86$, $\theta_W = 20^{\circ}$, Air, Hansen EOS, DMR.		190
18a - Interferogram		190
18b - Calculated isopycnics using the experimental fringes		190
18c - Wall plots for p/p_0 , ρ/ρ_0 with experimental data included,	e,	ũ 191
18d - Whole-flowfield contour-plots		192
18e - Blowup-frame plots		195
18f - Reproduction of the interferogram of Exp. 974, Ref. [14];		
M = 10.18 a - 200 Air		196

Figure	2-40
rigure	Page
19 - Transition set 1, $\theta_w = 45^{\circ}$, $\gamma = 1.4$	197
19.1a - M _s = 1.30, whole-flowfield contour-plots	197
19.1b - M _s = 1.30, blowup-frame plots	198
19.2a - $M_s = 1.40$, whole-flowfield contour-plots	200
19.2b - M _s = 1.40, blowup-frame plots	201
19.3a - M _s = 1.50, whole-flowfield contour-plots	203
19.3b - M _s = 1.50, blowup-frame plots	204
19.4a - M _s = 1.60, whole-flowfield contour-plots	206
19.4b - M _s = 1.60, blowup-frame plots	207
19.5a - M _s = 1.70, whole-flowfield contour-plots	209
19.5b - M _s = 1.70, blowup-frame plots	210
19.6a - M _s = 1.80, whole-flowfield contour-plots	212
19.6b - M _s = 1.80, blowup-frame plots	213
19.7a - M _s = 1.90, whole-flowfield contour-plots	215
$19.7b - M_s = 1.90$, blowup-frame plots	216
19.8a - M _s = 2.00, whole-flowfield contour-plots	218
19.8b - M_s = 2.00, blowup-frame plots	219
19.9a - M _s = 2.10, whole-flowfield contour-plots	221
19.9b - M _s = 2.10, blowup-frame plots	222
$19.10a - M_s = 2.20$, whole-flowfield contour-plots	224
19.10b - M _s ≈ 2.20, blowup-frame plots	225
19.11a - $M_s = 2.30$, whole-flowfield contour-plots	227
19.11b - $M_s = 2.30$, blowup-frame plots	228
19.12a - $M_s = 2.40$, whole-flowfield contour-plots	230
19.12b - M _s = 2.40, blowup-frame plots	231
$19.13a - M_S = 2.50$, whole-flowfield contour-plots	233
19.13b - M _s ≈ 2.50, blowup-frame plots	234
19.14a - M _s = 2.60, whole-flowfield contour-plots	236
19.14b - M _s = 2.60, blowup-frame plots	237

Figure	Page
20 - Transition set 1, $\theta_{\rm W}$ = 450, Hansen	239
20.1a - M _s = 1.50, whole-flowfield contour-plots	239
20.1b - $M_s = 1.50$, blowup-frame plots	240
20.2a - $M_s = 1.60$, whole-flowfield contour-plots	242
$20.2b - M_s = 1.60$, plowup-frame plots	243
20.3a - $M_s = 1.70$, whole-flowfield contour-plots	245
20.3b - M _s = 1.70, blowup-frame plots	246
20.4a - M _s = 1.80, whole-flowfield contour-plots	248
$20.4b - M_s = 1.80$, blowup-frame plots	249
20.5a - M_s = 1.90, whole-flowfield contour-plots	251
20.5b - $M_s = 1.90$, blowup-frame plots	252
20.6a - M _s = 2.00, whole-flowfield contour-plots	254
20.6h - $M_s = 2.00$, blowup-frame plots	255
20.7a - M_s = 2.10, whole-flowfield contour-plots	257
$20.7b - M_s = 2.10$, blowup-frame plots	258
$20.8a - M_s = 2.20$, whole-flowfield contour-plots	260
20.8b - M_s = 2.20, blowup-frame plots	261
20.9a - M_s = 2.30, whole-flowfield contour-plots	263
20.9b - M_S = 2.30, blowup-frame plots	264
21 - Transition set 2, $M_S = 4.0$, $\gamma = 1.4$	266
21.1a - $\theta_w = 29^{\circ}$, whole-flowfield contour-plots	266
21.1b $-\theta_{\parallel} = 29^{\circ}$, blowup-frame plots	267
21.2a - θ_u^2 = 30°, whole-flowfield contour-plots	269
$21.2b - \theta_{u}^{u} = 30^{\circ}$, hlowup-frame plots	270
21.3a - $\theta_{\rm w}$ = 310, whole-flowfield contour-plots	272
21.3b - θ_{u}^{u} = 310, blowup-frame plots	273
$21.4a - \theta_{u} = 320$, whole-flowfield contour-plots	275
21.4b - θ_{u}^{u} = 320, blowup-frame plots	276
21.5a - θ_{w}^{w} = 330, whole-flowfield contour-plots	278
21.5b - e = 330, blowup-frame plots	279
21.6a - θ_{w}^{w} = 340, whole-flowfield contour-plots	281
$^{\text{w}}$ 21.6b - θ = 340, blowup-frame plots	282

Figure	Pag e
22 - Transition set 2, $M_s = 4.0$, Hansen	284
22.la - ə = 250, whole-flowfield contour-plots	284
22.1b - $\theta_{\rm M}^2$ = 250, blowup-frame plots	285
22.2a - a = 260, whole-flowfield contour-plots	287
22.2b - 3 = 260, blowup-frame plots	288
22.3a - a = 270, whole-flowfield contour-plots	290
22.3b - ə = 270, blowup-frame plots	291
22.4a - ə = 280, whole-flowfield contour-plots	293
22.4b - $\frac{\pi}{2}$ = 280, blowup-frame plots	294
22.5a - ə = 290, whole-flowfield contour-plots	296
22.5b - a = 290, blowup-frame plots	297
22.5a - ə = 300, whole-flowfield contour-plots	299
22.6b - $\theta_{\rm w}^{\rm m}$ = 30°, blowup-frame plots	300
23 - Transition set 3, $M_S = 8.75$, $\gamma = 1.1$	302
23.1a - θ_{M} = 60, whole-flowfield contour-plots	302
23.1b - $\theta_{\rm m}^{\rm m}$ = 60, blowup-frame plots	303
23.2a - θ_{u}^{u} = 70, whole-flowfield contour-plots	305
23.2b - $\theta_{\rm w}$ = 70, blowup-frame plots	306
23.3a - $\theta_{u} = 80$, whole-flowfield contour-plots	308
23.3b - $\theta_0 = 80$, blowup-frame plots	309
23.4a - θ_u^2 = 90, whole-flowfield contour-plots	311
23.4b - 0 = 90, blowup-frame plots	312
23.5a - $\theta_0^{\prime\prime}$ = 100, whole-flowfield contour-plots	314
23.5b - θ_0 = 100, blowup-frame plots	315
23.6a - e = 220, whole-flowfield contour-plots	317
23.6b - e = 220, blowup-frame plots	318
23.7a - e = 230, whole-flowfield contour-plots	320
23.7b - θ = 230, hlowup-frame plots	321

Figure	Page
23.8a - e = 240, whole-flowfield contour-plots	323
23.8b - θ = 240, blowup-frame plots	324
23.9a - θ_{w}^{w} = 250, whole-flowfield contour-plots	326
23.9b - 9 = 250, blowup-frame plots	327
23.10a - ə = 260, whole-flowfield contour-plots	329
23.10b - $\theta_{\rm w}^{\rm v}$ = 260, blowup-frame plots	330
24 - Transition set 3, M _S = 8.75, Hansen	332
24.1a - θ_{ω} = 50, whole-flowfield contour-plots	332
24.1b - $\theta_{\rm w}^{\rm o}$ = 50, blowup-frame plots	333
24.2a - θ_{w}^{0} = 60, whole-flowfield contour-plots	335
24.2b - $\theta_{\rm w}^{\circ}$ = 60, blowup-frame plots	336
24.3a - $\theta_{\rm w}^{\circ}$ = 70, whole-flowfield contour-plots	338
24.3b - $\theta_{\rm w}$ = 70, blowup-frame plots	339
24.4a - $\theta_{\rm w}^{\rm w}$ = 80, whole-flowfield contour-plots	341
24.4b - e = 80, blowup-frame plots	342
24.5a - $\theta_{\rm w}$ = 90, whole-flowfield contour-plots	344
$24.5b - \theta_{\rm w} = 90$, blowup-frame plots	345
24.6a - $\theta_{\rm w}$ = 150, whole-flowfield contour-plots	347
24.6b - $\theta_{\rm w}^{\rm v}$ = 150, blowup-frame plots	348
24.7a - $\theta_{\rm w}$ = 160, whole-flowfield contour-plots	350
24.7b - $\theta_{\rm w}$ = 160, blowup-frame plots	351
24.8a - $\theta_{\rm w}$ = 170, whole-flowfield contour-plots	353
24.8b - $\theta_{\rm w}$ = 170, blowup-frame plots	354
24.9a - $\theta_{\rm w}$ = 180, whole-flowfield contour-plots	356
$24.9b - \theta_{\rm w} = 18^{\rm 0}$, blowup-frame plots	357
24.10a - $\theta_{\rm w}$ = 190, whole-flowfield contour-plots	359
$24.10b - \theta = 190$, blowup-frame plots	360

Figure	Page
25 - Transition set 4, $M_s = 7.10$, $\gamma = 5/3$, density contour-plots	362
25a - Interferogram, a = 490	362
25b - a = 190	362
25c - ə = 500	363
25d - 9 = 610	363
25e - ຈູ້ = 520	363
$25f - \theta_{W} = 52.75^{\circ}$	364
$25g - \theta_{W}^{W} = 53.0^{\circ}$	364
$25h - \theta_{w}^{w} = 53.10^{\circ}$	364
$25i - e_{w}^{w} = 53.20^{\circ}$	365
$25j - e_w^w = 53.30^\circ$	365
$25k - \theta_{W}^{W} = 53.40^{\circ}$	365
$251 - e_{w}^{w} = 53.50^{\circ}$	366
$25m - 9_w^w = 53.750$	366
25n - e = 540	366
25o - e = 55º	366

26 - Plot of DMR Mach stem height versus θ_w , extrapolated to zero height for RR(h/L = 0 for θ_w = 53.850), h/L = 0 for θ_w = 540 is a numerical result (see Figure 25n) Δ , experimental point; • numerical results.

SECTION 1

INTRODUCTION

A direct comparison is made for fifteen basic cases of oblique snock-wave reflections between interferometric results obtained at the University of Toronto Institute for Aerospace Studies (UTIAS) 10 cm x 18 cm Hypervelocity. Shock Tube and numerical results obtained by using a current computational method for solving the Euler equations of compressible flow. Additional parametrized sequences of calculations are presented to assess the utility of the present numerical method in constructing the various reflection-transition lines (RR - SMR, SMR - CMR, CMR - DMR; see Figs. 1 & 2) for inviscid flows in the shock-wave Mach-number, wedge-angle (Ms, $\theta_{\rm W}$) -plane. An additional parametrized sequence has been calculated in order to study the validity of the boundary-layer displacement theory to account for the "von Neumann paradox."

Over the past five years, extensive experimental and analytical data were obtained for these problems (Ben-Dor & Glass 1978, 1979, 1980; Ando & Glass 1981; Lee & Glass 1982; Shirouzu & Glass 1982; Deschambault & Glass 1983; Deschambault 1984; Hu 1984; Hu & Glass 1985; Hu & Shirouzu 1985; Wheeler and Glass 1985; Wheeler 1985). We refer the reader to these references for an extensive discussion of the theory of oblique shock-wave reflections, an introduction to the history of the field, and further references.

With the advent of modern computers, it has become possible to attempt the computation of such problems using finite difference schemes. The state-of-the-art in this area was surveyed in Ben-Dor & Glass 1978 and Deschambault & Glass 1983; these authors concluded that advances in numerical technique would be required before numerical results could be viewed with the same confidence as experimental data. The main object of this report is to demonstrate that the numerical method used herein is sufficiently accurate to be placed on a nearly equal footing with experimental methods in the analysis of perfect, inviscid, compressible flows. However, further development work is needed in the numerical modelling of nonequilibrium, viscous flow fields.

This study deals exclusively with results for air and argon; the experimental data for these results may be found in Deschambault 1984 where they are discussed in detail. Many other related calculations have been performed with our computer code, and we briefly describe them here. Recent experimental data for ${\sf SF}_6$ has also been obtained and is reported in Hu 1985

and Hu & Glass 1985. An analogous numerical-experimental study to the present report for SF $_6$ may be found in Glaz et al 1985. An interesting problem is bosed by assuming a polytropic gas, fixing M $_5$ and $\theta_{\rm w}$, and allowing veration of specific neats to be a varying parameter. This problem is well—suited to a numerical study and the results using our numerical method are presented in Colella & Glaz 1984 and Rerger et al 1985. Finally, a computer code has been developed for the problem of a spherical explosion reflecting off an ideal surface; the numerical method is virtually identical to that used in obtaining the results for this report. Calculations using this code are presented in Colella & Glaz 1984 and Colella et al 1985. The results of these calculations show that the high M $_5$, DMR flow fields of planar oblique shock wave reflection have a lot in common with the Mach stem region flow fields of the spherical explosion problem just after the RR-DMR transition, although there are significant structural differences downstream of the triple point, presumably due either to unsteadiness or the different boundary conditions.

A portion of the calculations which are studied in this report have appeared in Colella & Glaz 1982, 1984 and Glaz et al 1985. The latter paper includes an expanded discussion of some of the overall issues involved in comparing experimental results to approximate solutions of a perfect inviscid flow. In this report, we concentrate on presenting the complete set of calculations including a discussion of each comparison or parametric series. The plan of the report is as follows. In Section 2, the terminology of oblique shock-wave reflections is reviewed and some notation is defined. Sections 3 and 4 are devoted to experimental techniques and the numerical method, respectively. In Section 5, the results are presented and Section 6 is an extended summary.

SECTION 2

DBLIQUE SHOCK-WAVE REFLECTIONS

The four types of pseudo-stationary oblique snock-wave-reflection batterns are shown in Figure 1 and consist of (a) regular reflection (RR), (b) single Mach reflection (SMR), (c) complex Mach reflection (CMR) and (d) double Mach reflection (DMR). Figure 1 illustrates the definitions of wedge angle $\theta_{\rm W}$, triple-point trajectory angles, $\chi_{\rm X}\chi'$, various snock waves I. R. R', M, M', slip surfaces S,S' and the flow regions 1-5 produced by the foregoing reflections, the angle $\theta_{\rm W}$ between the incident I and reflected R shock waves is also shown as well as the angle $\theta_{\rm W}$ between R and the wall or R and the triple-point-trajectory angle $\chi_{\rm W}$. The bow shock stand-off distances s and the length L, between the wedge corner and the reflection point or Mach stem are also indicated. Such qualities can be measured experimentally or predicted numerically and provide important information on the state of the gas whether frozen, non-equilibrium or equilibrium (Shirouzu & Glass 1982; Hu 1985; Hu & Glass 1985).

The equations of gas dynamics are, in Cartesian coordinates,

$$\rho_{t} + (\rho u)_{x} + (\rho v)_{y} = 0$$

$$(\rho u)_{t} + (\rho u^{2} + p)_{x} + (\rho u v)_{y} = 0$$

$$(\rho v)_{t} + (\rho u v)_{x} + (\rho v^{2} + p)_{y} = 0$$

$$(\rho E)_{t} + (\rho u E + u p)_{x} + (\rho v E + v p)_{y} = 0$$
(1)

where p is the density, $\underline{u} = (u,v)$ is the velocity field, $E = \frac{1}{2}(u^2 + v^2) + e$ is the total specific energy, e is the specific internal energy, and p is the pressure. The system is closed by specifying an equation-of-state (EOS).

$$p = p(\rho, e). \tag{2}$$

We shall often use the polytropic EOS,

$$p = (\gamma - 1)pe, (3)$$

where y > 1 is the ratio of specific heats.

If real gas and viscous effects can be ignored [i.e., equations 1), (2) hold], the problem has no intrinsic length-scale, suggesting the use of the self-similar or pseudo-stationary coordinate system $(\xi,n) = [(x-x_0)/(t-t_0), (y-y_0)/(t-t_0)]$ where (x_0,y_0) are the coordinates of the wedge corner and t_0 is the time at which the incident shock wave reaches the corner. Following Jones et al 1951, the system (1) may be transformed to pseudo-stationary coordinates and becomes, in conservation form,

$$(\rho \vec{u})_{\xi} + (\rho \vec{v})_{\eta} = -2\rho$$

$$(\rho \vec{u}^{2} + p)_{\xi} + (\rho \vec{u} \vec{v})_{\eta} = -3\rho \vec{u}$$

$$(\rho \vec{u} \vec{v})_{\xi} + (\rho \vec{v}^{2} + p)_{\eta} = -3\rho \vec{v}$$

$$(\rho \vec{u} \vec{H})_{\xi} + (\rho \vec{v} \vec{H})_{\eta} = -\rho (\vec{u}^{2} + \vec{v}^{2}) - 2\rho \vec{H}$$
(4)

where

$$\tilde{u} = u - \xi, \tilde{v} = v - \eta,$$

$$\tilde{H} = \frac{1}{2}(\tilde{u}^2 + \tilde{v}^2) + h$$
(5)

and h = e + p/p is the specific enthalpy. We refer to (\tilde{u}, \tilde{v}) , \tilde{H} as the self-similar velocity field and self-similar total enthalpy, respectively. In addition we define

$$\tilde{M}^2 = (\tilde{u}^2 + \tilde{v}^2)/c^2 \tag{6}$$

where c = sound speed and we refer to \tilde{M} as the self-similar Mach number. The system (4) is, evidently, the steady Euler equations with the addition of source terms. We note that the ratio s/L is constant, for given initial conditions, for self-similar solutions of the non-stationary equations, just

as s is constant for steady supersonic flow. In this and other ways a change to pseudo-stationary coordinates is very useful in the analysis of these flow fields and will be used in this study.

In particular, the type of reflection pattern is a function of the incident shock-wave Mach-number M_s , the wedge angle θ_w , and the gas equation of state. The transition boundaries in the M_s , θ_w -plane for oblique shock-wave reflection are reproduced from Lee and Glass (1982) in Figure 2 for real air and a polytropic equation of state with $\gamma=1.40$. The analogous figure for argon ($\gamma=5/3$) may be found in this reference. The construction of the transition lines is based on various (heuristic) transition criteria and the numerical calculation of the jump conditions at reflection and triple points. These criteria, which have been the subject of extensive investigation in the literature, are summarized in Lee and Glass (1984) and Shirouzu and Glass (1983). In Sec. 5, the numerical results will be used to partly assess the validity of some of these criteria as well as the overall accuracy of the transition diagram, Figure 2.

The fourfold partition of the (M_S, θ_W) plane illustrated in Figure 2 is quite coarse relative to the rich phenomenology present in these flow fields. Some other features that may be similarly partitioned (see Ben-Dor and Glass 1979) are (a) whether or not the reflected shock is detached or attached to the wedge corner; (b) in the attached case, whether the flow at the corner is subsonic or supersonic; (c) for RR whether the flow is subsonic or supersonic (in pseudo-stationary coordinates) at the reflection point and (d) for SMR, CMR and DMR whether or not M "toes-out" or "toes-in".

A comprehensive study of these issues is beyond the scope of this report, but they will be discussed as appropriate in the comparison of experimental and numerical results in Sec. 5.

SECTION 3

FXPFRIMENTAL TECHNIQUES

The experiments for the present study were performed in the UTIAS $10 \text{ cm} \times 18 \text{ cm}$ Hypervelocity Shock Tube. A design, performance and calibration study of the original facility can be found in Boyer 1964. More recent and detailed descriptions of the shock tube appear in Bristow 1971 and Ben-Dor and Whitten 1979. Further details of the experiments associated with the present work can be found in Deschambault 1984.

3.1 EXPERIMENTAL FACILITY.

The basic shock tube facility consists of a 1.4m long driver and a 12.2m channel. The initial pressure in the channel can be easily varied from near vacuum to atmospheric conditions. At the end of the channel is a test section containing high-quality interferometric windows through which the shock tube flows may be observed. A 23-cm diameter field-of-view Mach-Zehnder interferometer (Hall 1954) in conjunction with a giant-pulse ruby-laser is used to record simultaneous dual-wavelength (λ =694.3nm and 347.2nm) infinite-fringe interferograms of the two-dimensional flow-fields. This allows the direct observation of the flow-field isopycnics (lines of constant density). The 15ns pulse generated by the ruby laser effectively freezes all motion, thereby producing sharp, clear images.

Two methods were used to produce the incident shock-wave Mach-numbers for the present study. For shock-wave Mach-numbers less than 6 a cold-gas driver was employed. The diaphragm consisted of several layers of mylar-polyester films. With the proper choice of driver gas, ${\rm CO_2}$ or He, and diaphragm thickness, the desired shock-wave Mach-number could be obtained in the test gas upon rupture.

For shock-wave Mach-numbers greater than 6 combustion-driver techniques were used. Specially scribed stainless steel diaphragms were burst by the constant-volume combustion of a stoichiometric mixture of 0_2 and H_2 diluted with 70% He. Combustion was initiated by the impulsive heating of a 0.38-mm diameter tungsten wire through the discharge of a $45\mu F$ 13kV capacitor.

The reflection patterns were generated by the impingement of normal shock

waves with steel wedges. The wedges were bolted firmly to the bottom wall of the facility to ensure rigidity. The sides of the wedges were flush with the inside walls and interferometric windows of the shock-tube test-section

3.2 DATA REDUCTION TECHNIQUES.

The infinite-fringe interferograms enabled the recording of small relative density changes of the various snock tube flows. The density difference $\Delta\rho$ between the two adjacent fringes of the same color is related to the wavelength λ of the light source (694.3nm and/or 347.2nm) and the Gladstone-Dale constant K (2.274 x $10^{-4} \text{m}^3/\text{kg}$ for air, λ = 589.6nm and 1.574 x $10^{-4} \text{m}^3/\text{kg}$ for Ar, λ = 694.3nm) and is expressed by the relation $\Delta\rho$ = λ/KL , where L is the depth of the test section (10.16cm).

To obtain quantitative values for the isopycnics the following method was employed. From the initial conditions of the experiment, i.e., shock-wave Mach-number, wedge angle, initial pressure and temperature, the thermodynamic states around the reflection point for RR and the triple point for MR were calculated using two- and three-shock theory (Ando 1981). These were used as reference states from which all other density values could be obtained using the above relation.

The wall-density distribution plots were obtained directly from the interferograms. The origin was defined to be the reflection point of a RR or the foot of the Mach stem of a MR. The corner of the wedge was defined to be a distance L from the origin. All absolute distances were then scaled by L giving a value of 1 to the distance from the origin to the wedge corner. Where possible the center of the isopycnic intersecting the wedge surface was used to locate the value of the density at that point.

For some of the experimental results presented here, it was necessary to use test gases with very low initial densities and pressures relative to atmospheric conditions. As a result, several interferograms show the effects of vibrational nonequilibrium which must be taken into account when analysing the corresponding interferograms. The relaxation zones are clearly visible and appear as additional fringe shifts in the post-shock flow-field parallel to the frozen incident shock front. Behind the reflected shock wave, the characteristic signature of a relaxing gas is the nearly tangential incidence of the isopycnics and the reflected shock wave.

SECTION 4

NUMERICAL METHOD

The numerical results presented in this paper have been calculated with a version of the Eulerian second-order Godunov scheme for nonstationary gas dynamics of a type considered by Colella and Woodward 1984. The version of the scheme used here is presented in Colella and Glaz 1982,1983, including the modifications required for non-polytropic gases.

The method is a finite-difference scheme for systems of hyperholic conservation laws in one space-like dimension: for multidimensional applications such as the shock-on-wedge problem, we employ operator splitting. Differencing is in conservation form and the numerical fluxes are computed by solving zone interface Riemann problems whose time-centred left and right states are computed from the characteristic form of the equations. This technique leads to second-order accuracy in smooth flow and ensures that the method is centred upstream. In practice, the method is very stable and robust. In the immediate vicinity of a strong shock, some dissipation is required; this has been accomplished by smoothly degrading the scheme to the first-order Godunov scheme in such regions. The degree of degradation is a function of the shock thickness and strength.

For argon, we have used a perfect (frozen) gas equation of state with $\gamma=5/3$. If the shock tube test gas was air, the equation of state was chosen to be either a perfect (frozen) gas with $\gamma=7/5$ or the Hansen 1959 real air equation of state as modified by Deschambault 1984 for the present application. The efficient solution of the Riemann problem in the context of our second-order Godunov method for an arbitrary equation of state is treated in Colella and Glaz 1982,1983. Also, these papers demonstrate that the choice of equation of state has a substantial influence on the quantitative numerical results, as might be expected.

As noted in the preceding section, vibrational non-equilibrium, which is only temperature dependent, can be significant for moderate to high Mach numbers when the test gas is air (at high Mach numbers dissociation effects are also density dependent): for the argon cases considered here we expect the gas to remain frozen. The choice of an appropriate equation of state for the air calculations depends mainly on the vibrational relaxation length $l_{\rm V}$, behind the shock waves I, R, M of Figure 1. If $l_{\rm V} > 1$ (where 1 is a characteristic flow length arising in the problem; for the present

experiments, 1 ~ J.1mm), then the gas is frozen and the perfect gas equation of state is correct. If $1 > 1_{\rm V}$, then the gas is in equilibrium and the Hansen equation of state for real air is used. Finally, if $1_{\rm V}$ ~ 1, then neither the frozen nor the equilibrium hypothesis is appropriate, and the flow is said to be in non-equilibrium. We have numerically treated such cases as equilibrium flow fields by using the Hansen equation of state, although the only correct procedure would be to solve an extra partial differential equation representing a rate equation for vibrational relaxation (and for dissociation at high Mach numbers). This decision will be an important issue in our discussion of these cases in Sec. 5.

The computational mesh and our problem initialization procedure is illustrated in Figure 3. Note that these figures are drawn from right to left to conform with the experimental interferograms. We have used a square (i.e., $\Delta x = \Delta y = \text{constant}$) mesh for all of the computations in Sec. 5. Because the flow is pseudo-stationary, the choice of Δx is immaterial.

The initial data are taken as U_0 , M_s where $U = (\rho, \rho, u, v,)^T$ is the state vector and M_s is the initial shock-wave Mach-number. From these data and the given equation of state, the post-shock state U_1 may be calculated. To initialize the two-dimensional calculation, these data are placed on the grid far upstream (ca. 60-75 zones) of the corner, as illustrated in Figure 3a; interpolation of conserved quantities [i.e. $U^{C} = (\rho, \rho u, \rho v, \rho E)^{T}$] is used for zones that straddle the incident shock. However, this is a very poor representation of the numerical shock because any shock-capturing scheme will diffuse a shock wave over two or more zones in the computational mesh. The resulting structure is referred to as a discrete travelling wave (i.e., a mesh function that depends only on x - Vt, where V is the vector velocity of the wave and equals the shock speed in magnitude for a discrete shock wave). Starting with any initial data (e.g., the one zone U_0 - U_1 jump described above) satisfying the Rankine-Hugoniot conditions, the solution will tend as the number of time-steps becomes large towards the appropriate discrete travelling wave, plus other low-amplitude waves that we refer to as "starting error", with the starting error separating from the travelling wave. For the present application, it is very important to ensure that the starting error is eliminated before the shock wave is allowed to reflect, and we proceed as follows. First, the computer code is allowed to run until the shock wave

reaches the corner, and the situation in Figure 3b is reached. In this figure, the region immediately behind the shock and about 2+3 zones thick is the discrete travelling wave and the small (less than 5%) relative amplitude disturbances further downstream is the one-time starting error. The computer code then arbitrarily changes the flow field to that illustrated in Figure 3c, i.e., the discrete travelling wave (arbitrarily set to exactly 4 zones in the computer code) is retained but the starting error is replaced by the post-shock state U_1 .

At this point, the flow field becomes truly two-dimensional and the computer code is now run without further interruption until the end of the calculation is reached.

The boundary conditions for this problem, which are standard, are discussed in detail in Colella and Glaz 1983. We remark here that our treatment of the intersection of the incident shock with the upper or left-hand boundary or both is not entirely consistent with the discrete travelling wave and leads to the introduction of a low relative amplitude (ca. 1%) wave behind the incident shock at its intersection with the boundary. This wave, which we call a boundary error, may lead to a rather unaesthetic structure in the contour plots and it can impinge on the disturbed flow field behind the reflected shock. Examples will be noted in Sec. 5.

All calculations were performed on a CRAY I at Los Alamos National Laboratory, Los Alamos, New Mexico. The computer code was designed to take significant advantage of the machine's vector architecture. Each calculation in Sec. 5 required 15-40 min. c.p. time with most in the range of 20-30 min. Much of this time is wasted on the extra grid points introduced to eliminate the starting error as well as grid points outside the reflected shock. Also a fine mesh is only really needed in the Mach-stem region. Thus, an intelligent adaptive mesh structure could reduce these times substantially.

JECTION 5

COMPUTATIONAL RESULTS

A direct comparison of experimental results and numerical calculations is presented in Sec. 5.1 for fifteen cases of oblique shock-wave reflections. For eight of the cases in air, the computation has been performed twice, once with a perfect gas EOS with $\gamma=1.4$ and once with the Hansen EOS. Thus, twenty-three computations are reported on in this part. In Sec. 5.2, the results of several parametrized sequences of calculations are presented to demonstrate the capability of our numerical method to compute the correct transition in the (M_S, θ_W) -plane. An additional sequence is presented in this part to demonstrate (upon comparison with experimental data) the effect of boundary-layer displacement on the RR-DMR transition.

5.1 COMPARISON OF EXPERIMENT WITH CALCULATION.

The initial conditions for the fifteen cases are listed in Table I along with the computational mesh (NX,NY) and the equation of state selected for each case (and it is noted where two choices of EOS were made for a case). All four wave configurations are represented in the range of (M_e, θ_u) considered. The following data are presented for each case: experimental isopycnics; computed isopycnics using the same density levels as were obtained in the experiment; wall distribution plots, q vs x/L, with q = p/p_0 , ρ/ρ_0 , e, u and with the ρ/ρ_0 plot including a comparison with experiment; whole flow field contour plots, using thirty equally spaced contours, of the quantities p,e,p,M,u,v,u,v,H; in a "blowup" frame in the vicinity of the Mach stem or reflection point, contour plots, using thirty equally spaced contours of ρ ,e, ρ ,M,H are shown along with the experimental isopycnics, self-similar streamlines, and a self-similar velocity vector field plot. For those cases involving a comparison of two calculations with differing EOS, the contour plots of actual isopycnics are shown together with the interferogram, and an additional wall distribution plot is added comparing ρ/ρ_0 vs. x/L for the two calculations and the experiment on the same graph.

In order to assist the reader in interpreting the graphical output, we make several general comments here, which are not repeated below for each case. It is regretted that many interesting phenomena are not commented on in the text, but we felt it useful to present the entire set of figures. First, for those cases involving two calculations, the subscript "P" refers to the

perfect gas calculation and the subscript "H" refers to the Hansen calculation; in the event that such a figure is referenced in the text without a subscript, the context determines which 'or both' figures are being discussed. Concerning the contour plots, the coordinate system is priented with the origin at the corner point, the k or 5 direction along the weage surface after reflection, the yor holicection perpendicular to the wedge and facing upwards in the figure; however, we have reversed prientation for the wall plots and have set c/L = 1.0 at the corner and c/L = 0.0 at the reflection point or the intersection of the Mach stem with the wedge surface. We regret that we have not matched the length L in the plots of the calculations with the corresponding interferograms. The contour plots of those quantities which may take on positive or negative values use solid (dashed) lines to represent positive (negative) contour levels. The zero level is always the last solid contour. In particular, the sonic line in the M plots can always be easily found. The most important feature of an equally spaced contour plot for compressible flow is that discontinuities are clearly visible because several contour levels overwrite each other on the plot, at the location of the discontinuity. When plotting density contours using the levels prescribed by the available experimental fringes, this effect is still present but degraded to varying degrees for the different cases. In particular, density levels between ρ_1 and ρ_2 , etc. may not be present at all, although in many cases we arbitrarily inserted extra contours for aesthetic reasons.

Referring to Figure 7 (i.e., Case 4), the generic features of the various plots are discussed (the notation "Figure N" is used when it would not be useful to consider Figure 7 as an example). First, most pages have a heading with certain information: "MS", "ALP" are M_s , θ_w in the notation of the text; NR and NZ are the number of mesh points in the calculation and correspond to Table I; " P_0 " is P_0 , the initial shock-tube pressure; KBEG is the first point (viewing from right to left) in the x-direction after the reflection point; and the word "PERFECT" or "HANSEN" appears to denote the EOS. Notice that (NR-KBEG) by NZ is the appropriate aspect ratio, rather than NR by NZ. Figures 7a and b are presented on the same page and all plots are uniformly labelled according to the table appearing along with these figures. When comparing an interferogram (Figure 7a in this example) with the associated calculated isopycnics (Figure 7b here), the effect of an EOS mismatch in the

talculation of amphandamph can be striking and misleading. Recall, see Jection 3, that these gensity values are calculated for the interferogram using a specific choice of aguilibrium EOS. A numerical computation using a different choice of EOS will automatically get different values and shift all of the contours away from their correct locations. An excellent example of this effect is Figure 12by where most of the isopychics have been shifted into the numerical snock layers associated with the reflected snock and the second Mach stem. Concerning Figures To and a, we note the following: 1' the M plots in both figures show that the disturbed flow is subsonic everywhere except in Region 2 where the flow is entirely supersonic, (2) contact or slip surfaces tend to show more clearly in the plots of e and H than in p and this is especially true of the boundary of the vortex rollup as may be noted by comparing these three plots in Figure 7e; of course, this effect is caused by the different effects the Rankine-Hugoniot jump conditions have on the number of contours appearing in the shock layers for the different quantities and, it should also be noted, the EOS or the value of y in the case of the polytropic EOS has a large effect, (3) shock waves can be distinguished from slip surfaces by comparing the pressure plots with plots of o, e, M, etc., (4) distinguishing compressions from rarefactions can usually be done with the pressure contours alone (e.g., if a compression steepens to form a shock) or in conjunction with the wall pressure plot, Figure 7c; determining the direction in which a wave faces is often of interest and is not usually obvious although sometimes the u plot in Figure Nc can be used for waves normal to the wall, (5) the H plots are not constant states because of the source terms in eqn. 4; it follows from the Rankine-Hugoniot conditions that H does not change across a shock wave and this can be seen clearly in Figures 7d and e for the second Mach stem and the reflected shock wave, although there is some slight nonmonotone variation inside the numerical shock layer. It is also true that H does not jump across the incident shock wave and the first Mach stem despite appearances to the contrary in Figures 7d and e. Close inspection (not obvious to the reader in most instances in Figures Nd and e) reveals that H is the same in each of Regions 0-3 at the first triple point and the variation of H in the numerical shock layers is enough to cause this layer to be filled in with several additional contour levels, (6) the visual appearance in Figure 7c of streamlines ending in the interior of the calculation is, of course, just a plotter error (the density of

streamlines is fixed for reasons of efficiency, 7° the velocity vector plot, figure 7° shows now this vector jumps across shock waves and aligns itself with slip surfaces and the wall boundary conditions.

The discussion of many of the cases refers to hand measurements of χ and χ' . The accuracy of these measurements is not usually very good; however, differences between measurements (e.g., regarding the two calculations of the same case using different choices of EOS) is much more reliable.

Case 1: $M_s = 2.05$, $\theta_w = 60^{\circ}$, RR, Argon. Comparison of the experimental and numerical isopycnics (Figures 4a, b) show them to be in good agreement with an error of about one fringe at the start of the subsonic region. The wall density distribution (Figure 4c) disagrees by about the same amount. It may be observed that the density contour levels curve sharply towards the reflection point just above the wedge surface, an effect that is not present in the experimental results. The blowup plots of M, H exhibit this effect as well, even in the supersonic region. This numerical error is referred to as "wall heating" and is commonly observed in shock capturing calculations as shown, for example, in Noh 1976. Wall heating affects only the density, temperature, etc., and not the pressure (Figure 4d). It may be seen to account for part of the observed error in this case, including the slight error in the value of the reflected shock wave density ρ_2 on the wall. In addition, the error in the stand-off distance of the bow shock s, relative to the experimental distance from the reflection point P to the corner L is about 6.2%.

Case 2: $M_S = 1.26$, $\theta_W = 450$, RR, Air. Figures 5b and c show that the quantitative agreement between experiment and calculation is very poor for this case, and the results are largely independent of the choice of EOS. Furthermore, the angle between the wedge and the reflected shock as well as other gross flow field quantities are in substantial error, even though the isopycnic patterns are in excellent qualitative agreement. A possible explanation for this severe error may be found by considering the ratio ρ_2/ρ_0 as a function of M_S , fixing $\theta_W = 450$. It turns out that ρ_2/ρ_0 and ρ_2/ρ_0 and ρ_2/ρ_0 and ρ_2/ρ_0 , whereas the calculation has $\rho_2/\rho_0 = 2.2$. In other words, the slope $d(\rho_2/\rho_0)/dM_S$ is so steep in the

region of interest that small errors in either the numerical method or experimental measurement can lead to large errors in ρ_2 . For further discussion of this type of consideration, see Hu and Shirouzu 1985. Noting that the distruped flow field is wholly subsonic, Figure δa , and that this case is close to the RR-SMR transition boundary, it is also possible to speculate that the experiment is actually an SMR despite this not being visible on the interferogram or in the calculations. The full resolution of this disagreement should be possible with the adaptive mesh version of our code, see Berger et al 1985, used in a region of parameter space around $(M_s, a_w) = (1.26, 450)$. A posteriori, it is seen that studying the results for both choices of EOS was not useful.

Case 3: $M_s = 1.50$, $\theta_w = 45^\circ$, SMR, Air. Comparison of the interferogram and the calculated isopycnics, Figures 6a and b, shows excellent qualitative agreement and approximately a one fringe error quantitatively. This agreement continues for x/L = 1.0 since the experimental corner flow field is inviscid. We measure $\chi = 1.0^\circ$ and 0.5° for the calculation and experiment, respectively. Since the Mach stem is only 3-4 computational zones high, possible explanations for this disagreement include numerical error due to lack of resolution and the existence of viscous boundary-layer effects in the experiment. The disagreement in the wall density profiles, Figure 6c, for x/L < 0.5 may be due to several causes: the possibility of viscous effects along the wedge, the possibility of the numerical wall-heating error interacting with the slip surface, and the difficulty in precisely locating the intersection of a fringe with the wedge in the interferogram. Some of the contour plots in the blowup frame, Figure 5e, illustrate the difficulties. The different choices of EOS proved not to be important for this case.

Case 4: $M_S=3.03$, $\theta_W=470$, DMR, Air. The calculated and experimental wave patterns, Figures 7a and b, are in excellent qualitative agreement, including a relatively sharp slip surface emanating from the second triple point. The interferogram shows a different orientation for fringe c and an extra fringe d under the reflected shock between the two triple points, which may be an indication that the gas is relaxing in this region. The differences due to the choice of EOS are small, but noticeable. In particular, the values of χ and χ' are close to the experimental result for

the Hansen EOS but are too large by about 1° for the v = 1.4 results. It should be noted, however, that $\theta_{\omega}=470$ is very close to the RR-DMR transition ine and the boundary layer defect may have had some effect on χ_{*X}^{*} in the experiment. Also, the vortex rollup is closer to the leading Mach stem for the Hansen calculation than for the perfect gas calculation, Figure 7e; the interferogram does not show the rollup moving ahead at all, presumably due to viscous effects. As is typical for DMR results, the flow field is of mixed type with region 2 being supersonic and the remainder being subsonic, Figure 7d. Also typical is the relative strength of the contact surface and vortex rollup in the $\bar{\mathrm{H}}$ plots, Figure 7e; the waviness of this surface in the numerical results is a hint of the physical Kelvin-Helmholtz instability apparent in the interferogram. Concerning the wall density plots, Figure 7e, the Hansen EOS calculation is a few percent high on the peak value even after correcting the discrepancy in ρ_3/ρ_0 . The interferogram, of course, cannot exhibit the sharp inviscid peak and valley in the rollup region of the calculation. The relative displacement of these structures between the two calculations follows directly from the differences in the calculated values of x'. The viscous corner region, as expected, is not reproduced well in the calculations.

Case 5: $M_s = 2.65$, $\theta_w = 30^\circ$, CMR, Air. Comparing the interferogram with the density contours in Figures 8a and d, it may be seen that excellent overall agreement was obtained for the wave system except in the corner region. The comparisons using the experimental isopycnics, Figures 8a and b, differ by a larger degree. The differences in the vortex rollup pattern are clearly the result of experimental viscous effects. Figures 8d and e show a small supersonic region at the triple point, which is typical of CMR results. Another interesting feature of this flow field is the presence of three points along the wedge surface where u = 0 (see the u contour plot. Figure 8d, and the \tilde{u} wall plot, Figure 8c); the first two occur at the leading and trailing edges of the vortex rollup pattern and the third appears much further downstream. This pattern is pervasive (except, see Case 7) for those Mach reflections with a vortex rollup. We measure $\chi = 8.60$ for the perfect gas calculation and $\chi = 8.30$ for the Hansen EOS calculation and the experiment. The calculated values of ρ/ρ_0 in the region 0.0 < x/L < 0.2 are in good agreement with the interferogram, Figures eta_{a} and c, once the

talculation of s_3/ρ_0 is corrected for χ and the choice of EOS, and the different rollup patterns are taken into account. The larger disagreements in comparing the experimental isopycnics, Figure 3a, and the wall density distributions, Figure 3c, downstream of the vortex might be explained by the viscous effects providing different boundary conditions for the subsonic inviscid flow field between the rollup and the corner region (where these effects are substantial).

Case 6: $M_S = 5.07$, $\theta_W = 30^\circ$, CMR, Argon. The isopycnic patterns are in excellent agreement, despite the availability of relatively few fringes, Figures 9a and d, except for the corner region and the details of the vortex rollup pattern. The quantitative agreement, Figures 9a, b and c and measurements of χ , are also very good except in the corner.

Case 7: $M_s = 10.37$, $\theta_w = 10^{\circ}$, CMR, Air. The experimental results, Figure 10a, show strong relaxation effects in the disturbed flow field behind the reflected shock (this is indicated by the near tangential incidence of the fringes to the shock), and the incident shock jump appears almost in equilibrium. Also, the wedge surface does not appear to be perfectly straight in the photograph, which indicates that the sidewall boundary-layerdiffraction effects may be significant. There is reasonably good qualitative agreement (disregarding the real-gas effects) in the isopycnic patterns, Figure 10d, although the kink is more pronounced in the experiment than in the calculation. In evaluating the wall density plots, Figure 10c, it should be noted that the data points were evaluated assuming frozen-triple-point conditions while the calculation implicitly used the equilibrium Hansen EOS for the same task. Also, we estimate χ ~ 13.00 for the experiment and measure $\chi = 15.0^{\circ}$ for the calculation; the corner attachment angle is 20.5° for the experiment and 25.5° for the calculation. The latter difference is very large and is clearly the result of the difference between an equilibrium shock jump and a strongly relaxing shock jump at the corner. The former difference is probably also a real gas effect, and would have a strong influence on the kink structure. The vortex rollup patterns are in remarkably close qualitative agreement, although we note the rollup is closer to the leading Mach stem, which has a somewhat greater toe-out, in the experiment than in the calculation. An unusual feature of this flow field is that u has

just one zero on the wedge surface, located at the leading edge of vortex rollup. Comparing with the discussion in Case 5, this suggests that as the (M_s, θ_u) - plane is traversed from the low M_s , high θ_u region to the high $M_{\rm e}$, low $a_{\rm in}$ region and restricting to cases for which a vortex rollup pattern is present, the number of zero crossings of u along the wedge surface smoothly bifurcates between one and three. The results for Cases 14 and 15 substantiate this conjecture; the former lies near the transition point and has three zero crossings while the latter lies just beyond the transition (and the lone zero crossing ahead of the vortex rollup is pushed foward into the shock layer). Also, the contact surface is more unstable in the calculation than in the experiment. These two effects are opposite to those usually holding in our results. Thus, it seems that quite good quantitative agreement could be obtained for x/L small in the wall density plots if the rollup patterns could be spatially lined up, the Hansen EOS used in evaluating the data, and the corner jump conditions changed to provide the correct downstream boundary condition for the subsonic portion of the flow field, Figure 10d. The dip in Figure 10c at $x/L \sim 0.25$ is due to the boundary error.

Case 8: $M_s = 1.66$, $\theta_w = 400$, SMR, Air. The isopycnic patterns, Figures 11a, b and d, as well as the wall density plots, Figure 11c, are in excellent qualitative and quantitative agreement. The only noticeable difference between the two calculations is that the value of ρ_3/ρ_0 is in better agreement when using the Hansen EOS and this aligns the overall wall density plots closer to the experiment. There is a larger error for x/L large which is probably explained by viscous effects in the corner region for the experiment. The EOS effect on the values of ρ_3/ρ_0 is worth commenting on in detail, since the small value of M_{ς} precludes significant real-gas effects. Assuming a perfect gas and a Mach stem normal to the wedge surface at the triple point, one may compute $\rho_3/\rho_0 = [(\gamma+1)M_0^2]/[(\gamma-1)M_0^2 + 2]$ where $M_0 = \frac{1}{2}$ M_s csc ψ_0 and $\psi_0 = \pi/2 - (\theta_w + \chi)$ which implies that ρ_3/ρ_0 is sensitive to the value of χ at low shock-wave Mach numbers and/or high values of $(\theta_{\omega}^{+}+\chi)$ and that $d(\rho_3/\rho_0)/d\chi < 0$. For this case, $\chi = 3.50$ in both calculations but is slightly less in the Hansen calculation which is enough to account for the wall density results presented in Figure 11c. Noting that the calculations compute ρ_3/ρ_0 and account for the EOS, χ , and any deviation from normality of the Mach stem automatically, while the experimenter must make the assumptions

above and measure χ by hand, one sees that the differences in the various results are outweighed by the agreements. Also, boundary layer-displacement may be a factor because of the relatively low value of χ .

Case 9: $M_s = 2.87$, $\theta_w = 400$, DMR, Air. Comparing the density contour plots, Figure 12d, with the interferogram, we see that there is excellent overall qualitative agreement for both calculations. This agreement is maintained only for the perfect gas calculation when comparing the experimental isopycnics, Figure 12b; of course, the experimental data reduction used a frozen triple point analysis. The vortex rollup pattern, the corner flow field, and the second triple point flow field differ considerably, however. Taking up the latter point first, we note that this case is near the CMR-DMR transition boundary, irrespective of the choice of EOS. Also, we measure $\chi = 5.30$ for the perfect gas calculation, $\chi = 5.00$ for the Hansen EOS calculation and $\chi = 4.50$ for the experiment. Thus, it is not unreasonable for the calculations to contain a much stronger second Mach stem and sharper second triple point than the experiment, which is close to CMR. The effects of boundary-layer displacement might also play a role. The experimental contact surface is very diffusive and this effect may prevent the vortex from moving forward towards the Mach stem, in conjunction with the presumed boundary-layer effects. The peak stagnation density (see the u vs. x/L plots in Figure 12c) behind the vortex is substantially higher in the calculations; we can conjecture that this is due to the sharper DMR structure and the nondiffusive contact surface of the calculation. In view of these effects and the differing boundary conditions at the upstream stagnation point and the corner, the wall density results, Figure 12c, are actually in good agreement. The Hansen EOS results must be corrected for the data reduction technique and there is otherwise little difference in the two calculations. In particular, the bunching of the fringes, in Figure $12b_H$, at the second Mach stem is due to the mismatch in Regions 2 and 3 between the experiment and the Hansen calculations.

<u>Case 10</u>: $M_S = 3.72$, $\theta_W = 40^\circ$, DMR, Air. The analysis for this case follows closely that for Case 9, although the interferogram, Figure 13a, is clearly DMR as are the calculations. It is likely that there is a relaxation fringe underneath the reflected shock between the two triple points. We

measure $\chi = 5.50$ for the perfect gas calculation, $\chi = 5.20$ for the Hansen EOS calculation, and $\chi = 5.00$ for the experiment; the differences in the measurements of χ' are similar. Referring to the M contour plots. Figures 13d and e, one sees that the sonic line is coincident with the second Mach stem; this always occurs in our clear DMR results and it is a useful criterion in distinguishing the CMR-DMR transition. Other typical flow field features are (1) the transition of the second Mach stem to a continuous compression near its intersection with the main contact surface, Figures 13d and e, and (2) the existence of two stagnation points Q_1 and Q_2 , one helind the vortex and the other just below the S-M' intersection and above the wedge, see the (u,v) vector field plots. Figure 13e; note that the self-similar streamlines are singular at these two points. Also, the pressure attains local maxima at these two points, Figure 13e, and the u contour plots, Figure 13d, show u = 0 at Q_1 and Q_2 . Concerning the wall density plots, Figure 13c, the agreement is closer than it appears because the data points in the range 0.18 < x/L < 0.425 need to be shifted to the right to account for the different relative locations of the second triple point; such a shift lines up the plots but the peaks are still off as in Case 9.

Case 11: $M_c = 4.62$, $\theta_w = 400$, DMR, Air. The analysis is similar to Cases 9-10. There is probably a relaxation fringe underneath the reflected shock which is stronger in the interferogram, Figure 14a, then for some of the other cases. We measure x = 6.00 for the perfect gas calculation and x = 5.00 for the Hansen EOS calculation and the experiment; there are similar differences for x'. The quantitative agreement between the experiment and the perfect gas calculation is very good away from the corner in both the isopycnic plot, Figure 14b_p, and the wall density plot, Figures 14c and c_p . A shift of data points as in Case 10 leads to nearly exact agreement for 0.16 < x/L < 0.35 and the peak density error at the stagnation point is very small, relative to Cases 9-10. This is perhaps due to a reduced relative influence of viscous effects in the region, although the interferogram shows significant instabilities in the contact surface and a vortex rollup pattern similar to these two cases. It is interesting to note that in Cases 9-11 the Hansen EOS provides better agreement with the experiment in terms of gross flow field features (e.g., $\chi_{,\chi}$ ') but worse agreement on quantitative details such as wall density curves (we are discussing the situation, of course, after

the data have been corrected for the choice of EOS). The present case exemplifies this fact in that the Hansen EOS calculation shows exact agreement on χ but is badly off on beak density along the wall.

Sase 12: $M_s = 2.03$, $\theta_w = 270$, SMR, Air. The agreement between calculation and experiment is extremely strong in all respects, Figures 15a, b and c, and is the best of all the fifteen cases. Quantitatively, the isopycnics are off by about one fringe and the wall density plot shows similar agreement except in a small region near the corner. The contact surface spreads out in the experiment and does not rollup as much as in the calculation.

Case 13: $M_s = 8.70$, $\theta_w = 270$, DMR, Air. The interferogram, Figure 16a, exhibits substantial real gas effects and even the Hansen EOS does not model the isopycnic shapes and locations very well. The relaxation length, t_{y} , is about 0.1 x L for the incident shock and the fringes are at nearly tangential incidence to the reflected shock. Also, the relaxing gases in the Mach stem region have obscured the contact surface and part of the roll-up pattern. The density contour plot, Figure 16d, and the interferogram show very good agreement. The rollup patterns substantially agree, although the contact surface normal to the wall S_n at x/L = 0.02 and the backwards facing shock wave W_h normal to the wall at $x/L \sim 0.065$ in the calculation, Figure 16e, are either not resolved or are lost due to viscous effects in the interferogram. Both calculation and experiment exhibit a strong toe-out of the first Mach stem; the kink on this shock surface may be near transition to a new triple point in view of the possible existence of an extra slip surface So emanating from this point, see the u contour plot, Figure 16d, and the M plots, Figures 16d and e. The vector field plot, Figure 16e, shows the existence of a pseudo-stationary stagnation point θ_2 near the intersection of the two slipstreams, in addition to the one at the center of the vortex O_1 ; indeed, there is a two-dimensional region around Q_2 where the flow appears to be stagnated. Our measurements show that $\chi = 9.60$ and $\chi' = 9.00$ for the calculation, and that $\chi = 7.50$ and $\chi' = 7.80$ for the experiment. The measured corner attachment angles are 33.5° and 23.0° for the calculation and experiment, respectively. This nonequilibrium effect (which apparently is poorly modelled with the equilibrium Hansen EOS) explains the large

disagreement near x/L=1.0 in the wall density plots, Figure 16c. After correcting the data for the Hansen EOS, there are large errors in the wall density plot in the range x/L < 0.5. Possible explanations include the large error in downstream boundary condition at x/L=1.0, the large difference in χ and χ' , viscous effects and differences in rollup pattern, and general relaxation effects, of course, the nonequilibrium flow field likely contributes to the other three effects. Overall, real-gas effects have an extensive impact on the flow field dynamics for this case and the equilibrium calculation was unable to reproduce many of the details.

Case 14: $M_s = 7.19$, $\theta_w = 20^{\circ}$, C/DMR, Air. The interferogram, Figure 17a, shows clearly that the experimental flow field is neither frozen nor in equilibrium, including the disturbed flow beneath the reflected shock. A more detailed discussion of equation of state and nonequilibrium effects in the numerical analysis of this case is available in Colella and Glaz 1985. The triple point angle χ is nearly in exact agreement, and the rollup patterns and Mach stem toe-out agree qualitatively, Figure 17d. The attached shock wave at the corner is hifurcated in the interferogram and supersonic in the calculation, a possible relaxation effect, although viscous effects may be important too. The experiment and calculation both show this case lying near the CMR-DMR transition boundary. Lee and Glass 1982 conjecture that this transition occurs when the sonic line just reaches the kink; the $\tilde{\text{M}}$ contours. Figures 17d and e, bear this out quite well. After allowing for the EOS correction of the wall data, a possible small shift for the vortex location, and the different corner structures, the wall density plots show surprisingly strong agreement; the dip at x/L = 0.35 is an excellent example of the computational boundary error.

<u>Case 15</u>: $M_S = 8.86$, $\theta_W = 20^{\circ}$, DMR, Air. The calculated density contours, Figure 18d, and the interferogram, Figure 18a, show good overall agreement, including many flow field details. The vortex rollup patterns are very close, although viscous effects in the experiment preclude detailed agreement. The vortex is pushed foward very close to the Mach stem in both calculation and experiment; the calculation shows a wave interaction W in this region which does not appear in any of the other cases. This is seen most clearly in the blowup plots, Figure 18e. The details of this portion of the

flow field are lost in the interferogram and are underresolved in the calculation. However, this flow field pattern is reproduced in the interferogram of Experiment 974 from Deschambault 1984 for which $\rm M_S=10.18$, $\rm M_S=10.18$, $\rm M_S=10.18$, $\rm M_S=10.18$. Vibrational relaxation effects are pervasive in the experiment, Figure 186. Vibrational relaxation effects are pervasive in the experiment, Figure 18a, including the Mach stem region. The failure of the fringes to merge into the second Mach stem as the contours do in the calculation, Figure 18d, is probably a real-gas effect. The corner attachment angle is 27° for the calculation and about 21° - 23° for the experiment. The calculation has a supersonic corner and relaxation effects dominate the experimental results in the corner region. We measure $\rm M_S=10.00$ for the calculation and $\rm M_S=100$ for the experiment, and $\rm M_S=10.00$ for the experiment. Overall, nonequilibrium effects preclude a realistic quantitative comparison for this case.

5.2 TRANSITION SEQUENCES.

Four sets of parametrized sequences of calculations are presented in this section. The purpose of the first three sets of calculations is to assess the potential of detailed computational results in constructing oblique shockwave-transition boundaries (see Figure 2) and in validating theories explaining these transitions. Each set contains two sequences of calculations, one for a perfect gas with $\gamma = 1.4$ and one using the Hansen EOS. The following data is presented for each case: whole flow field contour plots, using thirty equally spaced contours, of the quantities ρ , \tilde{M} ; in an appropriate "blowup" frame in the vicinity of the triple point or reflection point, contour plots, using thirty equally spaced contours, of the quantities p,e,p,M,H,u, along with the streamlines and vector field associated with the pseudo-stationary velocity (u,v). The purpose of the fourth set is to demonstrate the boundary-layer defect theory by presenting a parametrized sequence of inviscid calculations for argon (treated as a perfect gas with $\gamma = 5/3$) near the RR-DMR transition boundary and comparing with an experimental result. For this set, only whole flow field density contours are presented.

Set 1: Here, an attempt is made to locate the SMR-CMR and CMR-DMR boundaries for θ_w = 450, Air; 1.30 < M_S < 2.60, perfect gas with γ = 1.40;

1.50 < M_s < 2.30, Hansen EOS; in increments of Δ M_s = 0.1. The results are presented in Figures 19 and 20. Considering the M plots in the vicinity of Me = 1.70, we see that the sonic line has moved into region 2 for the cases with $M_{\rm g} > 1.70$ and that the extent of the supersonic region increases with increasing shock-wave Mach number. Assuming that the SMR-CMR transition occurs when region 2 becomes supersonic at the triple point (see Lee and Glass 1982), it follows that the $M_s = 1.70$ case is a CMR and the cases where 1.30 < M_c < 1.70 are SMR's because region 2 is entirely subsonic for these cases. It may be noted that for $M_{\rm S}$ = 1.30, the Mach stem M and the slipstream S are only barely visible and the case appears like an RR. The differences due to EOS effect are not marked at these M_{ς} values, but the Hansen M_c = 1.60 results provide a slightly earlier CMR than the perfect gas calculation. The results agree reasonably well with the analytic transition diagram, Figure 2. Also, it would not be unreasonable for the reader to view Figures 19 and 20 and take these transitions at slightly higher M_s values, which would have the effect of making the comparison with Figure 2 somewhat less close.

In view of the small values of χ in this region, it would be useful to restudy these cases with a refined mesh in the triple-point region (using an adaptive mesh algorithm, Berger et al 1985), thereby substantially eliminating the effects of numerical error near the wall boundary and allowing sufficient resolution to separate the results for the two choices of EOS. Also, the severe slope $d\theta_w/dM_s$ of the transition curves at $M_s = 1.70$ argues for increased resolution.

We now consider the ρ , \tilde{M} plots in the range 2.20 < M_S < 2.40. One theory for the CMR-DMR transition (see Lee and Glass 1982) is that the flow at the first triple point should be supersonic with respect to the motion of the kink. Because the flow immediately beneath the reflected shock and between the two triple points is constant, this criterion is equivalent to requiring that the sonic line (in pseudo-stationary coordinates) intersect the kink. Also, the sonic line should have the same tangent at the kink as the second Mach stem, because the flow is supersonic ahead and subsonic behind this discontinuity. Finally, the density contours may be expected to hegin coalescing as the shock wave is about to form. Using these criteria, the calculations show that the M_S = 2.30, perfect gas case is a weak DMR and that

the $\rm M_S=2.40$, perfect gas case is a clear-cut DMR; for the Hansen calculations, the $\rm M_S=2.20$ can be considered a DMR and the $\rm M_S=2.10$ case is a CMR. These results are also in reasonable agreement with the analytic results for the perfect-gas transition at $\rm A_W=450$, see Figure 2. Note that Figure 2 indicates that no DMR can exist in this range of $\rm M_S$ for $\rm A_W=450$. However, it has been found experimentally that the CMR-DMR transition line meets the SMR-CMR transition line where it joins the RR-MR line. The exact shape of this curve is not known, although it would be expected to lie much closer to the present numerical values. Insofar as this observation is due to inviscid, equilibrium effects, the numerical results are further corroborated. It would be of great interest to pursue the numerical studies in the neighborhood of the coincidence of the SMR-CMR and CMR-DMR lines.

It is also worth noting that in this set, the isopycnic shapes and distributions resemble those for RR until M $_{\rm S}$ 1.60, where a loop exists at the wedge corner and the next fringe away from this loop is bowed towards it. This effect becomes increasingly prominent as M $_{\rm S}$ increases through the CMR range, loops begin to form near the slipstream as DMR approaches, and prominently so as M $_{\rm S}$ increases through the DMR range. For smaller values of $\theta_{\rm W}$, such isopycnic distributions can occur for smaller values of M $_{\rm S}$ (see Figure 15, M $_{\rm S}$ = 2.03, $\theta_{\rm W}$ = 270). The foregoing gives some insight into the changing overall wave patterns as the (M $_{\rm S}$, $\theta_{\rm W}$) - plane is traversed.

Set 2: The CMR-DMR transition is studied for $M_S = 4.0$, Air; $29^{\circ} < \theta_w < 34^{\circ}$ perfect gas with $\gamma = 1.40$; $25^{\circ} < \theta_w < 30^{\circ}$, Hansen EOS; in increments of $\Delta\theta_w = 1^{\circ}$. The results are presented in Figures 21 and 22. The analytic CMR-DMR transition, Figure 2, for $M_S = 4.0$ takes place at $\theta_w = 32^{\circ}$ for a perfect gas and $\theta_w = 26^{\circ}$ for the Hansen EOS. The EOS effect is predicted correctly, that is, $\theta_w = 32^{\circ}$ for a perfect gas and 20° for the Hansen EOS so that this transition line is shifted up by about 3° . It is worth noting that the calculated Mach stems are not perpendicular to the wedge at the triple point; this is an assumption in the analytic calculations leading to Figure 2 (see Lee and Glass 1984). Also, it would not be unreasonable to require the calculated kinks to clearly sharpen up to a new triple point before assuming a DMR transition.

Set 3: The SMR-SMR and SMR-DMR transitions are studied for $M_c = \frac{1}{2}$ 3.75, Air: 5° < 3_{\circ} < 10° and 22° < 3_{\circ} < 26° , perfect gas with γ = 1.40; $53 < 3 \le 3$ and $150 < 3 \le 190$, Hansen EOS; in increments of $20 \le 10$. The results are presented in Figures 23 and 24. The analytic SMR-CMR transition Figure 1' for $M_{\pi} = 3.75$ takes place at $\theta_{\pi} = 30$ for a perfect gas and at agraduate for the Hansen EOS. According to our criteria involving the M sonic line, none of the reported calculations with $\epsilon_{\rm c} < 10^{
m o}$ are GMR with the possible exception of the $\theta_w = 90$ Hansen EOS result. Thus, the calculated transitions differ from the analytic results by at least 3°. Once again, none of the calculated Mach stems are perpendicular to the wedge at the triple point. The analytic CMR-DMR transition, Figure 2, for $M_c = 8.75$ takes place at θ_{ω} = 230 for a perect gas and at θ_{ω} = 160 for the Hansen EOS. The calculations show transition at $\theta_{\underline{u}}$ no greater than 22° for a perfect gas and at θ = 15-160 for the Hansen EOS. This represents close agreement. Here as well, the Mach stems are not perpendicular to the wedge at the triple point. It should be noted that the experimental results are not in close agreement with either of the two transition lines (see Figure 2) at such high values of Me. Consequently, new criteria may have to be found so that better agreement can be obtained for the SMR-CMR-DMR transition lines (see Deschambault and Glass 1983, and Hu and Glass 1985).

Set 4: $M_S = 7.10$, Argon (perfect gas with $\gamma = 5/3$); $490 < \theta_W < 550$, in increments of $\Delta\theta_W = 10$; $\theta_W = 52.750$, 53.750; $53.10 < \theta_W < 53.50$ in increments of $\Delta\theta_W = 0.10$. The purpose of this set of calculations is to estimate the inviscid RR-DMR transition boundary and, by comparison with experimental results, to demonstrate and quantify the well-known disagreement between theory and experiment for this issue (see, for example, Shirouzu and Glass 1982). An experimental interferogram for $\theta_W = 490$ and all of the computational results are presented in Figure 25. Noting the results in the range $53.00 < \theta_W < 53.50$ and comparing with the experiment, a value of $\Delta\theta_W = 4.0 - 4.50$ may be inferred as the "boundary-layer defect" (see Hornung and Taylor 1982; Shirouzu and Glass 1982; Wheeler and Glass 1985) for the $M_S = 7.10$ RR-DMR transition. We are referring, in particular, to the substantial disagreement concerning the extent of the Mach stem region relative to the entire flow field. We have attempted to calculate the precise RR-DMR transition point by plotting the height of the Mach stem relative to L

against $\theta_{\rm w}$ for the computations and extrapolating the curve to zero height. Figure 26. The result is $\theta_{\rm w}=53.35^{\circ}$ which disagrees moderately with the theoretical results of $\theta_{\rm w}=54.4^{\circ}$ in Lee and Glass, 1982. We remark that this error may be caused by an unnoticed bias in our measuring technique done by simply using a ruler on the computer-generated contour plots of the blow-up Mach stem region (not snown)), lack of numerical resolution when the Mach stem is only 1-2 zones high, or a numerical error in the post-shock flow field at the wall. In any case, the error is small relative to the viscous-inviscid difference and it is also possible that the theoretical inviscid prediction of $\theta_{\rm w}=54.4^{\circ}$ does not apply when the entire disturbed flow field is taken into account. Higher resolution calculations using an adaptive mesh scheme, Berger et al 1985, will be carried out in an effort to settle this issue. This set of calculations also illustrates the dramatic collapse of the complex DMR-pattern into the simple RR-pattern as $\theta_{\rm w}$ changes by a fraction of a degree (see Figures 251 and m).

SECTION 5

CONCLUSIONS

A computer tode has been developed for the inviscid, perfect gas shock-on-wedge problem and the results have been compared with the best available experimental data. The code is based on contemporary methodology in the numerical analysis of hyperbolic conservation laws, and has only recently been available.

Good to excellent qualitative agreement has been obtained in all cases of direct comparison, and this applies to structures beneath the reflected shock such as the vortex roll-up as well as coarser criteria such as the reflection pattern. Quantitatively, the results are very good for flow fields without observable nonequilibrium or viscous effects, except for Case 2. The error in this case is probably a result of the relatively large variation of the solution with respect to small increments in the problem parameters in the vicinity of the parameter values defining this case. When nonequilibrium or viscous effects are present, the quantitative error can be 10-15% and we may recall Case 11 which has a much larger, and unexplained error.

Although not entirely proven, it appears that the computer code represents a substantial predictive capability for the shock-on-wedge problem restricted to inviscid, perfect gases. Even for viscous, real gas flow fields, the computational results provide a significant amount of information, including highly resolved flow-field structures.

Significant non-equilibrium and viscous effects have been demonstrated in the shock wave diffraction experiments. Much of this could be inferred without the numerical study, but the latter can provide a quantitative estimate of the various effects. In particular, vibrational relaxation is observed in the high shock wave Mach number cases, and this can have large-scale effects on criteria such as the corner attachment angle and type 'subsonic or supersonic) and viscous effects are important in determining the vortex roll-up pattern and the wedge corner flow field. Although these effects occur in thin layers or small regions, they may have an effect on the quantitative results in the inviscid portion of the flow field.

The capability of the computer code to discriminate between very small increments in problem parameters (M_s , θ_w , and the equation of state, although the latter has not been treated here) has been demonstrated.

By using parametrized sequences of calculations, it would be possible to construct transition boundaries in the (M_s, θ_w) -plane. Of tourse, the transitions obtained would be dependent on the transition criteria used in their construction; our use of the sonic criterion in self-similar coordinates shows now the infinite amount of data potentially available from a calculation can be invaluable in evaluating one of the proposed criteria.

The discussion of transition set 1 in Section 5 illustrates how parametrized numerical calculations can be used to elucidate details of the flow field transition not otherwise available. It is quite possible that such results will prove useful in the discovery of more precise analytic transition criteria, in the future. For the high $\rm M_S$ transitions, the inviscid numerical results provide a guide for the analysis of inviscid transition criteria in a parameter regime where analytic-experimental agreement has been relatively poor and where nonequilibrium phenomena are hard to avoid in the experiments. Of course, the formulation of transition criteria for viscous, nonequilibrium flow fields is not assisted by the present computer code.

Also we have been able to validate the conjecture that the RR-DMR. transition is offset in experiments by a boundary-layer defect.

In Section 5, several calculations were noted where our analysis could be greatly improved with a more efficient adaptive mesh in the vicinity of the Mach stem. Obtaining the necessary resolution with the present computer code would be overly expensive if carried out on a production basis for a large number of calculations. Using the methods of Berger and Colella 1985, Berger et al 1985, we expect to overcome this problem and revisit some of the cases discussed in this report. Additionally, we are working on techniques to reduce further or eliminate the starting error and boundary error from our results.

In future work, we intend to modify our computer code and include an approximation for vibrational relaxation. We expect that this work will settle some of the questions raised in this paper. The results presented here demonstrate, however, that a valid approximate solution method for the Navier-Stokes equations will be required if complete agreement between experiment and calculation is demanded. Despite these shortcomings, the comparison of the present numerical simulations with interferometric data from RR, SMR, CMR and DMR experiments are probably the best available to date.

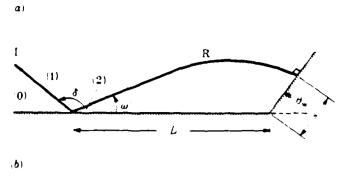
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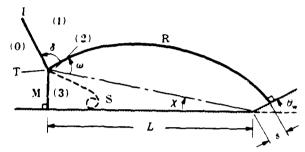
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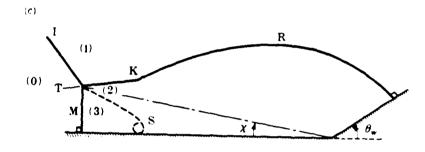
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Table 1. Initial conditions, equation of state, and computational mesh for each experimental case.

CASE	FIGURE	GAS	TYPE	Эw	Ms	Po (k Pa)	ρο (gm/cc)	EOS	NX	NY
1	4	Argon	RR	60°	2.05	20.00	3.23 X 10 ⁻⁴	y = 5/3	355	90
2	5	Air	RR	45 ⁰	1.26	101.12	1.146 X 10 ⁻³	γ= 1.4 Hansen	350	160
3	6	Air	SMR	45°	1.50	50.66	5.73 X 10 ⁻⁴	γ= 1.4 Hansen	375	160
4	7	Air	DMR	47 ⁰	3.03	3.33	3.77 X 10 ⁻⁵	γ= 1.4 Hansen	500	120
5	8	Air	CMR	30°	2.65	13.33	1.52 X 10 ⁻⁴	γ= 1.4 Hansen	390	125
6	9	Argon	CMR	30o	5.07	4.00	6.45 X 10 ⁻⁵	y = 5/3	420	140
7	10	Air	CMR	10°	10.37	6.67	7.53 X 10 ⁻⁵	Hansen	400	140
8	11	Air	SMR	40°	1.66	33.33	3.8 X 10 ⁻⁴	γ= 1.4 Hansen	375	135
9	12	Air	DMR	40°	2.87	16.67	1.9 X 10 ⁻⁴	γ= 1.4 Hansen	420	110
10	13	Air	DMR	40°	3.72	6.00	6.87 X 10 ⁻⁵	γ= 1.4 Hansen	420	100
11	14	Air	DMR	40°	4.62	2.80	3.19 X 10 ⁻⁵	γ= 1.4 Hansen	420	90
12	15	Air	SMR	27 ⁰	2.03	33.33	3.87 X 10 ⁻⁴	y= 1.4	350	130
13	16	Air	DMR	27°	8.70	4.10	4.76 X 10 ⁻⁵	Hansen	440	85
14	17	Air	C/DMR	20°	7.19	8.00	9.29 X 10 ⁻⁵	Hansen	420	120
15	18	Air	DMR	20°	8.86	4.10	4.65 X 10 ⁴	Hansen	500	110







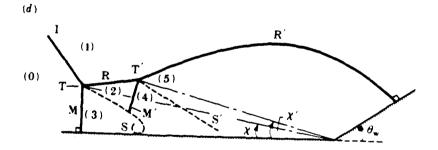


Figure 1. Schematic diagrams of types of oblique shock-wave reflections: (a) RR; (b) SMR; (c) CMR; (d) DMR; also definitions of L and s.

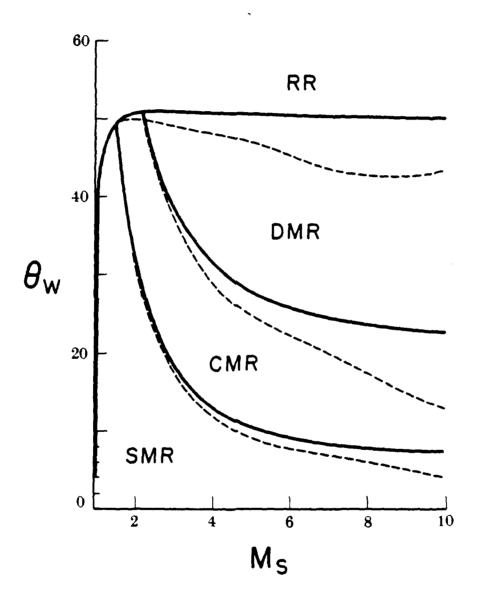


Figure 2. Regions of RR, SMR, CMR, and DMR and their transition boundaries in the (M_S, θ)-plane for perfect (frozen) air solid lines and imperfect (equilibrium) air broken lines, $p_0 = 2.00$ kPa, $T_0 = 300$ K, $\gamma = 1.40$.

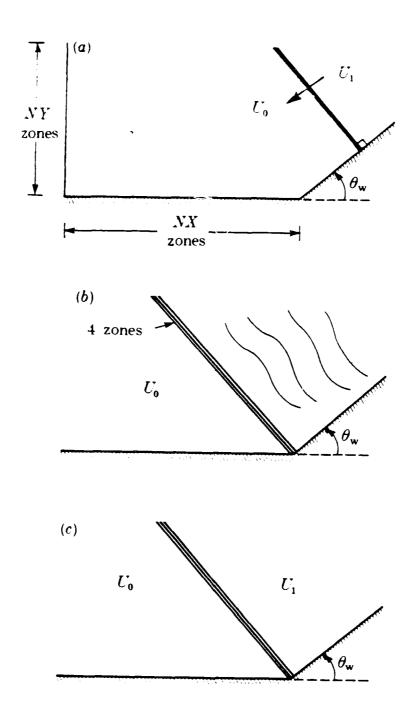
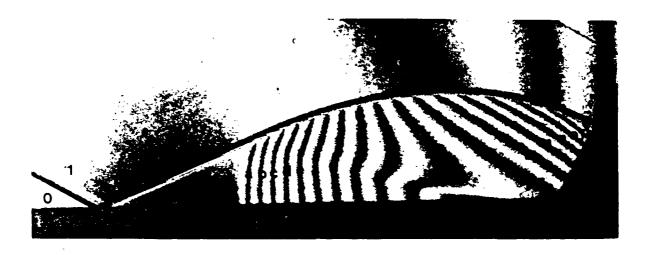


Figure 3. Numerical scheme for flow initialization; (a) starting procedure; (b) shock reaching corner; (c) elimination of small disturbances.

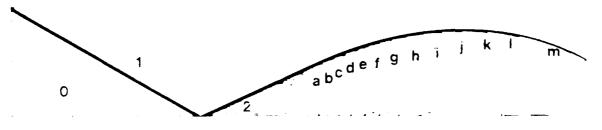


Region	\mathcal{A}_{v}	Region	:/::
j.	1.00	h	3.86
!	2.33	i	3.80
2	4.38	j	3.73
a	4.32	k	3.67
<u>.</u>	4.25	1	3.60
d.	4.19	m	3.54
ä	4.12	n	3.47
ę.	4.06	0	3.41
÷	3.99	p	3.34
	3.43	q	3.28

Figure 4a. Interferogram

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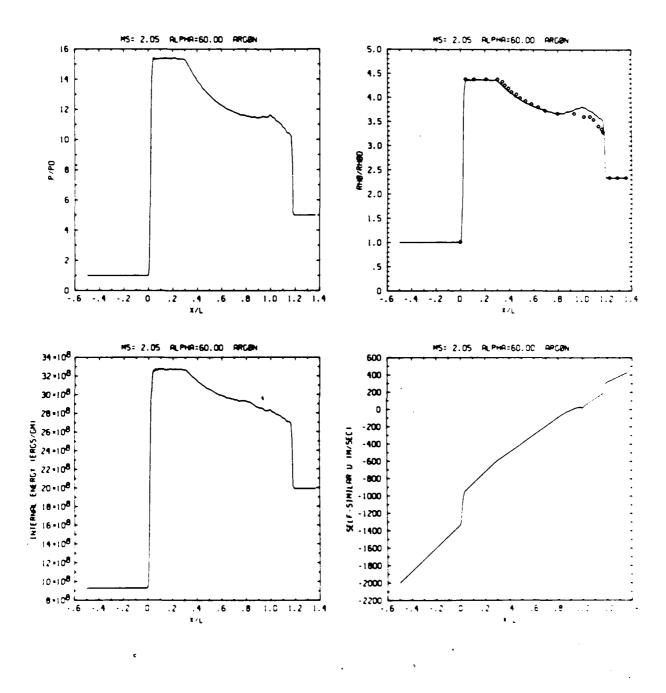


Figure 4c. Wall plots for $p/p_0,\; \rho/\rho_0$ with experimental data included, e, \bar{u}

Figure 4. Case 1, M_s = 2.05, θ_w = 60°, Argon, γ = 5/3, RR - Continued.

MS= 2.35 ALP=60.30 NR=400 NZ= 30 KBEG= 45 PD=2.305+35 AFS21.

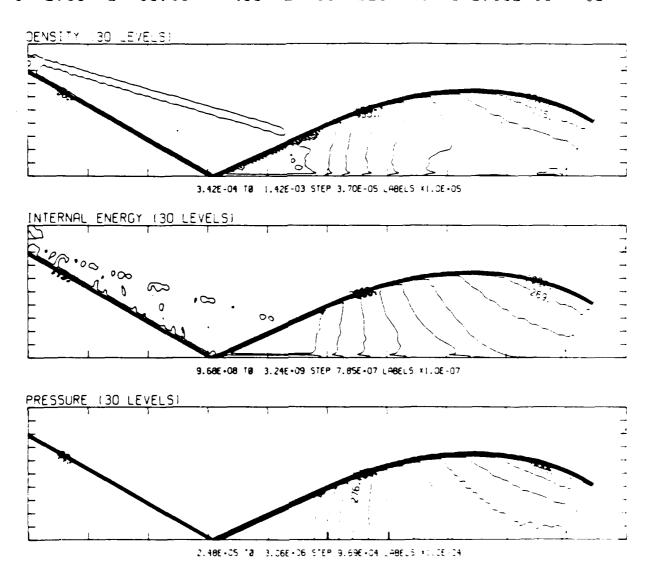


Figure 4d. Whole-flowfield contour-plots

Figure 4. Case 1, $M_S = 2.05$, $\theta_W = 60^\circ$, Argon, $\gamma = 5/3$, RR - Continued.

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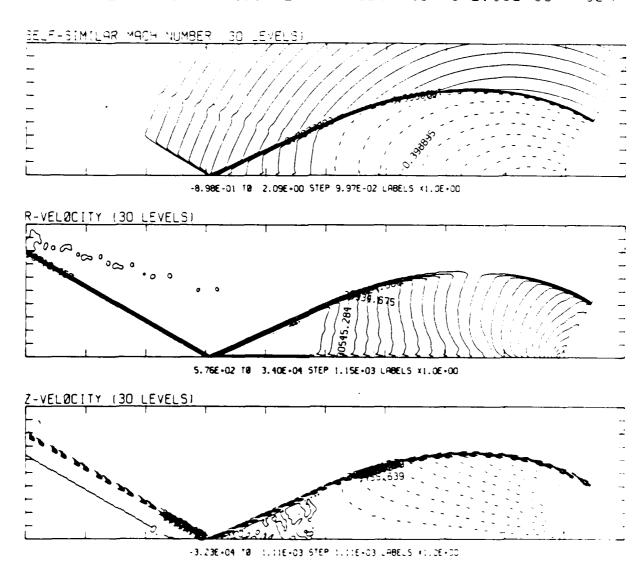
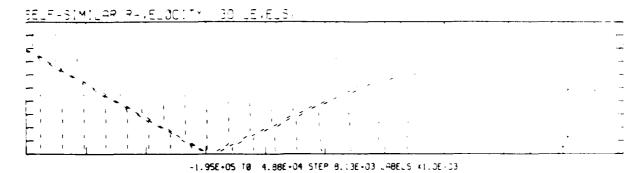
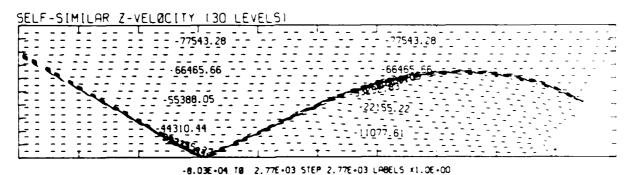


Figure 4d. Whole-flowfield contour-plots - continued.

Figure 4. Case 1, $M_s = 2.05$, $\theta_w = 60^{\circ}$, Argon, $\gamma = 5/3$, RR - Continued.

MS= 2.05 ALP=60.00 NF=400 NZ= 30 KBE3= 45 PS=2.00E+35 440.0%





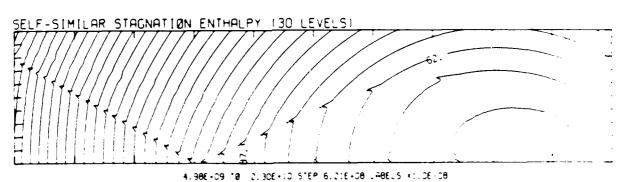
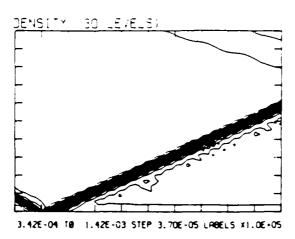
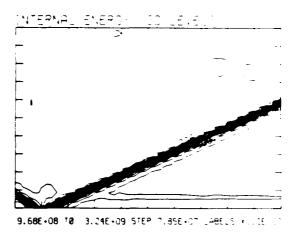


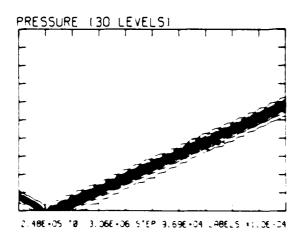
Figure 4d. Whole-flowfield contour-plots - continued.

Figure 4. Case 1, M_S = 2.05, θ_W = 60°, Argon, γ = 5/3, RR - Continued.

MBF 2.05 ALP460.00 (L#282 184288 UT+ 36 F0#2.00E408 44)25







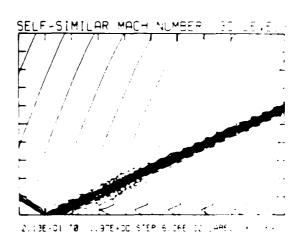
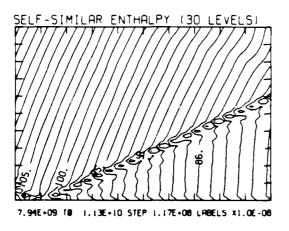
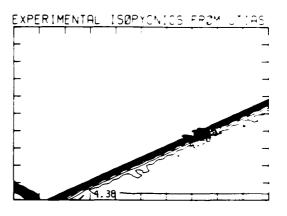


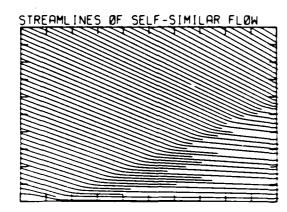
Figure 4e - Blowup-frame plots

Figure 4. Case 1, M_S = 2.05, θ_W = 60°, Argon, γ = 5/3, RR - Continued.

MS= 2.05 ALP=60.00 IL=232 IR=283 JT= 36 P0=2.00E+05 AP021.







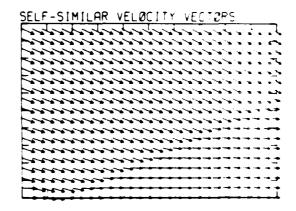
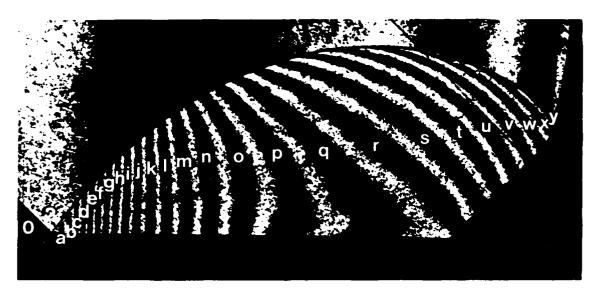


Figure 4e. Blowup-frame plots - continued.

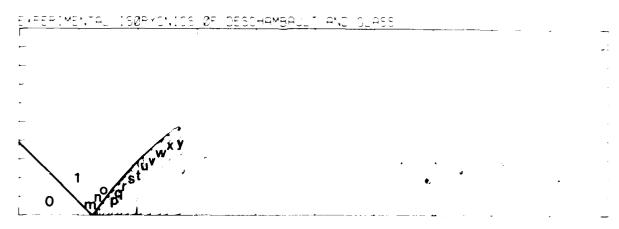
Figure 4. Case 1, M_S = 2.05, θ_W = 60°, Argon, γ = 5/3, RR - Continued.



Region	s/s _s	Region	٥/٥,	Region	e/e,
0	1.00	h	2.29	r	2.03
1	1.45	i	2.26	s	2.00
2	2.49	j	2.23	t	1.97
a	2.47	k	2.21	u	1.95
Ь	2.44	1	2.18	v	1.92
С	2.42	m	2.16	$\boldsymbol{\omega}$	1.89
d	2.39	n	2.13	х	1.87
e	2.36	0	2.10	у	1.84
f	2.34	P	2.08		
g	2.31	q	2.05		

Figure 5a. Interferogram

MS= 1.26 ALF=45.00 NR=500 NZ=160 MBEG=150 PC=1.11: 4 / / / / /



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Figure $5b_p$. Calculated isopycnics ($\gamma=1.4$) using the experimental fringes

Figure 5. Case 2, M_s = 1.26, θ_w = 450, Air, γ = 1.4 and Hansen EOS, RP.

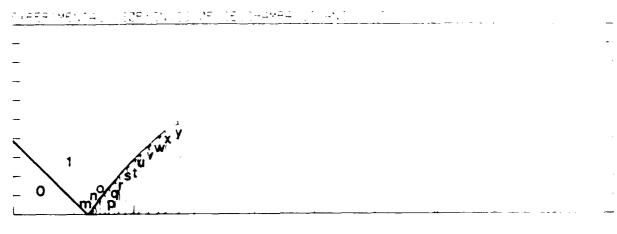


Figure 5bH. Calculated isopycnics (Hansen) using the experimental fringes $45 \pm 1.26 + 42 \pm 43.22$

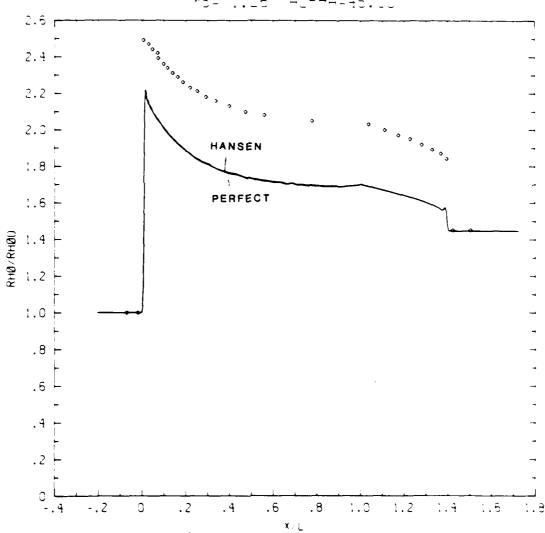


Figure 5c. Wall plot for ρ/ρ_0 , $\gamma=1.4$ and Hansen calculations, with experimental data

Figure 5. Case 2, M_s = 1.26, 9 = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

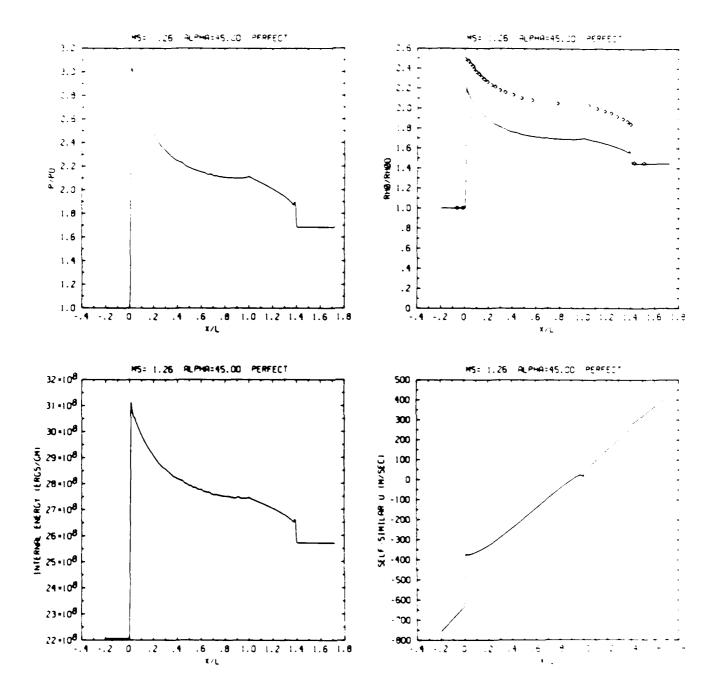


Figure $5c_p$. Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e, u; γ = 1.4.

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

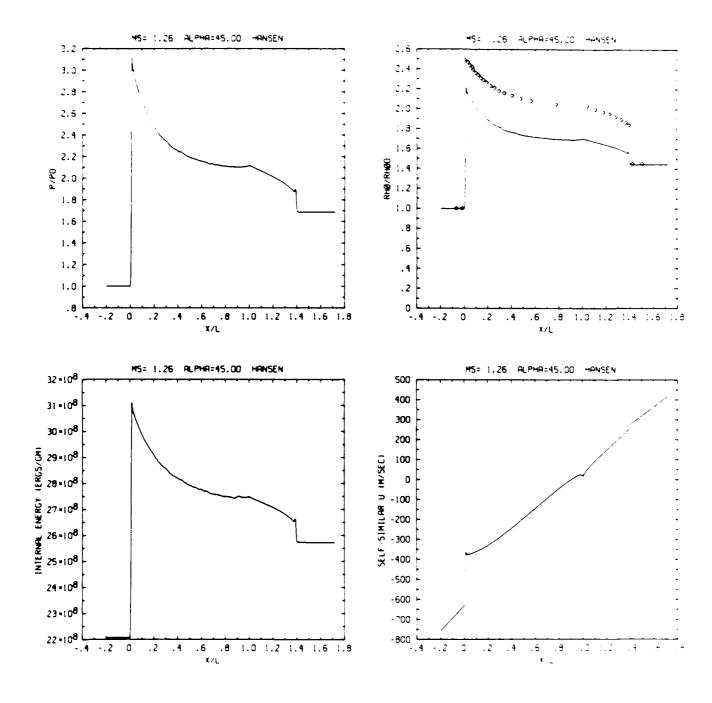
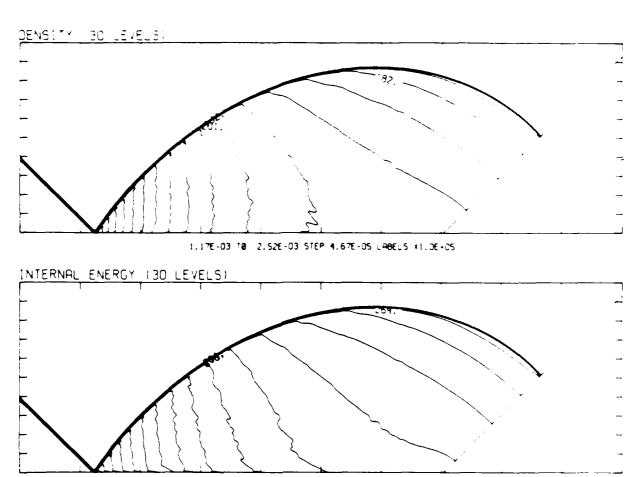
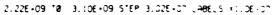


Figure 5c_H. Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen.

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.





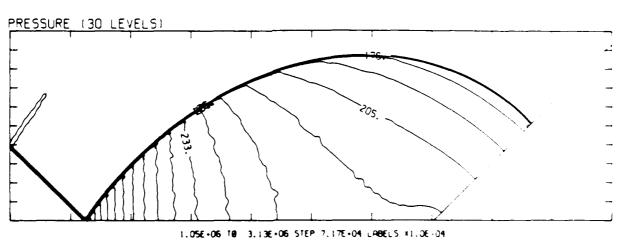
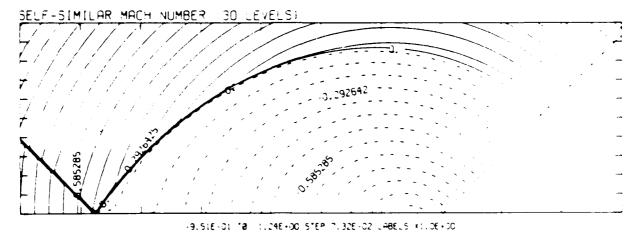
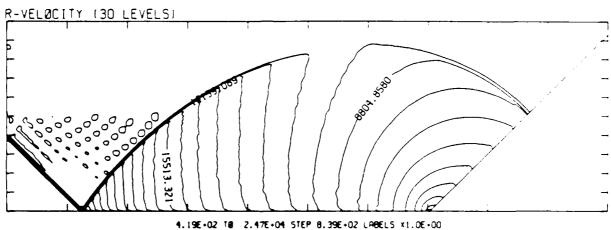


Figure $5d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$.

Figure 5. Case 2, M_S = 1.26, 9 = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.





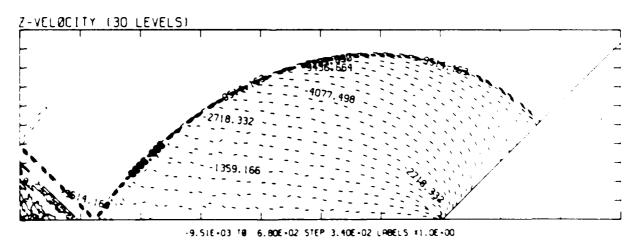
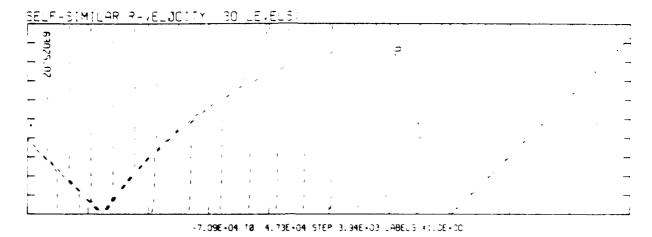
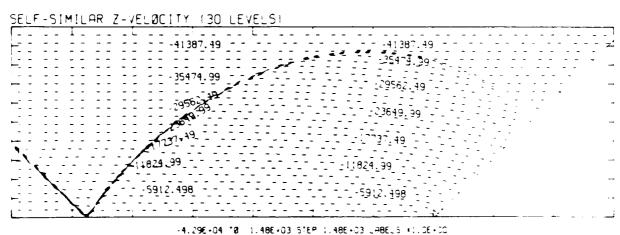


Figure $5d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.





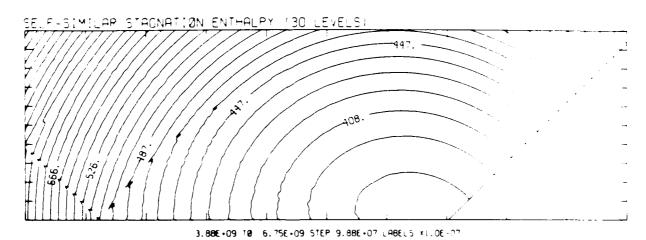


Figure $5d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

MS= 1.26 ALP=45.00 [L=395 [R=442 UT= 44 PO=1.01E+36 PERFEST

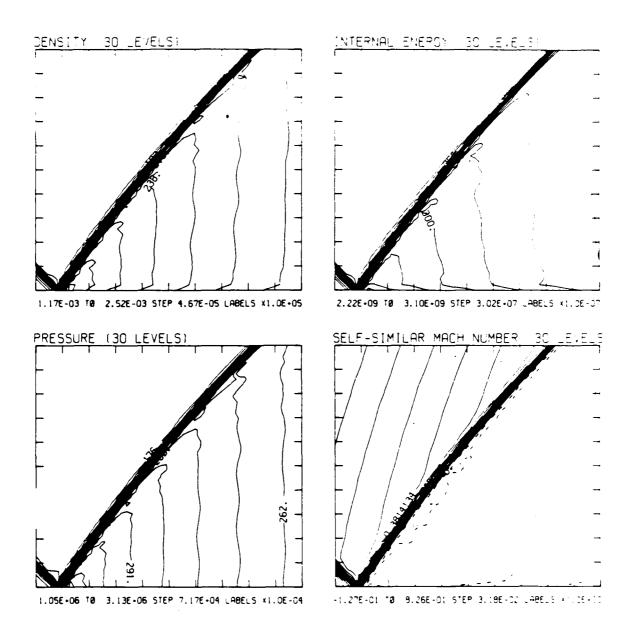


Figure $5e_p$. Blowup-frame plots; $\gamma = 1.4$

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

MS= 1.26 ALP=45.00 (L=395 (P=442 UT= 44 PD=1.015+18 -8885)

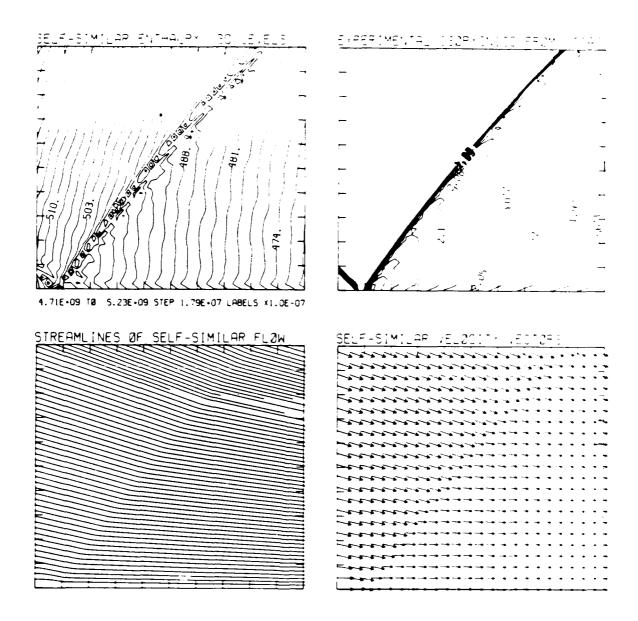
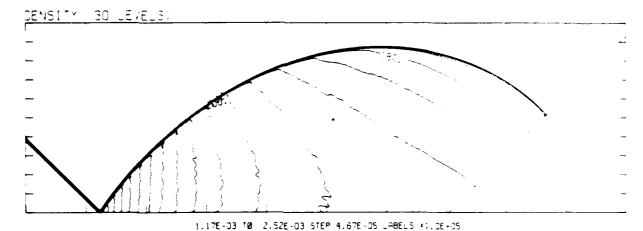


Figure $5e_p$. Blowup-frame plots; $\gamma = 1.4$ - continued

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

MS= 1.26 ALP=45.00 NR=500 NZ=160 MBEG=150 AD=1.115+ W HAN N



INTERNAL ENERGY (30 LEVELS)

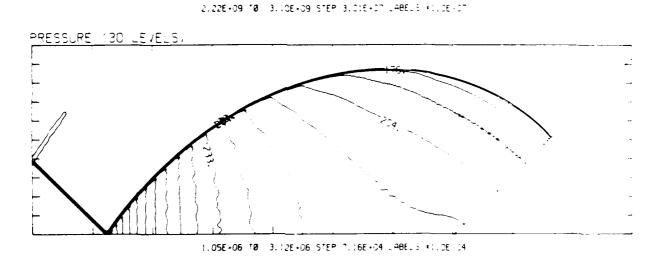
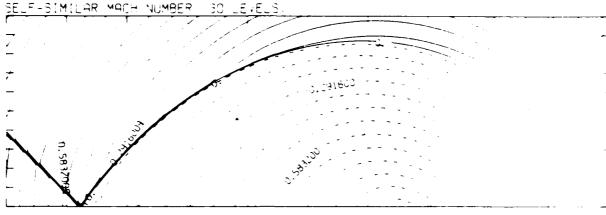


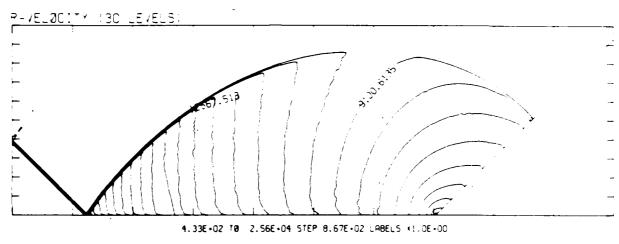
Figure 5d_H. Whole-flowfield contour-plots; Hansen

Figure 5. Case 2, M_S = 1.26, 9 = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

MS= 1.26 ALP=45.00 NR=500 NZ=160 MBE0=150 F171.115474 HAN N



-9.48E-01 18 0.24E+00 STEP 7.29E-02 148EUS *0.0E+00



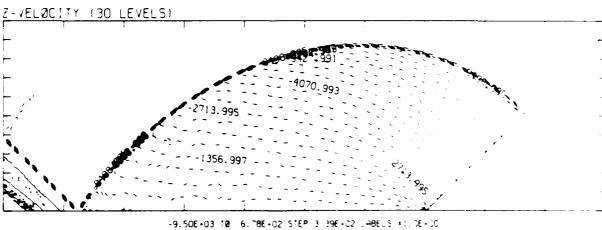


Figure $5d_{\mathrm{H}}$. Whole-flowfield contour-plots; Hansen - continued

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

48= 1.28 ALP=45.00 NF=500 NZ=080 HBE0=150 F0=1.7174 A HAN H

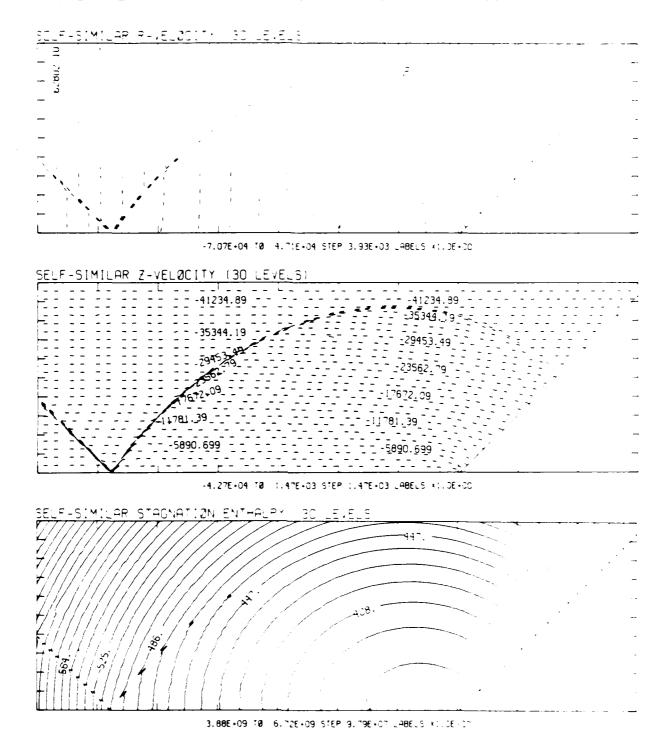


Figure $5d_{\mbox{\scriptsize He}}$. Whole-flowfield contour-plots; Hansen - continued

Figure 5. Case 2, M_S = 1.26, 9 = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

MRE 1.08 ALPE45.00 (Le398 (RE443 LTE 44 ADE).0084(8 HAN)

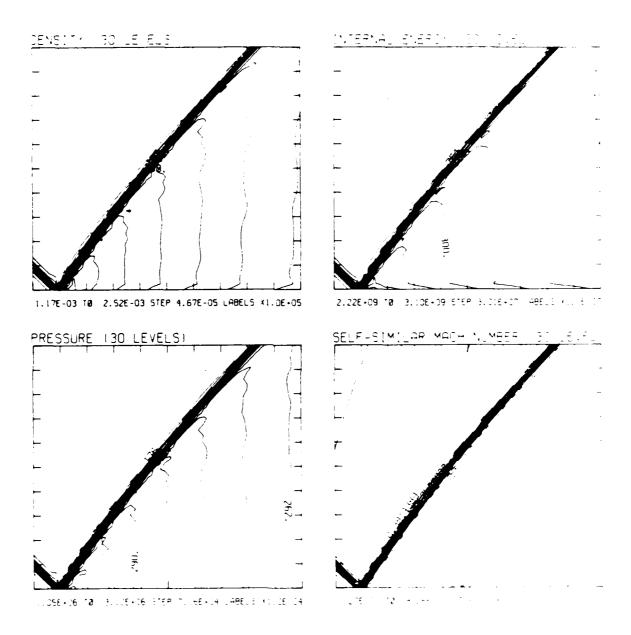


Figure 5e_H. Blowup-frame plots; Hansen

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.

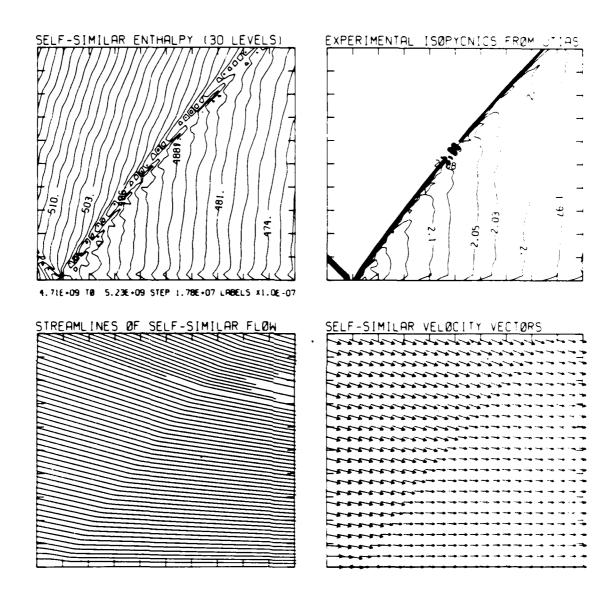


Figure 5e_H. Blowup-frame plots; Hansen - continued

Figure 5. Case 2, M_S = 1.26, θ = 45°, Air, γ = 1.4 and Hansen EOS, RR - continued.



Region	٥/٥٥	Region	٥/٥ _°
0	1.00	i	2.75
1	1.86	j	2.69
2	3.22	k	2.64
3	2.93	1	2.59
a	3.17	m	2.54
ь	3.11	n	2.48
С	3.06	0	2.43
d	3.01	p	2.38
e	2.96	q	2.33
f	2.90	r	2.28
g	2.85	S	2.22
h	2.80	t	2.17

Figure 6a. Interferogram

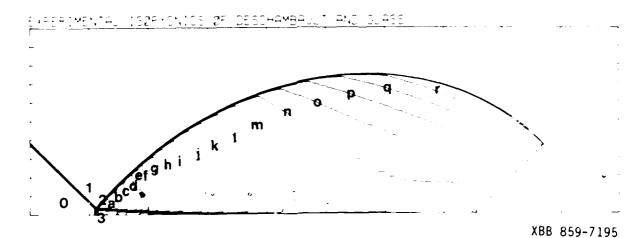


Figure $6 b_{\rm p}$. Calculated isopycnics (γ =1.4) using the experimental fringes.

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued 58

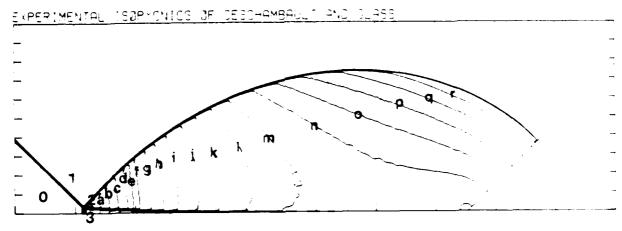


Figure $6b_{\mbox{\scriptsize H}}$. Calculated isopycnics (Hansen) using the experimental fringes

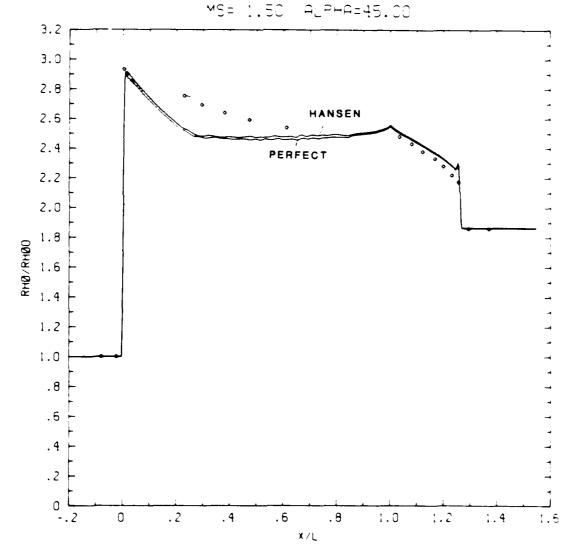


Figure 6c. Wall plot for ρ/ρ_0 , $\gamma=1.4$ and Hansen calculations, with experimental data.

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

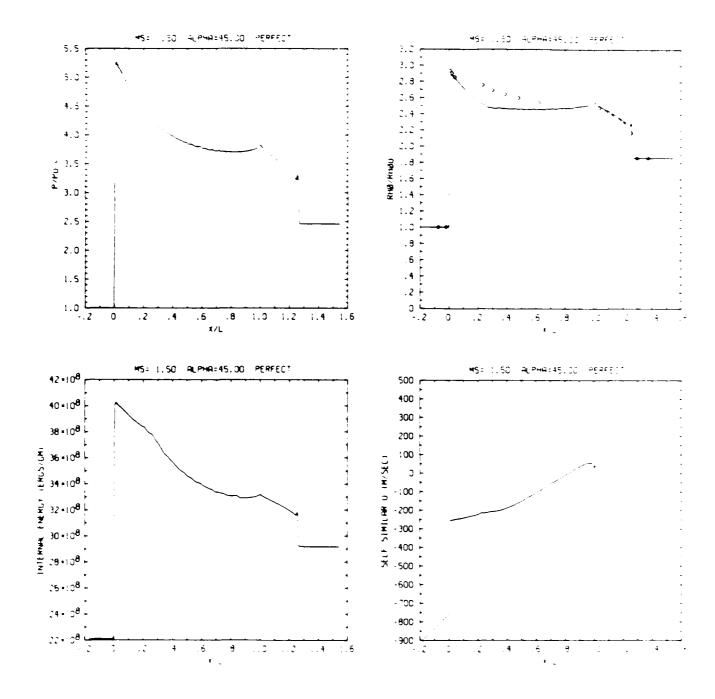


Figure 6cp. Wall plot for p/p_0 , p/p_0 with experimental data included, e, u; γ = 1.4.

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

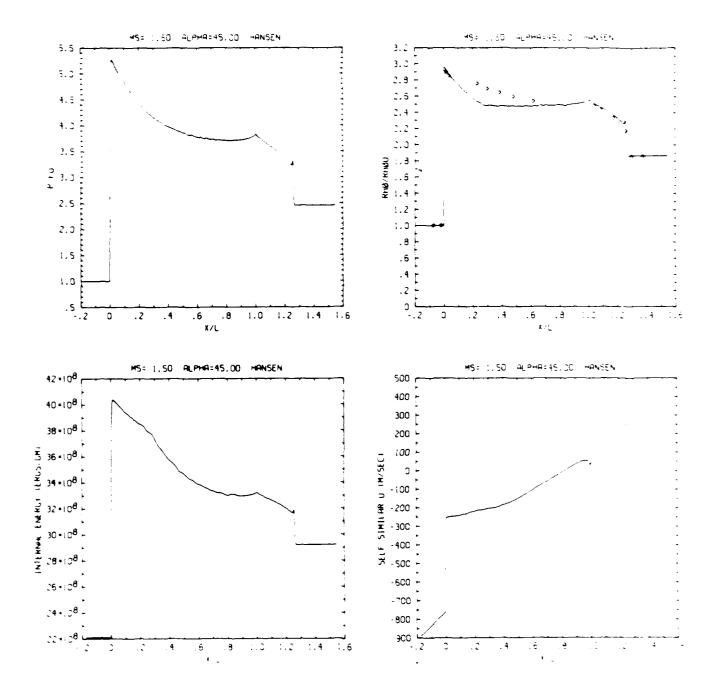


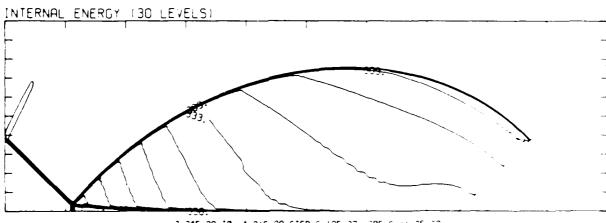
Figure 6c $_{\rm H^{*}}$ Wall plot for p/p $_{\rm 0}$, $_{\rm 0}/_{\rm 0}$ with experimental data included, e, u; Hansen

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

MS= 1.50 ALP=45.00 NR=500 NZ=180 48E0=125 F0=5.07E+15 AEAFE17



5.94E-04 18 1.79E-03 STEP 4..3E-05 LABELS (1.0E-05



2.24E-09 '0 4.01E-09 STEP 6.10E-07 LABELS 41.0E-07

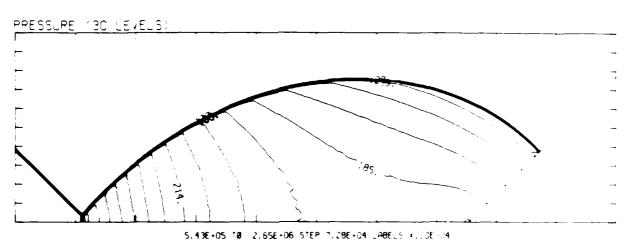


Figure 6dp. Whole-flowfield contour-plots; y = 1.4

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

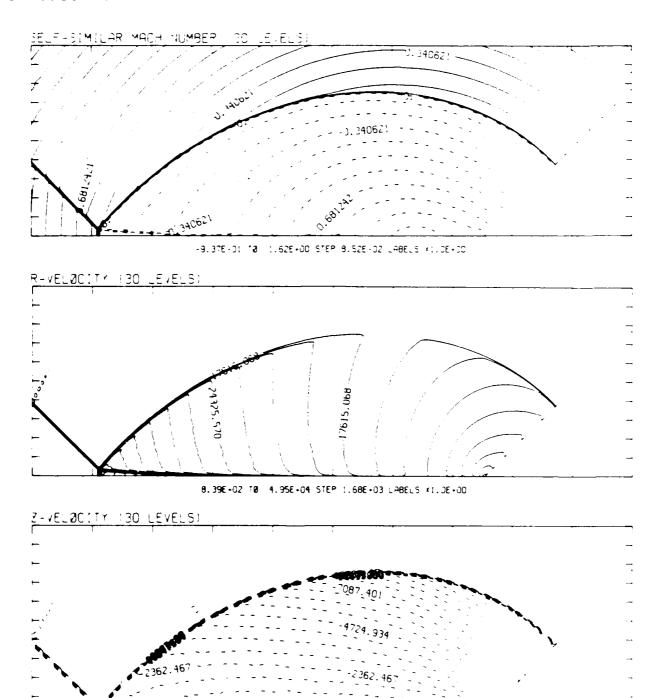
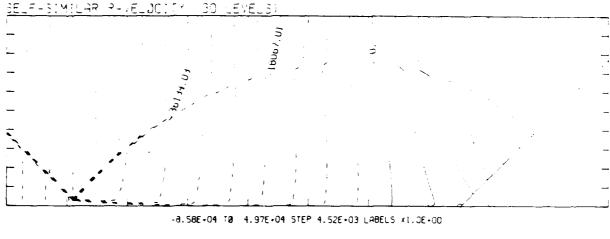
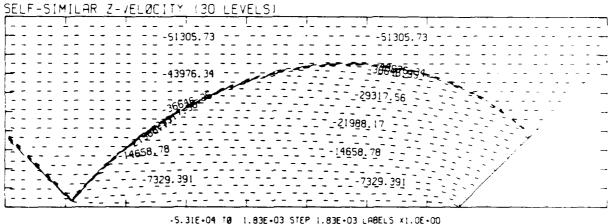


Figure $6d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued

-1.71E+04 70 5.91E+02 STEP 5.91E+02 LABELS K1.0E+00

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EDS, SMR + continued





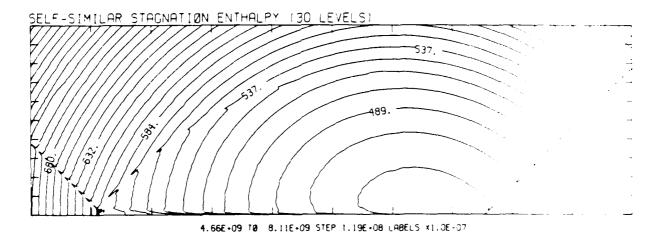


Figure 6dp. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

MS= 1.50 ALP=45.00 (L=395 (R=447 UT= 49 PD=5.07E+05 PERFE)

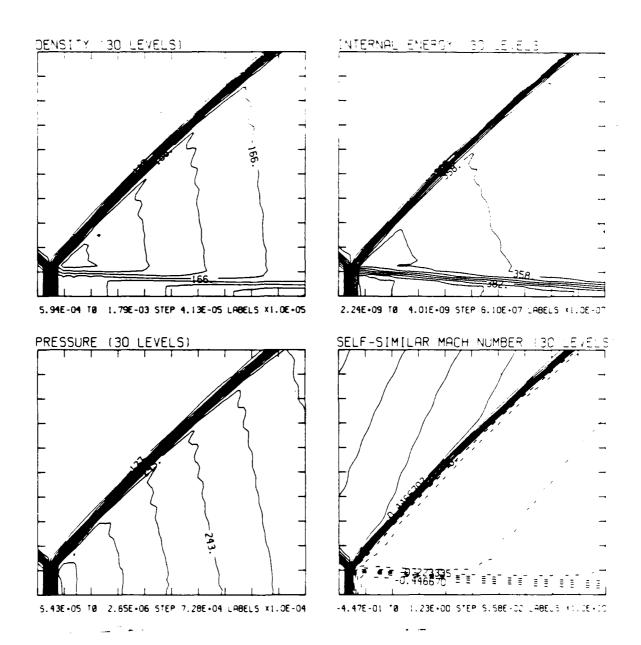


Figure $6e_p$. Blowup-frame plots; $\gamma = 1.4$

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

MS= 1.50 ALP=45.00 (L=395 (R=447 UT= 49 PC=5.0TE+65 PERFECT

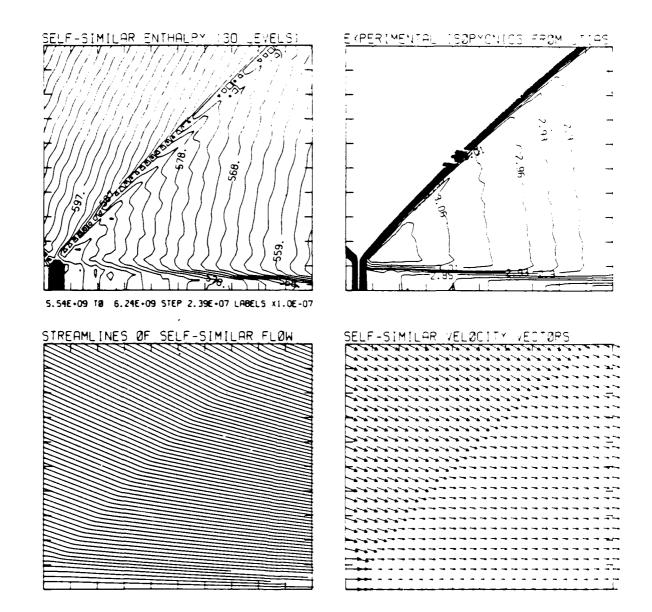
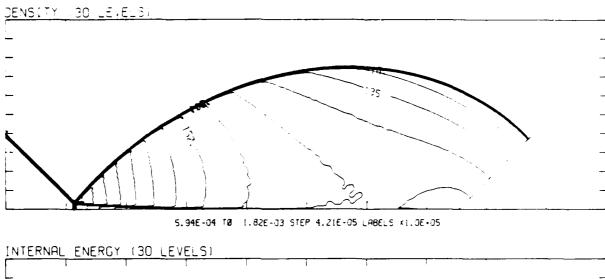
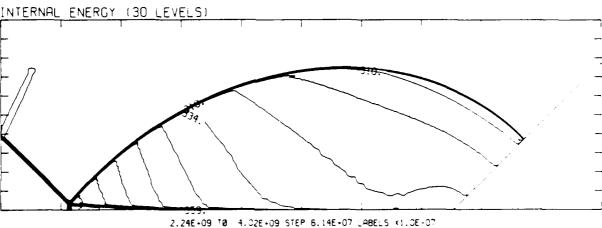


Figure 6ep. Blowup-frame plots; $\gamma = 1.4$ - continued

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

MS= 1.50 ALP=45.00 MR=500 NZ=160 M8E3=105 PD=5.0TE+15 HANGEN





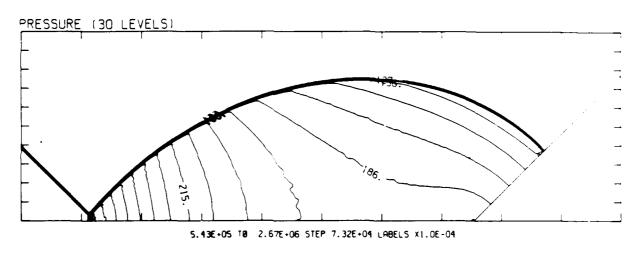


Figure $6d_{\mbox{\scriptsize H}}$ - Whole-flowfield contour-plots; Hansen

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued

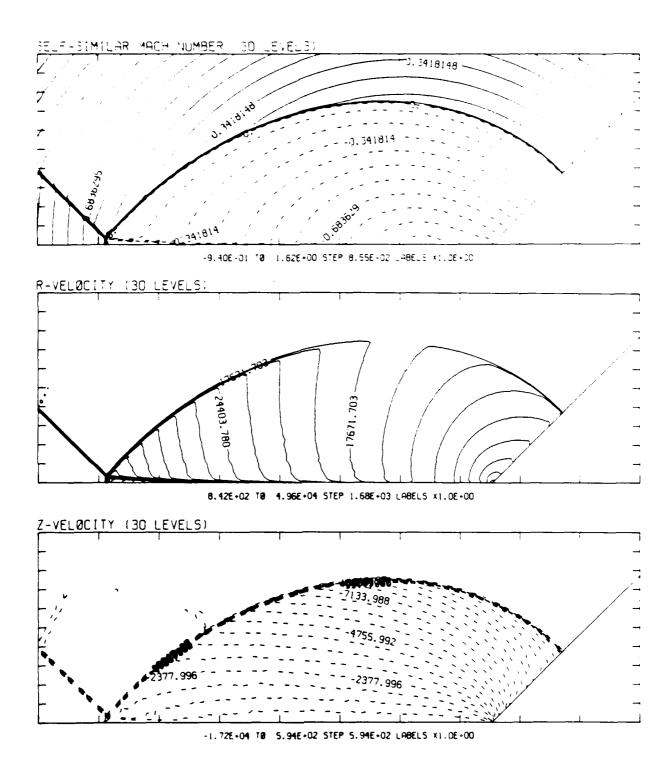
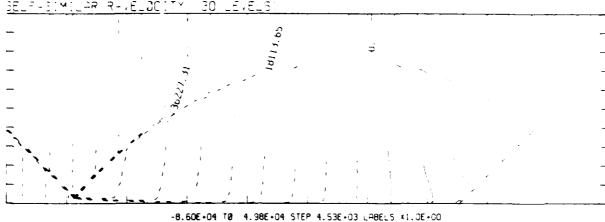
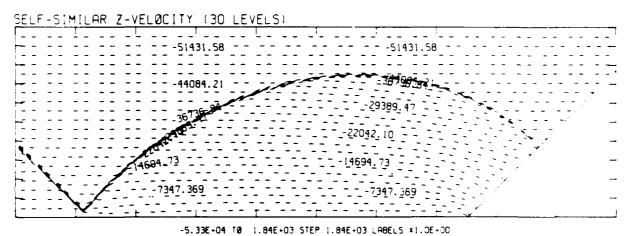


Figure 6d_H. Whole-flowfield contour-plots; Hansen - continued. Figure 6. Case 3, M_S = 1.50, $\theta_{\rm W}$ = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.50 ALP=45.00 NR=500 NZ=180 MBED=105 PD=5.078+05 H4NJEN





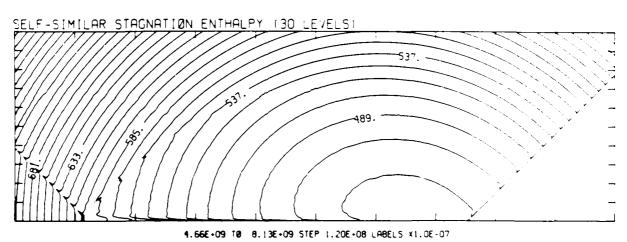


Figure 6d_H. Whole-flowfield contour-plots; Hansen - continued.

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.50 ALP=45.00 [L=394 [R=446 UT= 49 PO=5.0TE+05 HANSEN

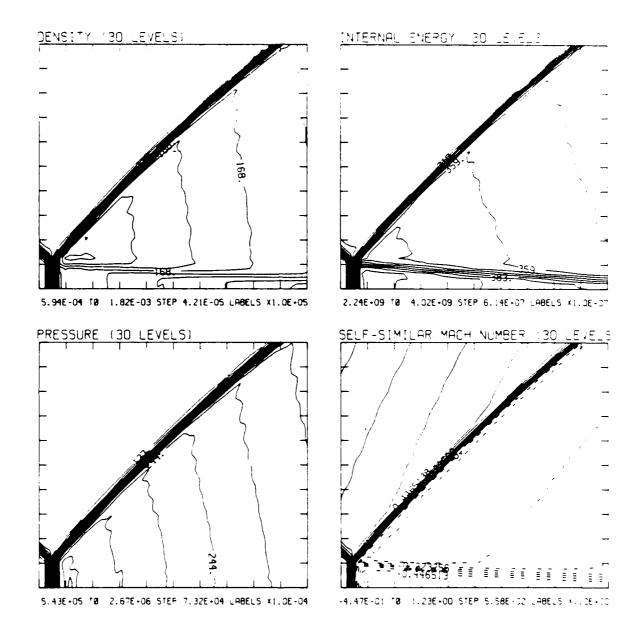


Figure 6e_H. Blowup-frame plots; Hansen.

Figure 6. Case 3, M_S = 1.50, θ_W = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

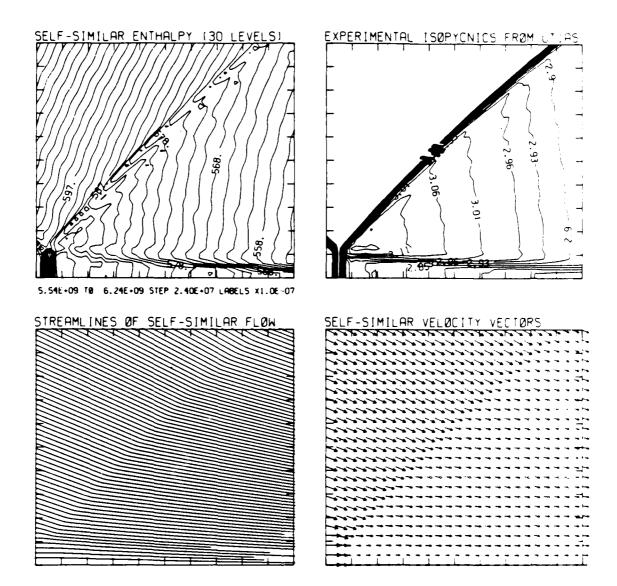
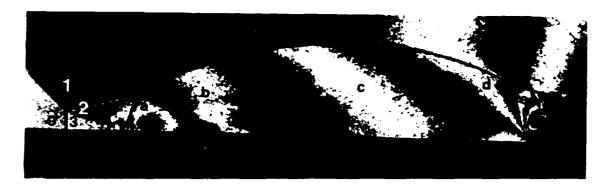


Figure 6e_H. Blowup-frame plots; Hansen - continued.

Figure 6. Case 3, M_S = 1.50, $\theta_{\rm W}$ = 45°, Air, γ = 1.4 and Hansen EOS, SMR - continued.



Region	c/c,	Region	۵/۵٫
0 1 2 3 a b	1.00 3.88 7.21 4.90 9.60 8.80	c d e f g	8.00 7.21 6.41 5.61 4.82

Figure 7a. Interferogram

Mus 3.03 ALP=47.00 NR=575 NZ=100 KEED= TE Hove.exe 4 4 4 4 4

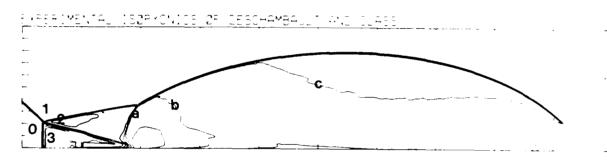


Figure $7b_p$. Calculated isopycnics (γ = 1.4) using the experimental fringes $9.03 \times 10^{-2} = 47.00 \times 10^{-2} = 120 \times 10^$

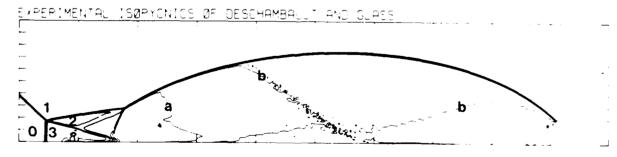


Figure 7b_H. Calculated isopycnics (Hansen) using the experimental fringes

XBB 859-7196

Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR.

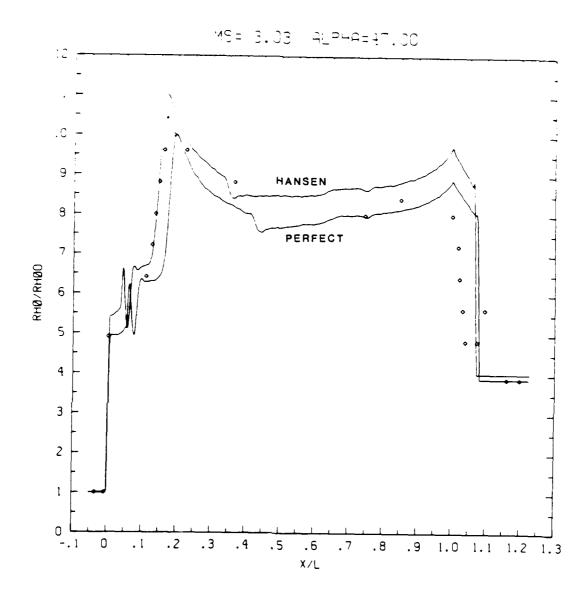


Figure 7c. Wall plot for p/p_0 , $\gamma=1.4$ and Hansen calculations, with experimental data.

Figure 7. Case 4, M = 3.03, $\theta_{\rm W}$ = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

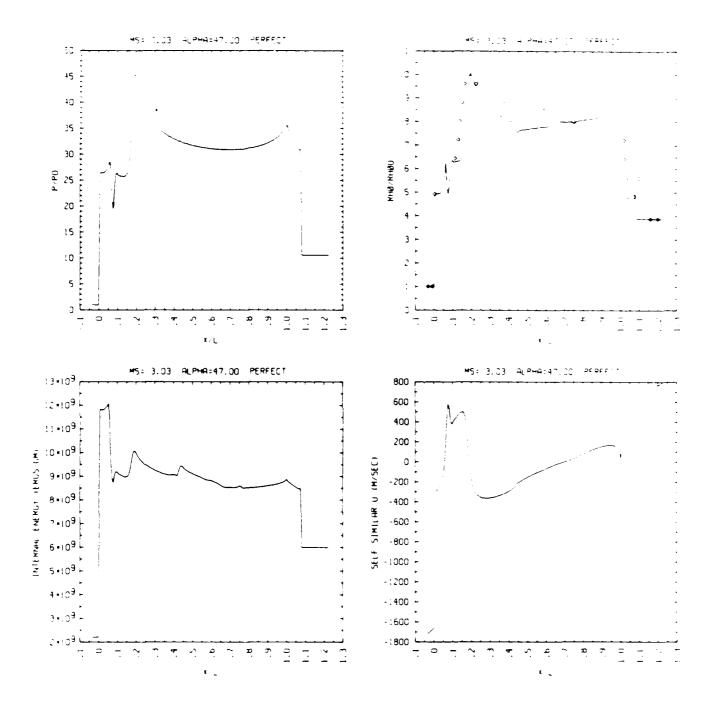


Figure 7c_p. Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e, u; γ = 1.4

Figure 7. Case 4, M_S = 3.03, $\theta_{\rm W}$ = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

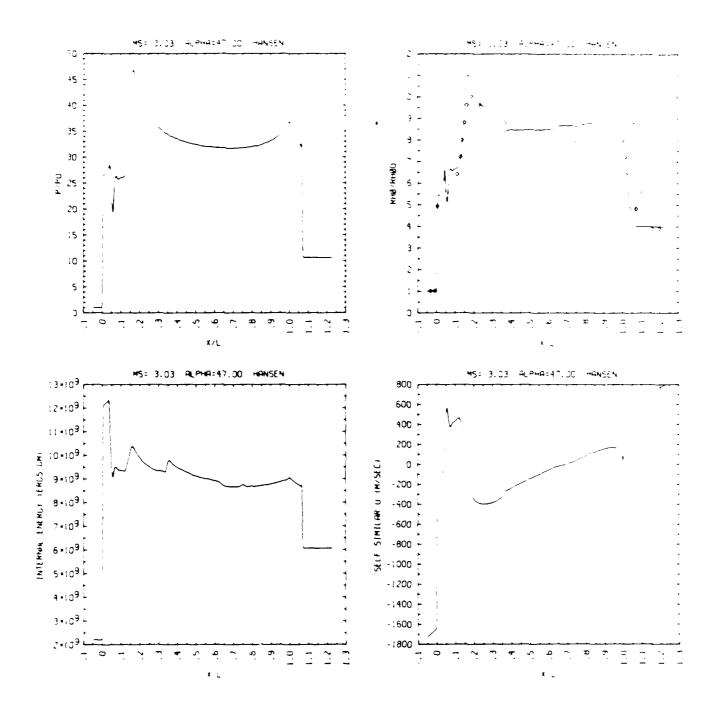
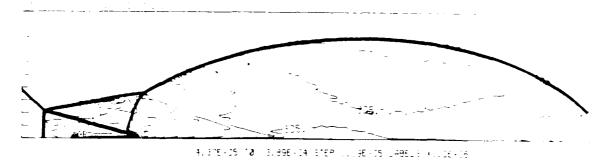


Figure 7c_H. Wall plots for $p/p_0,\; \rho/\rho_0$ with experimental data included, e, u; Hansen.

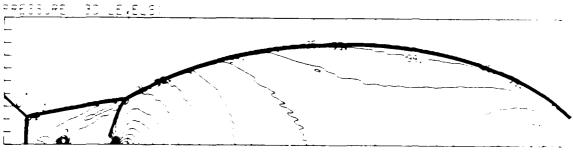
Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.



Y,≛84Y,4 FY,840+ 35 F,8 S



2.38E+09 10 1.23E+10 STEP 3.44E+08 LABELS #1.0E+08

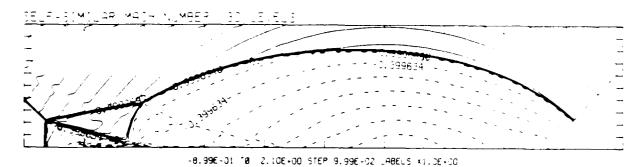


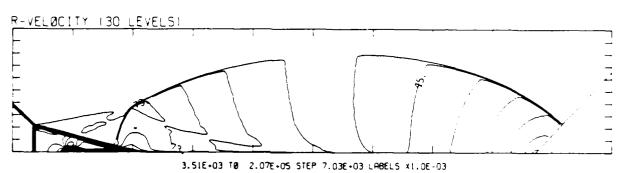
5.808+04 18 0.498+06 518P 4.938+04 048808 # 008 04

Figure $7d_p$. Whole-flowfield contour-plots; y = 1.4

Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

M94 3.03 ALPHAN.00 MRHANE M24.00 HEEGA NE POAS.335404 AERAEDT





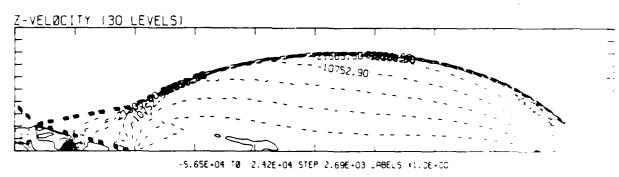
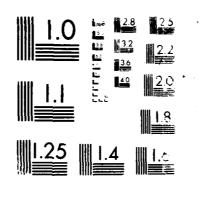


Figure 7d_p. Whole-flowfield contour-plots; v = 1.4 -

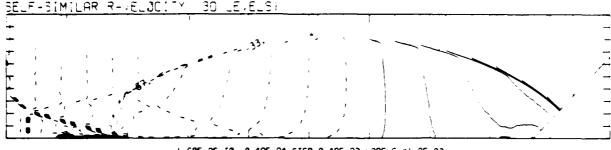
Figure 7. Case 4, $M_s = 3.03$, $\theta_w = 47^\circ$, Air, $\theta_w = 1.10^\circ$ continued.

A DETAILED NUMERICAL GRAPHICAL AND EXPERIMENTAL STUDY
OF OBLIQUE SHOCK MA (U) TORONTO UNIV DOMNSVIEW
(ONTRRIO) INST FOR AEROSPACE STUDIES H M GLAZ ET AL
01 AUG 86 UTIAS-285 DNA-TR-86-365 F/G 20/4 AD-A186 440 2/5 UNCLASSIFIED NL .

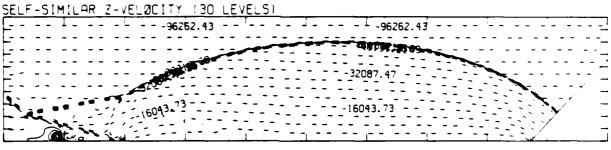


MURCHAR RESOLUTION TEST OF HE NATIONAL AREA OF THE ARE

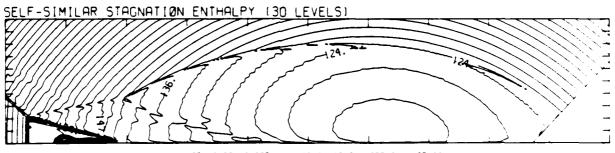
MS= 3.03 ALP=47.00 NR=575 NZ=120 KBEG= 75 PO=3.33E+04 PEFFEDT



-1.68E+05 TO 8.40E+04 STEP 8.40E+03 LABELS X1.0E-03



-9.63E+04 TØ 2.41E+04 STEP 4.01E+03 LABELS X1.0E+00.

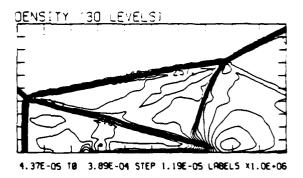


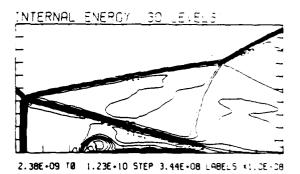
1.196+10 TØ 2.026+10 STEP 2.886+08 LABELS *1.06-08

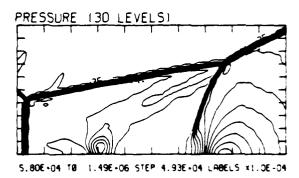
Figure $7d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.33 ALP=47.30 [L=434 [R=558 UT= 61 PD=3.33E+34 PERFECT







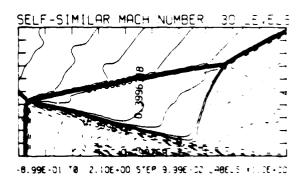
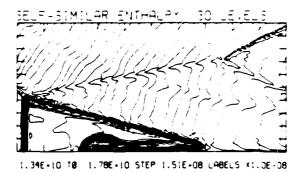
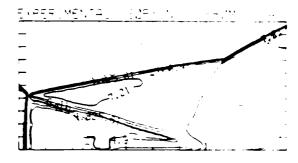


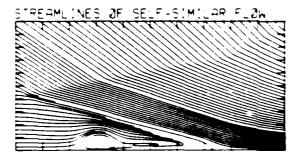
Figure 7e_p. Blowup-frame plots; $\gamma = 1.4$

Figure 7. Case 4, M_s = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.03 ALP=47.00 (L=434 (R=558 UT= 8) F0=3.3 54.4 PARE







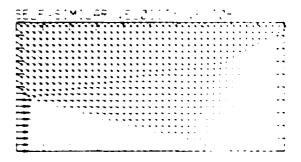


Figure $7e_p$. 3lowup-frame plots; v = 1.4 - continued.

Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.03 ALP=47.00 NP=575 NZ=120 KBEG= 75 F0=3.335474 HANDE

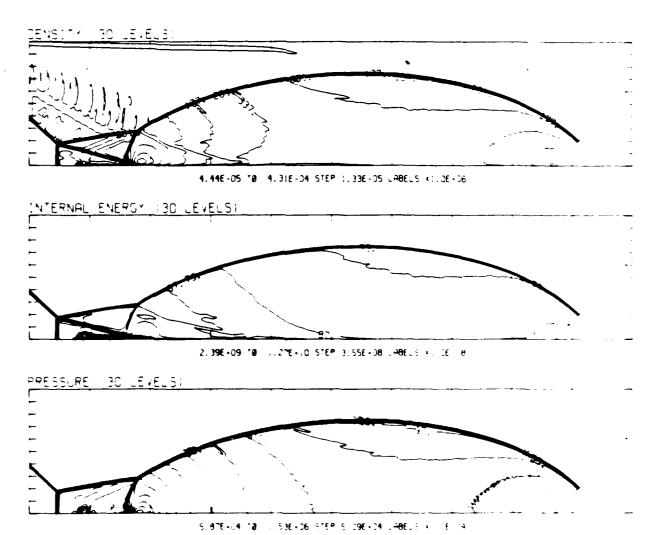


Figure $7d_{H^+}$. Whole-flowfield contour-plots; Hansen

Figure 7. Case 4, M_s = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.03 ALP=47.00 NR=575 NZ=120 ABEG= 75 F0=3.33E+0= 540 A

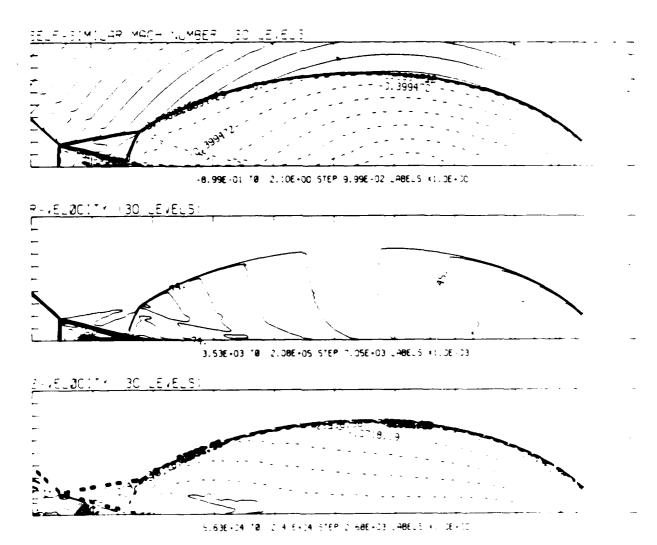


Figure 7d_H. Whole-flowfield contour-plots; Hansen - continued.

Figure 7. Case 4, M_s = 3.03, θ_W = 47°, Air, v = 1.4 and Hansen EOS, DMR - continued.

MS= 3.08 ALP=47.00 NP=575 NZ=120 KBE3= 75 F0=3.33E4,4 HAY

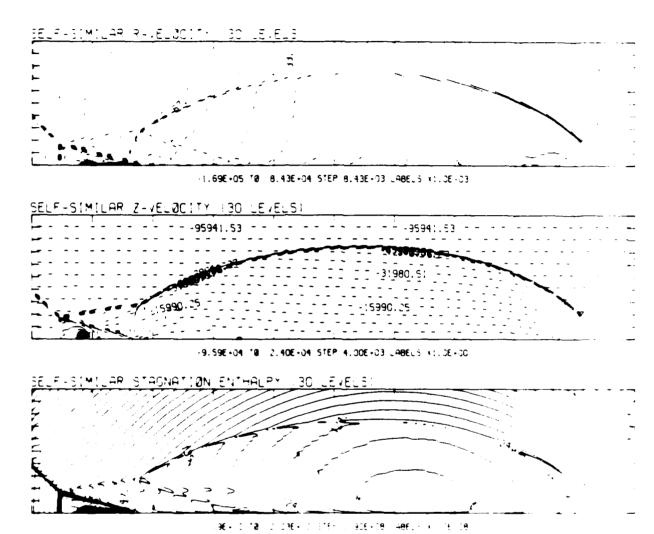
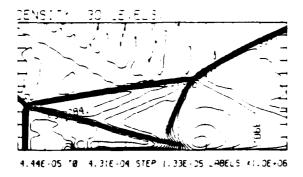
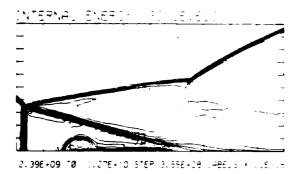


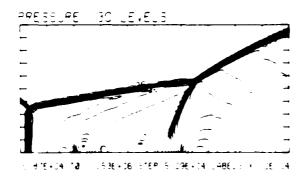
Figure $7d_{\text{H}}$. Whole-flowfield contour-plots; Hansen - continued.

Figure 7. Case 4, M_s = 3.03, θ_w = 47°, Air, γ = 1.4 and Hansen EDS, DMR - continued.

MS= 3.03 ALP=47.00 (L=403 (R=550 UT= 30 A0%), 875+14 H47 H.







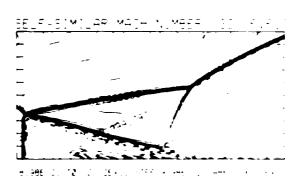
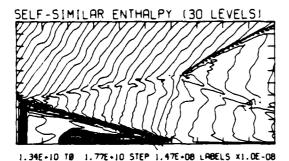
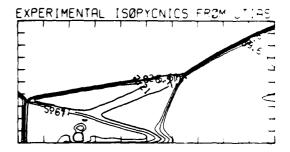


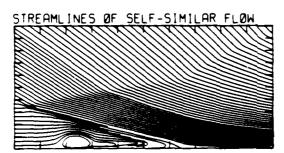
Figure $7e_{\mathrm{H}}$. Rlowup-frame plots; Hansen

Figure 7. Case 4, M_S = 3.03, θ_W = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.03 ALP=47.00 IL=429 IR=552 JT= 60 P0=3.33E+04 HANSE'.







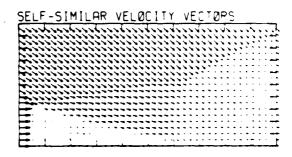
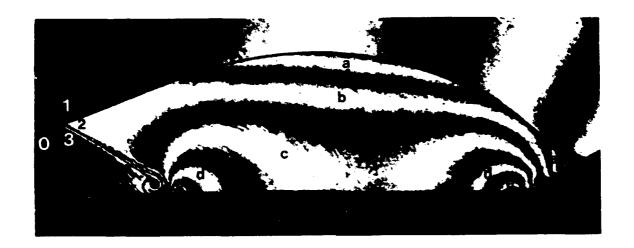


Figure 7e_H. Blowup-frame plots; Hansen - continued.

Figure 7. Case 4, M_S = 3.03, $\theta_{\rm W}$ = 47°, Air, γ = 1.4 and Hansen EOS, DMR - continued.



Region	٥/٥
0	1.00
1	3.51
2	4.89
3	4.14
a	4.69
Ъ	4.89
С	5.09
d	5.28
e	5.48
f	4.49

Figure 8a. Interferogram

MB= 0.65 ALP=30.00 NP=480 NZ=125 KBEG= 90 PC=1.33E+15 FFAFF

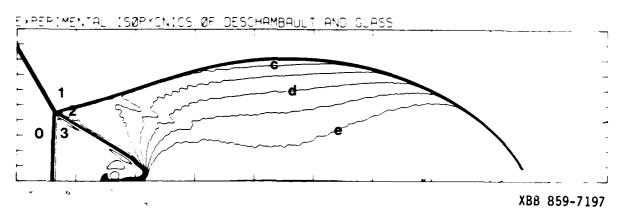


Figure $8b_p$. Calculated isopycnics ($\gamma = 1.4$) using the experimental fringes

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR.

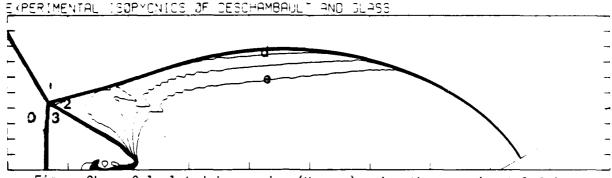


Figure $8b_{\mathrm{H}}$. Calculated isopycnics (Hansen) using the experimental fringes.

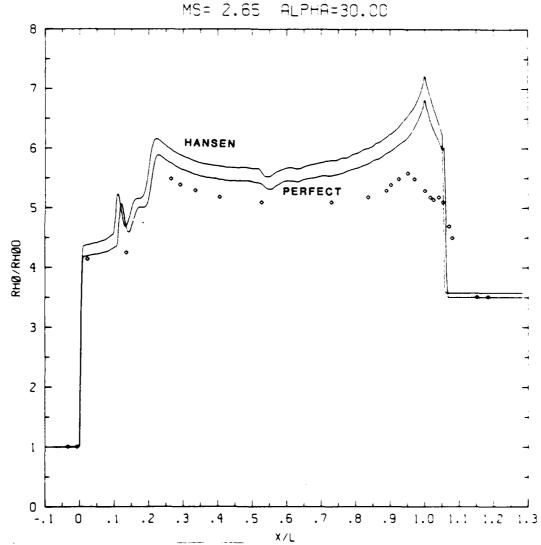


Figure 8c. Wall plot for p/p_0 , $\gamma = 1.4$ and Hansen calculations with experimental data.

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

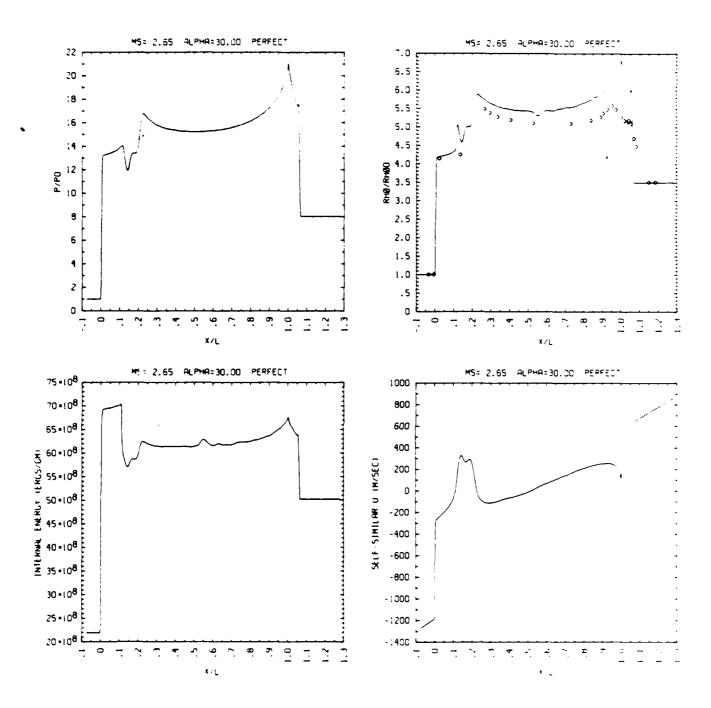


Figure 8cp. Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e, u; γ = 1.4.

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

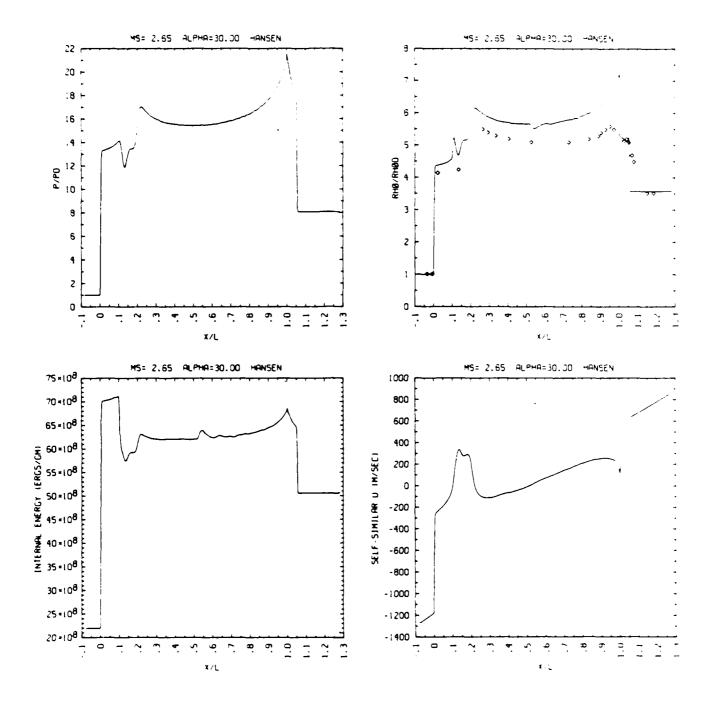


Figure 8c_H. Wall plot for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

MS= 0.65 ALP=30.00 NR=480 NZ=105 KBEO= 90 PD=1.33E+13 AFRA-51

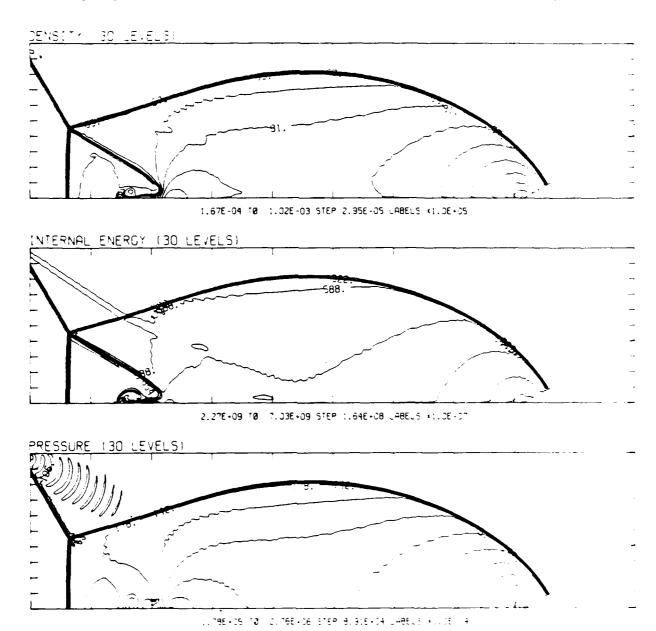
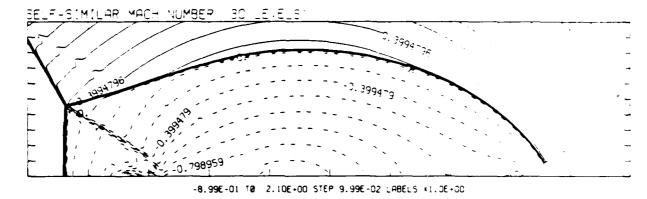
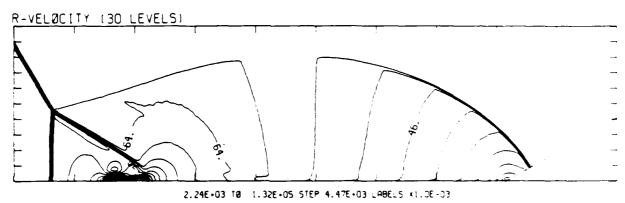


Figure 8d_p. Whole-flowfield contour-plots; $\gamma = 1.4$

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

MS= 2.65 ALP=30.00 NR=480 NZ=105 ABEG= 90 F0=1.835+05 ABEAE1





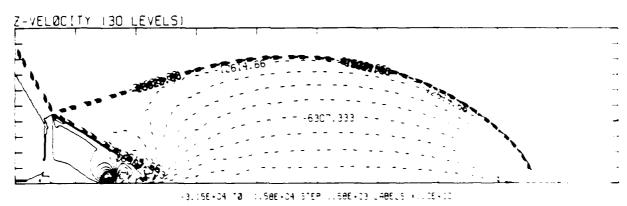
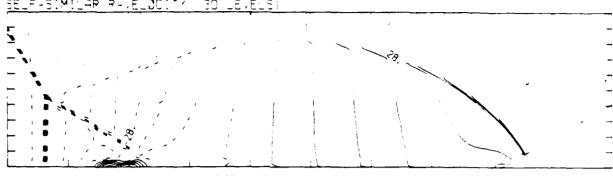


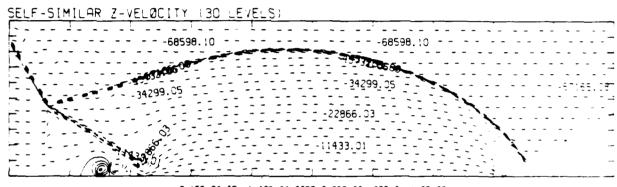
Figure $8d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 8. Case 5, M = 2.65, $\theta_{\rm W}$ = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

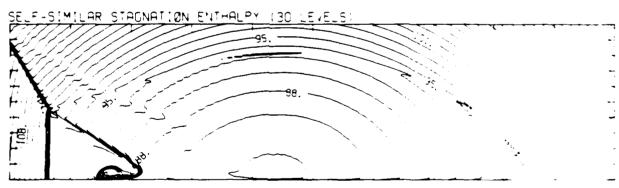
MS= 0.85 ALP=30.00 NR=480 NB=105 KBEG= 90 PD=1.33E+35 AEAA511



-1.21E+05 TØ 9.27E+04 STEP 7.13E+03 LABELS *1.0E-03



-7.15E+04 T0 1.43E+04 STEP 2.86E+03 LABELS X1.0E+00



8.5:E+39 18 ...34E+10 51EP ...69E+38 ..48ELS (1.3E ...

Figure $8d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

MS= 0.85 ALP=30.00 (L=378 .P=454 UT= 3) P1=.// P3=.//

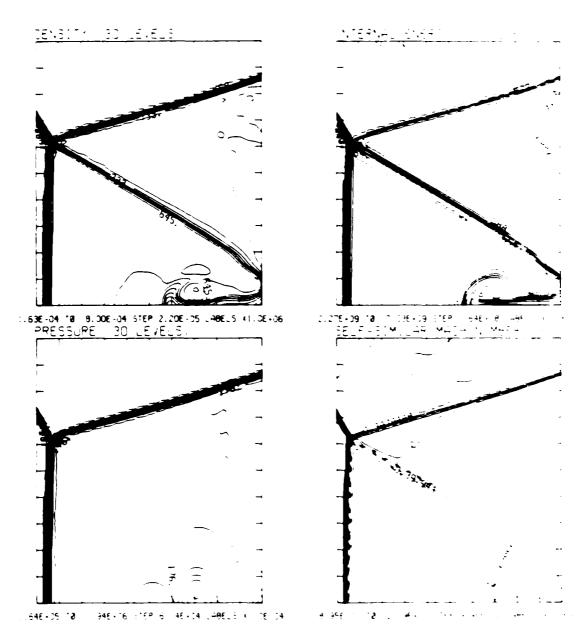


Figure $8e_p$. Blowup-frame plots; y = 1.4

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

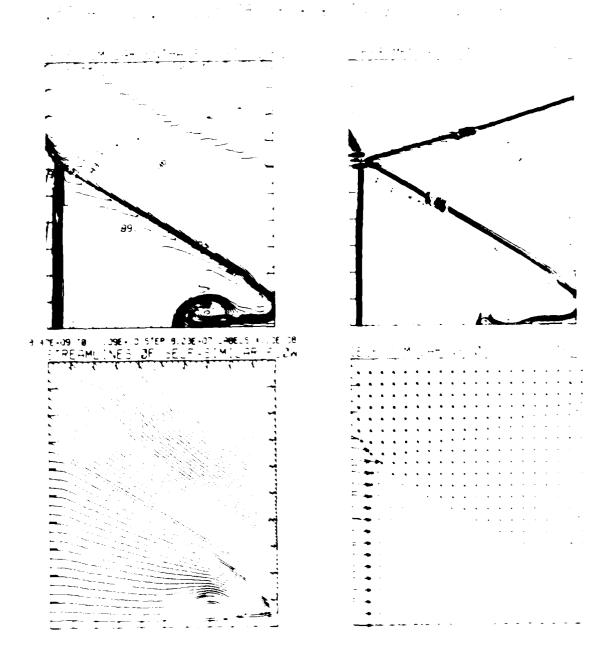


Figure Bep. 3lowup-frame plots; v = 1.4 - continued.

Figure 8. Case 5, M_S = 2.65, A_M = 30°, Air, v = 1.4 and Hansen EDS, CMR - continued.

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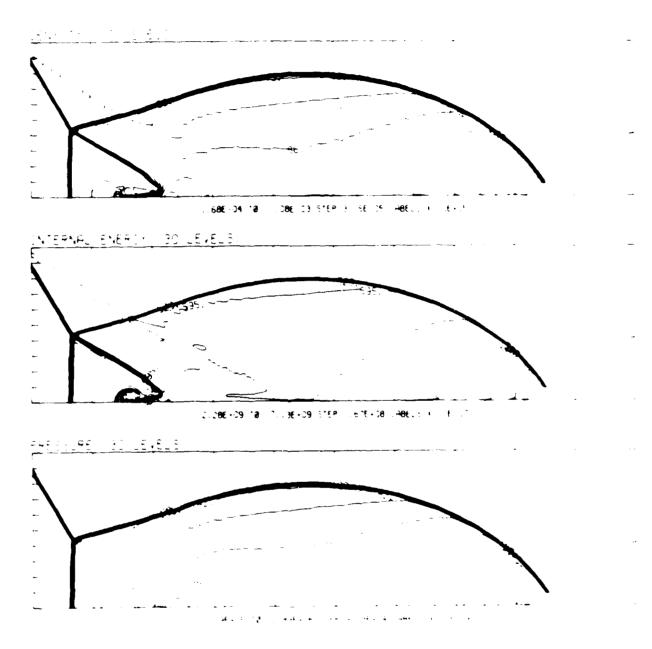


Figure 8d_H. Whole-flowfield contour-plots; Hansen

Figure 8. Case 5, M_s = 2.65, θ_w = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

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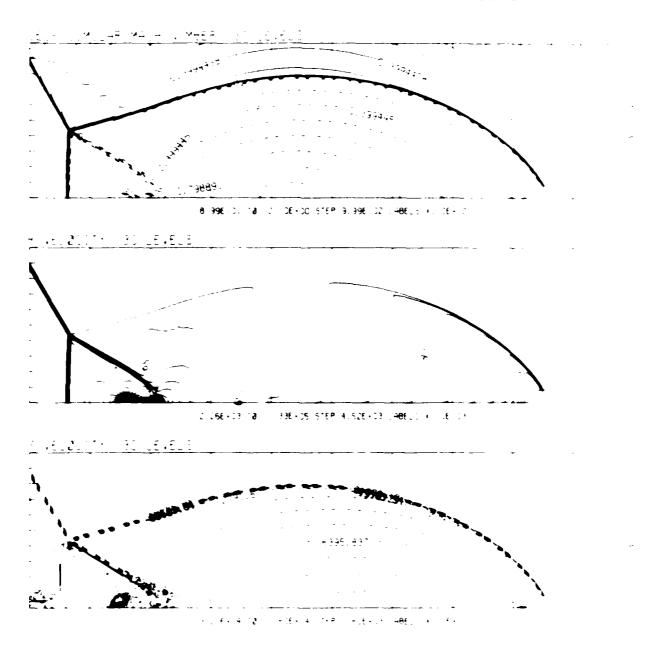


Figure 8d_H. Whole-flowfield contour-plots; lansen - continued.

Figure 8. Case 5, M_S = 2.65, θ_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

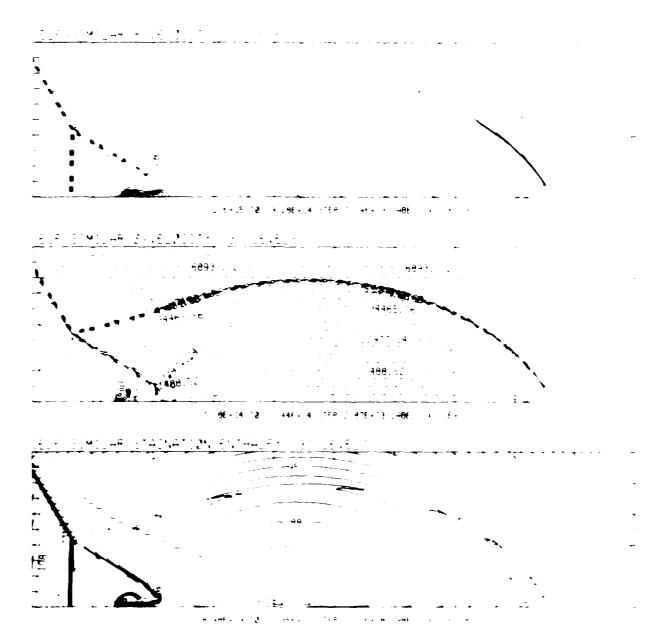


Figure 8d_H. Whole-flowfield contour-plots; Hansen - continued.

Figure 8. Case 5, M_s = 2.65, θ_w = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.

MRS (기유등 및 유료되었다. 이 100 기 기 유료교육의 기계를 보고 유기하는 사고 있다.

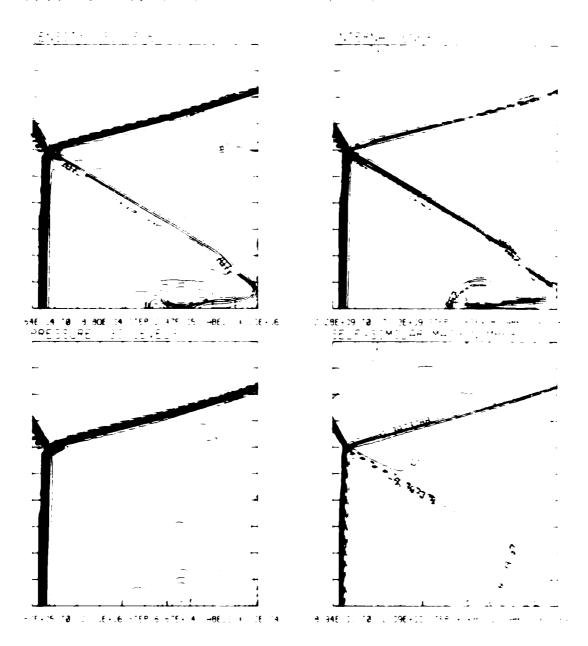


Figure Se_u. Blowup-frame plots; Hansen

Figure 8. Case 5, M_S = 2.65, A_W = 30°, Air, v = 1.4 and Hansen EOS, CMR - continued.

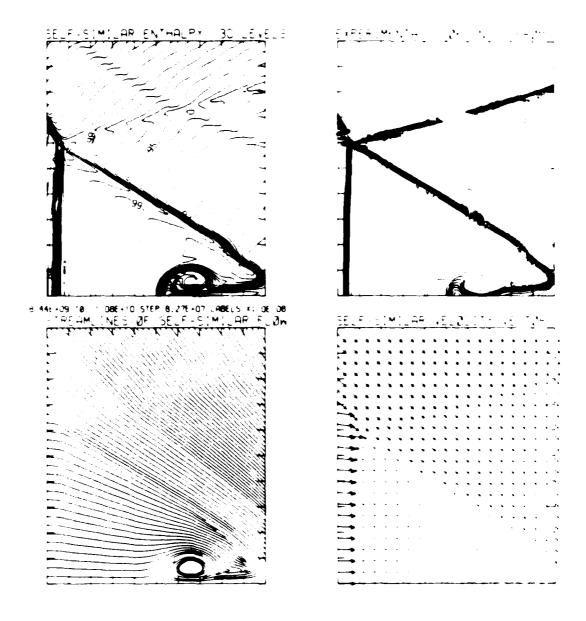
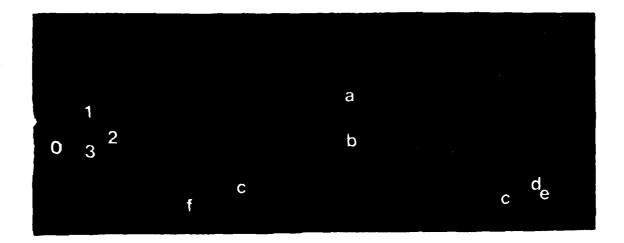


Figure 8e_N. Blowup-frame plots; Hansen - continued.

Figure 8. Case 5, M_S = 2.65, H_W = 30°, Air, γ = 1.4 and Hansen EOS, CMR - continued.



Region	0/00
0	1.00
1	3.59
2	4.83
3	3.74
а	4.83
ь	5.16
c	5.48
d	4.50
e	4.18
f	5.80

Figure 9a. Interferogram

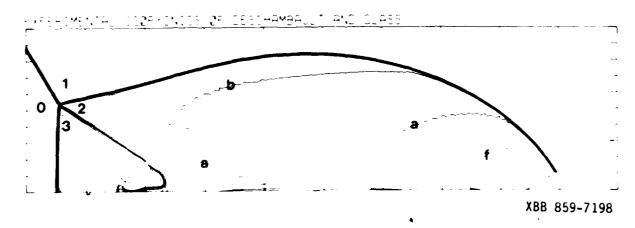


Figure 9b. Calculated isopycnics using the experimental fringes.

Figure 9. Case 6, $M_S = 5.07$, $A_W = 30^\circ$, Argon, $A_V = 5/3$, CMR.

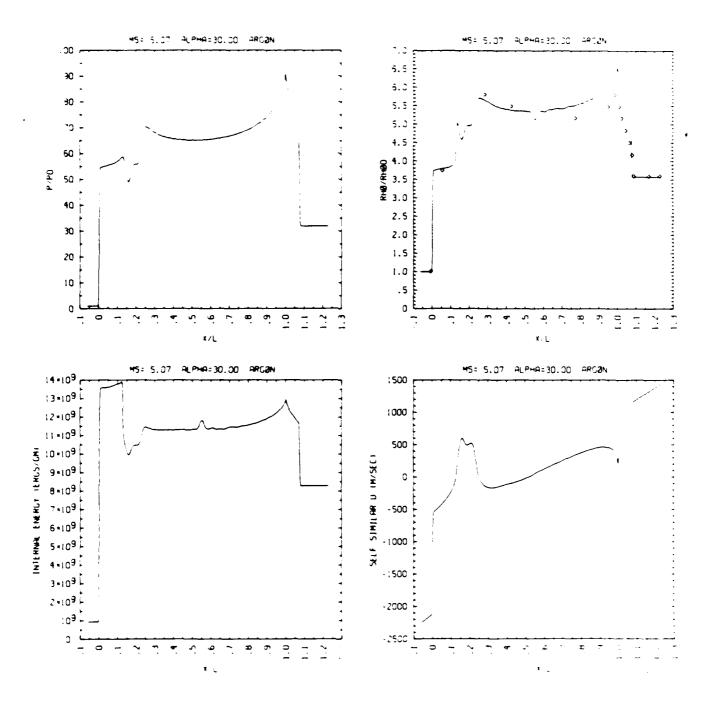


Figure 9c. Wall plots for p/p_0 , p/p_0 with experimental data included, e, \bar{u} . Figure 9. Case 6, M_S = 5.07, θ_W = 30°, Argon, γ = 5/3, CMR - continued.

MS= 5.0T ALF=30.00 NR=500 NZ=140 KBED= 80 PO=4.00E-04 AFIDY.

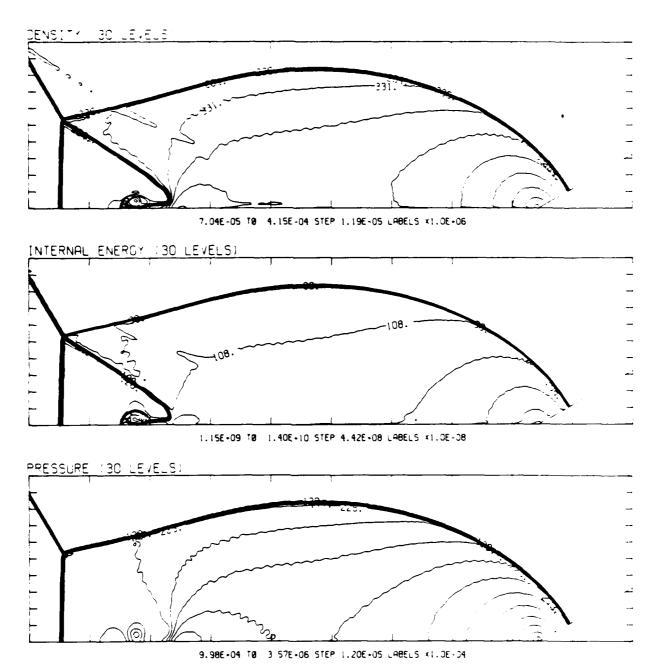
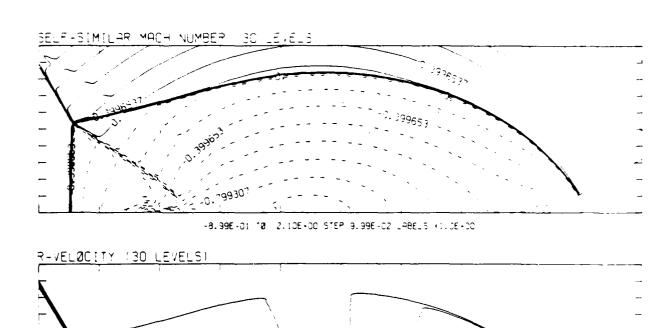
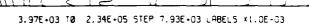


Figure 9d. Whole-flowfield contour-plots.

Figure 9. Case 6, M_S = 5.07, θ_W = 30°, Argon, γ = 5/3, CMR - continued.





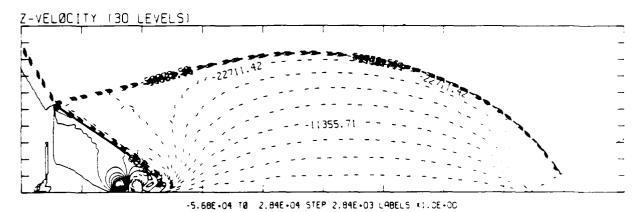
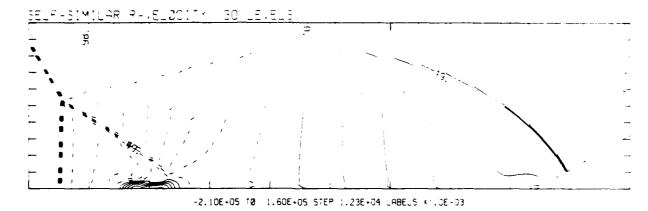


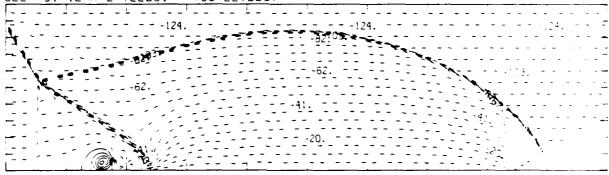
Figure 9d. Whole-flowfield contour-plots - continued.

Figure 9. Case 6, $M_S = 5.07$, $\theta_W = 30^{\circ}$, Argon, $\gamma = 5/3$, CMR - continued.

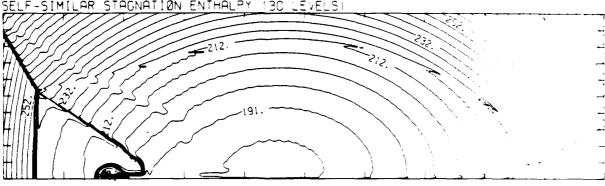
MS= 5.07 ALP=30.00 NR=500 NZ=140 MBEG= 80 F0=4,005+04 A-



<u> ELF-SIMILAR Z-VELØCITY (30 LEVELS)</u>



-1.30E+05 TØ 2.59E+04 STEP 5.18E+03 LABELS X1.0E-03



1.82E+10 10 3.28E+10 STEP 5.05E+08 LABELS x1.0E-08

Figure 9d. Whole-flowfield contour-plots - continued.

Figure 9. Case 6, M_S = 5.07, θ_W = 30°, Argon, γ = 5/3, CMR - continued.

MS= 5.07 ALP=30.00 [L=374 [P=477 LT=100 F0=4.005+14 AF104.

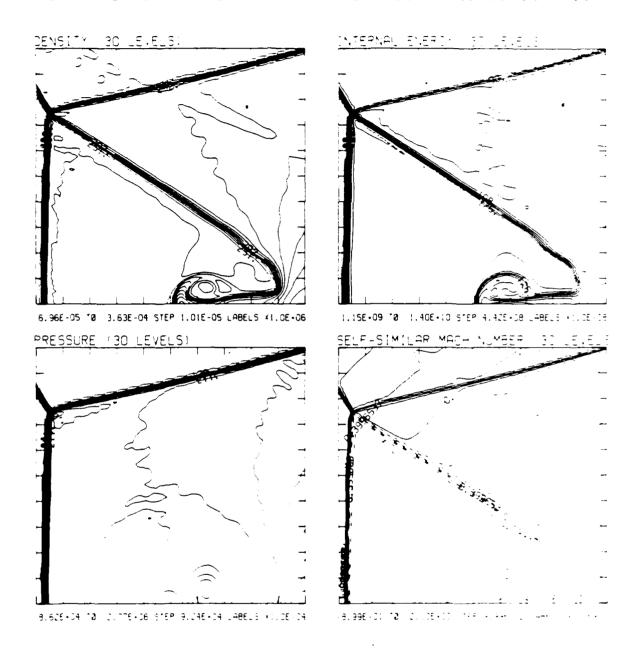


Figure 9e. Blowup-frame plots.

Figure 9. Case 6, M_S = 5.07, θ_W = 30°, Argon, γ = 5/3, CMR - continued.

MS= 5.07 ALP=30.00 (L=374 [P=477 UT=100 F0=4.01E+14 A+ M

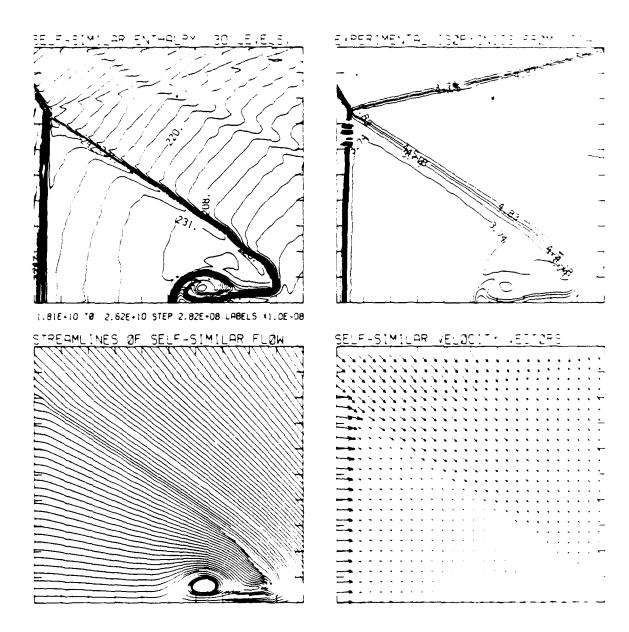


Figure 9e. Blowup-frame plots - continued.

Figure 9. Case 6, $M_S = 5.07$, $\theta_W = 30^\circ$, Argon, $\gamma = 5/3$, CMR - continued.

```
1 1'
0 2
3 a b c d e f g h;
```

Region	2/2 ₀
0	1.00
1	5.73
1'	6.33
2	6.30
3	5.77
a	6.70
Ъ	7.10
С	7.50
d	7.90
e	8.29
f	8.69
g	9.09
h	9.49
i j	9.89
j	10.29

Figure 10a. Interferogram.

MS=10.37 ALP=10.00 NP=475 NZ=140 KBE3= 75 P0=8.875+04 H400+7

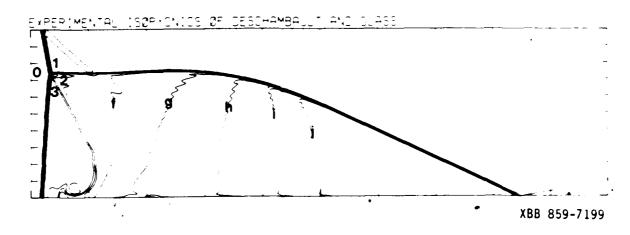


Figure 10b. Calculated isopycnics using the experimental fringes.

Figure 10. Case 7, M_S = 10.37, θ_W = 10°, Air, Hansen EOS, CMR.

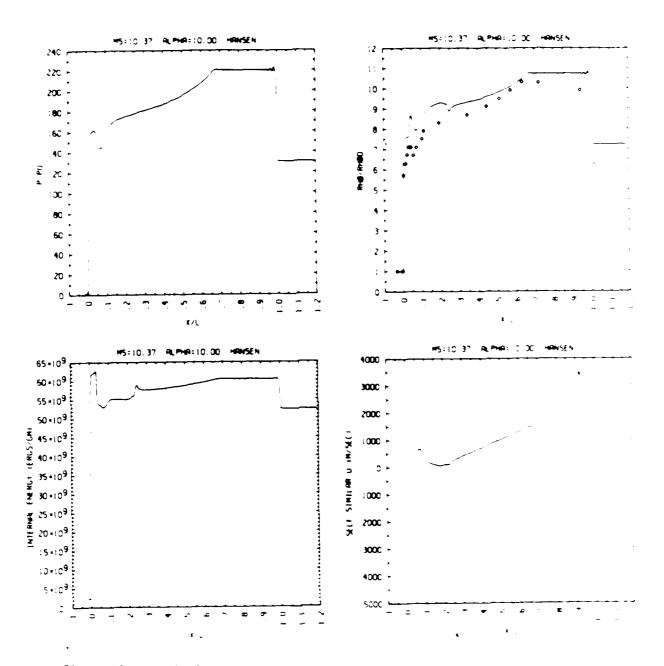


Figure 10c. Wall plots for p/p_0 , p/p_0 with experimental data included, e, \bar{u} .

Figure 10. Case 7, M_S = 10.37, θ_W = 10°, Air, Hansen EOS, CMR - continued.

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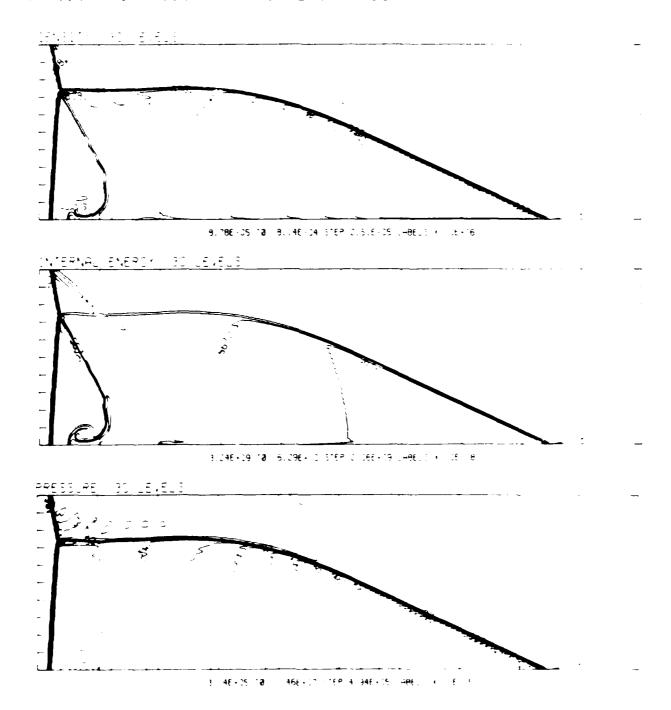


Figure 10d. Whole-flowfield contour-plots.

Figure 10. Case 7, $M_s = 10.37$, $A_w = 10^\circ$, Air, Hansen EOS, CMR - continued.

Mineral (97) 4 (Fet) 2.00 (AB 4978 (N.24), 40 (ABBC4) 78 (Ft) (ABBC4) (ABBC4)

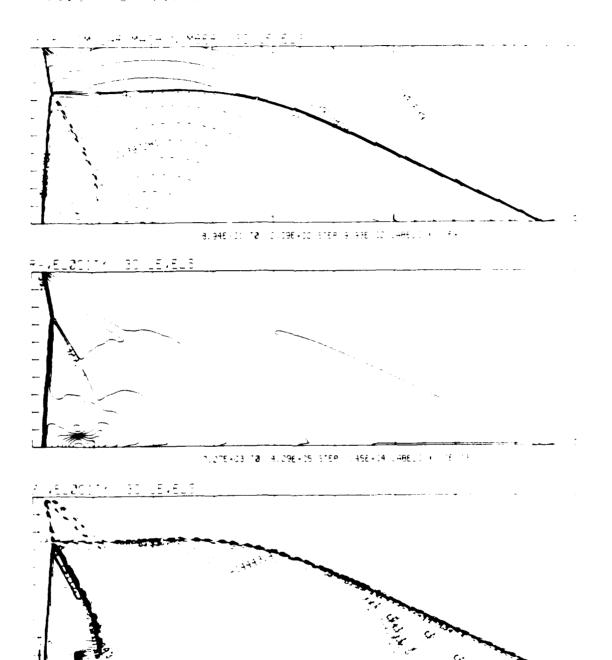


Figure 10d. Whole-flowfield contour-plots - continued.

Figure 10. Case 7, M_S = 10.37, α_W = 10°, Air, Hansen EOS, CMR - continued.

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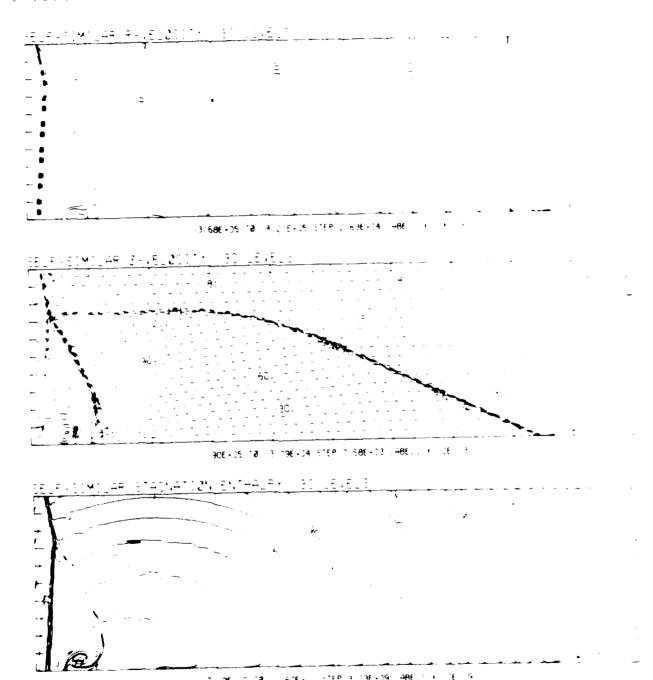


Figure 10d. Whole-flowfield contour-plots - continued.

Figure 10. Case 7, M_S = 10.37, A_W = 10°, Air, Hansen EOS, EMR - continued.

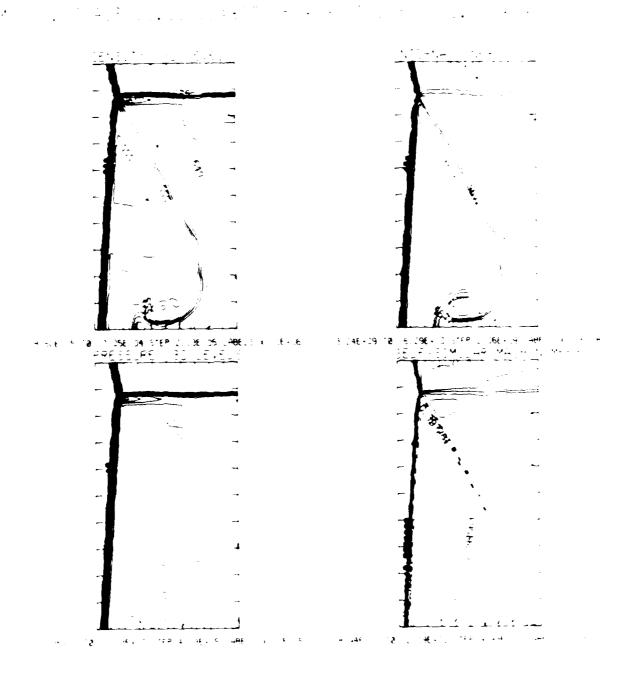


Figure 10e. Blowup-frame plots.

Figure 10. Case 7, M_S = 10.37, A_M = 10°, Air, Hansen EOS, CMR - continued.

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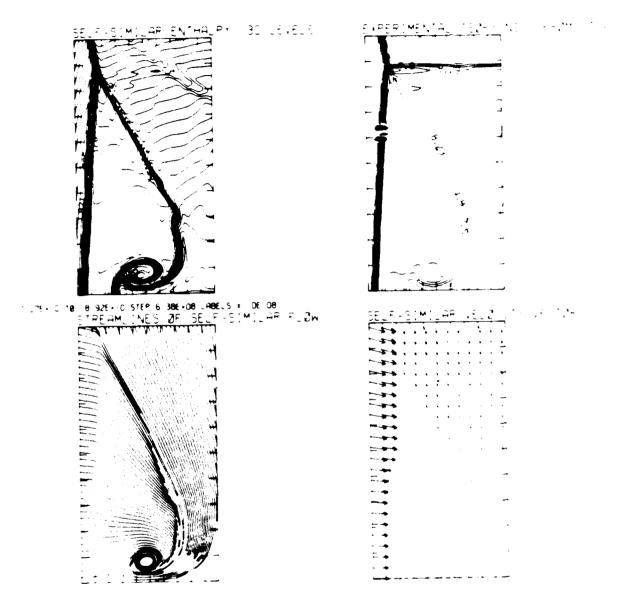
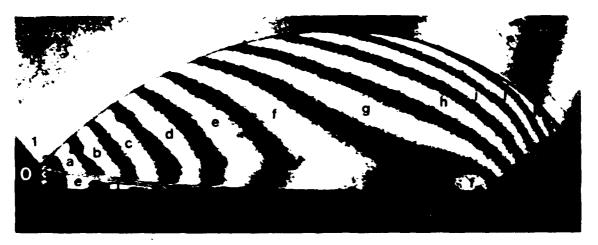


Figure 10e. Blowup-frame plots - continued.

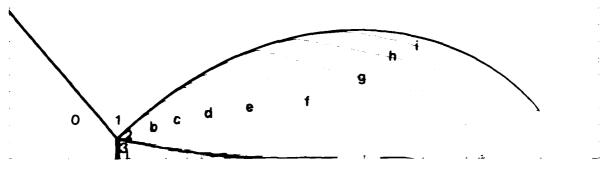
Figure 1). Case 7, M_S = 10.37, A_W = 10°, Air, Hansen EOS, CMR - continued.



Region	./.c	Region	e/e
0	1.00	e	3.04
1	2.13	f	2.96
2	3.44	g	2.88
3	3.09	ĥ	2.80
а	3.36	i	2.72
Ь	3.28	j	2.65
С	3.20	k	2.57
d	3.12	1	2.49

Figure 11a. Interferogram.

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XBB 859-7200

Figure 11th. Calculated isopycnics (v=1.4) using the experimental fringes.

Figure 11. Case 3, M_S = 1.66, n_W = 40°, Air, v = 1.4 and Hansen EOS, SMR.

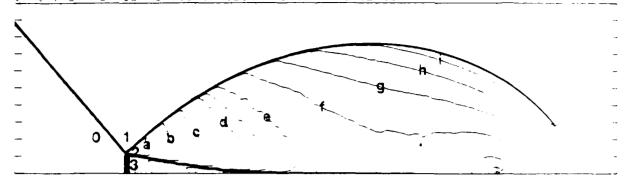


Figure $11b_{\mbox{\scriptsize H}}$. Calculated isopycnics (Hansen) using the experimental fringes.

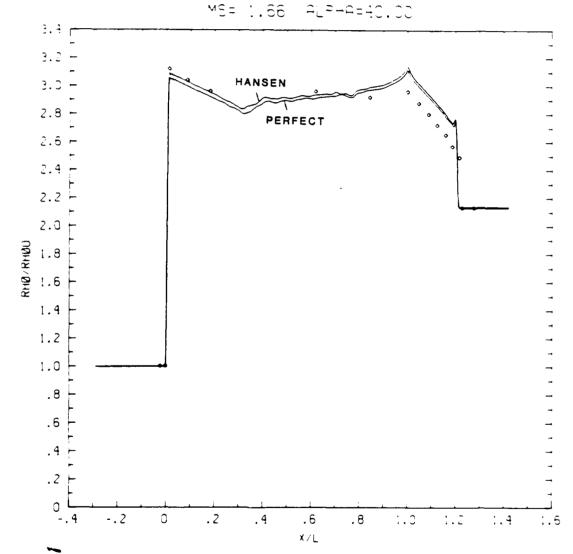


Figure 11c. Wall plot for p/p_0 , γ = 1.4 and Hansen calculations, with experimental data.

Figure 11. Case 8, M_s = 1.66, θ_w = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

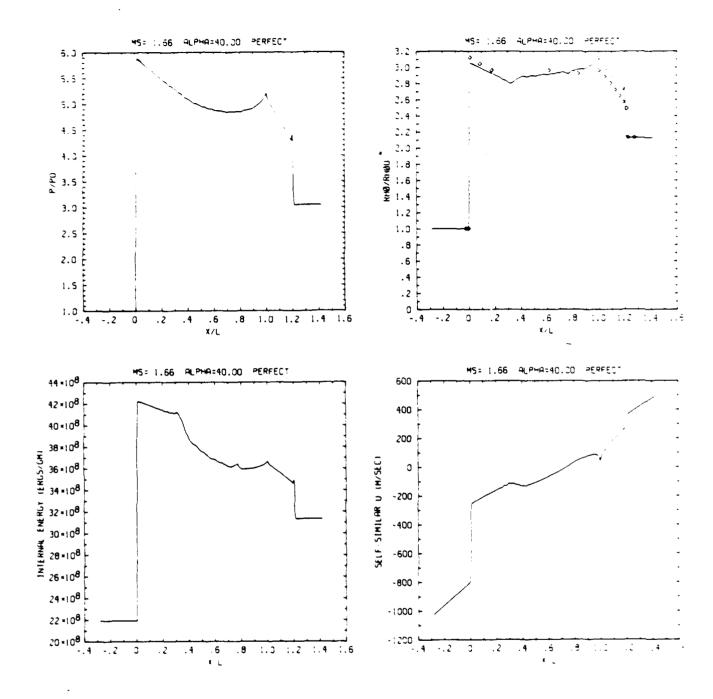


Figure 11cp. Wall plots for p/p_0 , p/p_0 with experimental data included, e, u; γ = 1.4.

Figure 11. Case 8, Ms = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

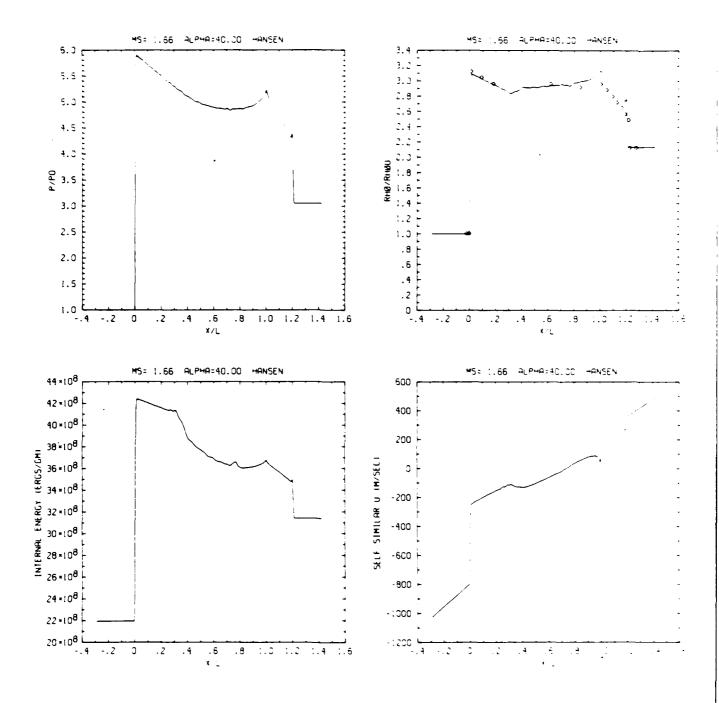


Figure 11c_H. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.66 ALP=40.00 NR=470 NZ=135 KBEG= 95 PD=3.33E+35 PEFFEDT

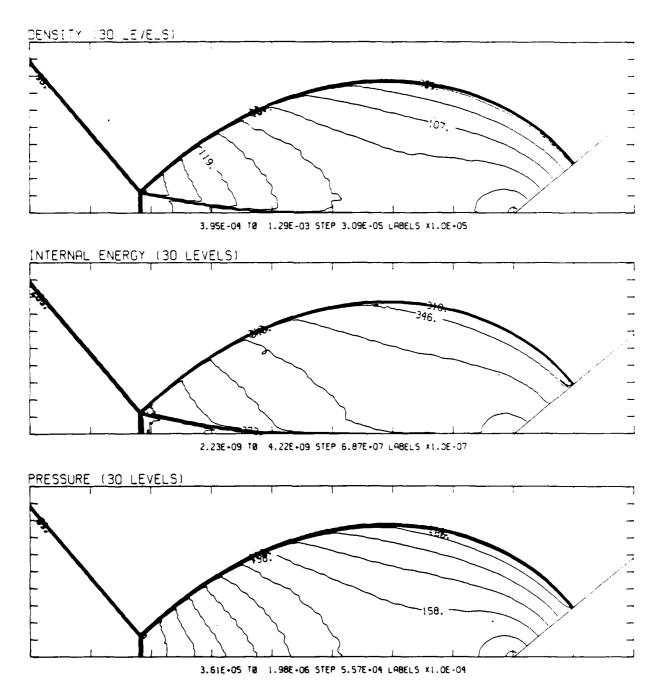
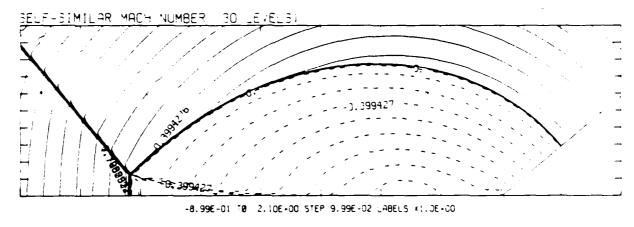
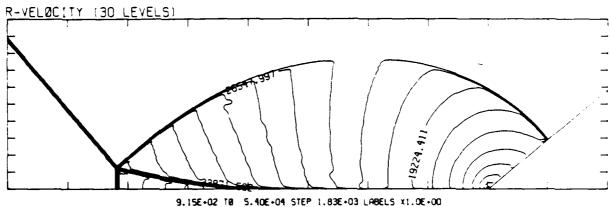


Figure 11dp. Whole-flowfield contour-plots; $\gamma = 1.4$.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.66 ALP=40.00 NR=470 NZ=135 KBEG= 95 PG=3.33E+05 PEFFEDT





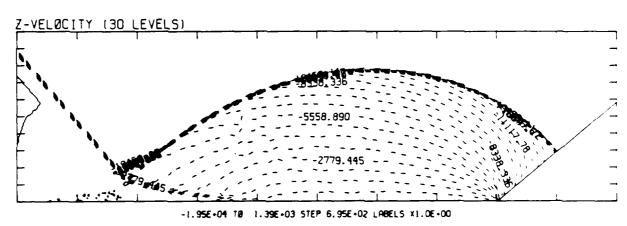
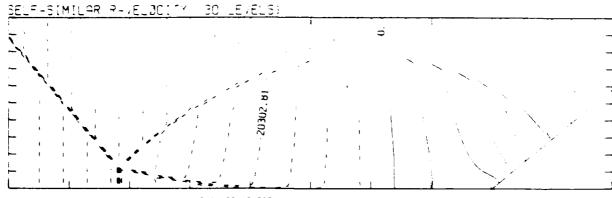
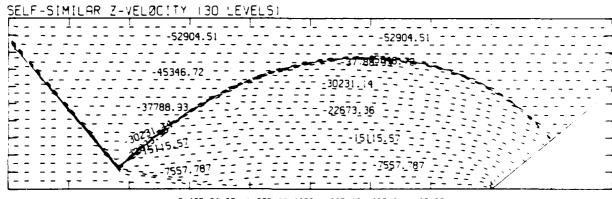


Figure 11d_p. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.



-9.64E+04 T0 5.58E+04 STEP 5.08E+03 LABELS X1.0E+00



-5.48E+04 T0 1.89E+03 STEP 1.89E+03 LABELS X1.0E+00

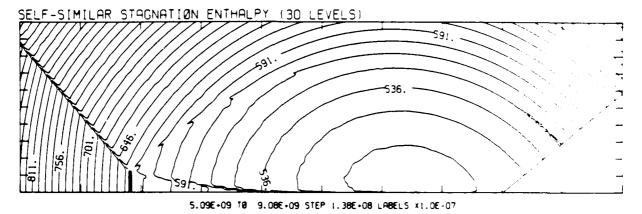
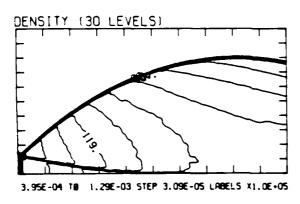
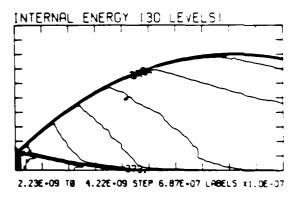


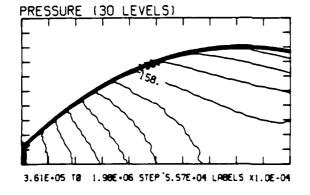
Figure 11dp. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.66 ALP=40.00 IL=150 IR=387 UT=130 PC=3.33E+05 PERFECT







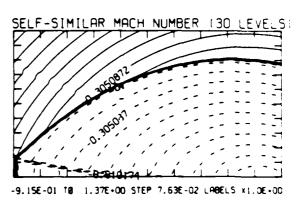
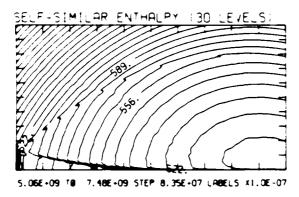
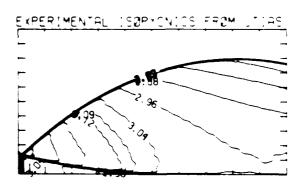


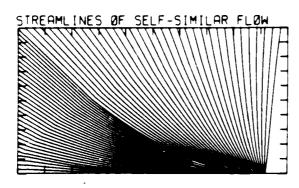
Figure 11ep. Blowup-frame plots; $\gamma = 1.4$.

Figure 11. Case 8, M_S = 1.66, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.86 ALP=40.00 (L=150 (P=387 UT=130 PD=3.33E+05 PERFEDT







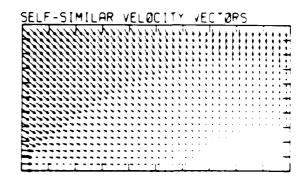


Figure $11e_p$. Blowup-frame plots; $\gamma = 1.4$ - continued.

Figure 11. Case 8, M_S = 1.66, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.88 ALR=40.00 NR=470 NZ=135 W8574 N8 A PRIVATE

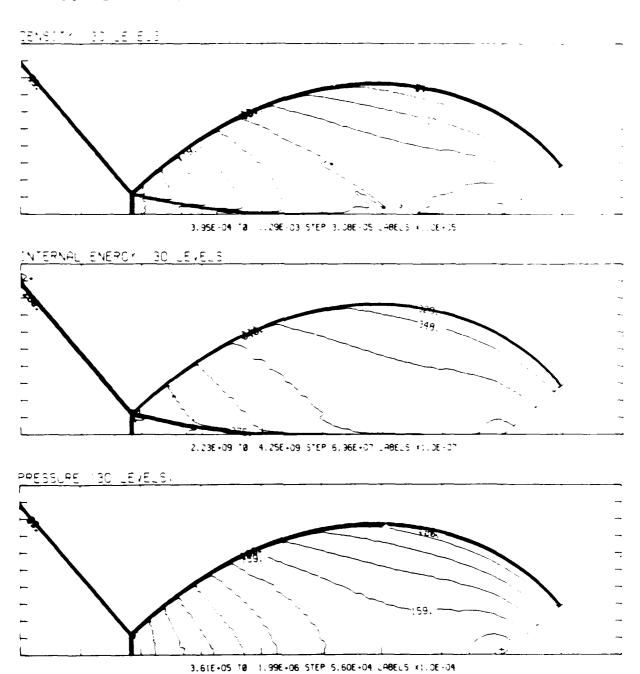


Figure 11d_H. Whole-flowfield contour-plots; Hansen.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}\approx$ 40°, Air, γ = 1.4 and Hansen EQS, SMR - continued.

M35 1.68 ALF643.00 NF6470 NZ6135 KEED4 35 ACMA, Abb A HAV

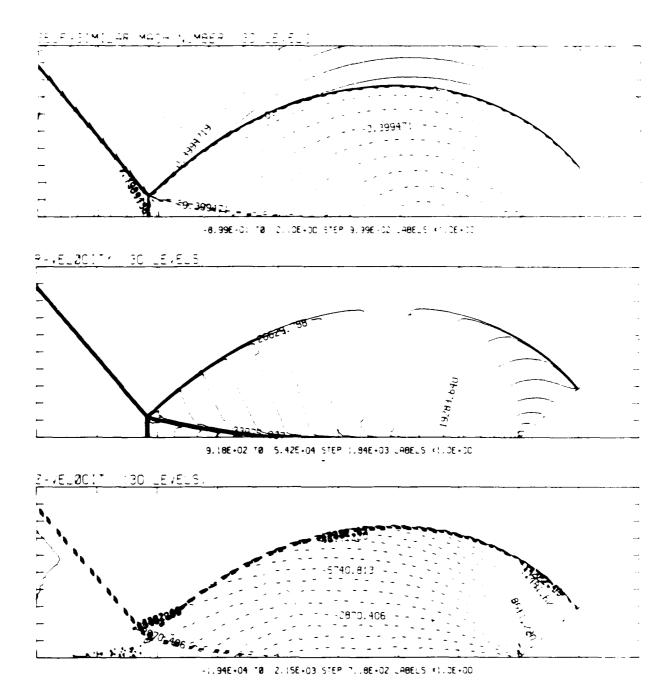
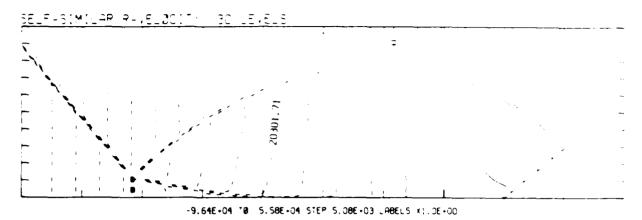
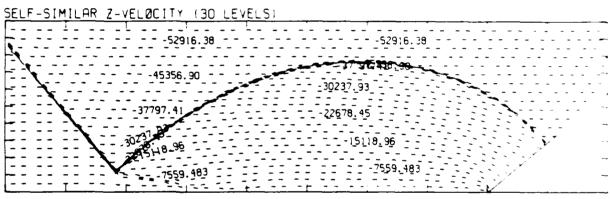


Figure 11 d_H . Whole-flowfield contour-plots; Hansen - continued.

Figure 11. Case 8, M_S = 1.66, $\theta_{\rm W}$ = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.88 ALP=40.00 NR=470 NZ=185 MBED= 85 F0=3.885 AB HANGE







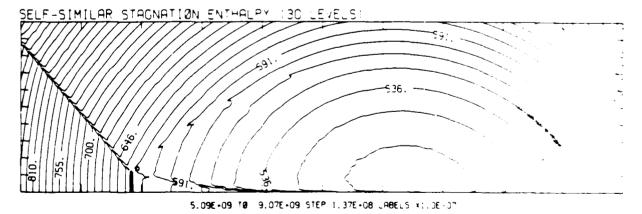
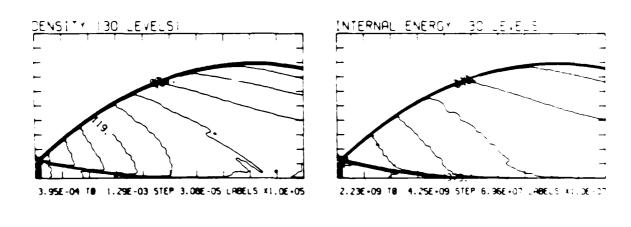


Figure 11d_H. Whole-flowfield contour-plots; Hansen - continued.

Figure 11. Case 8, M_S = 1.66, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.86 ALP=40.00 (L=150 (R=387 LT=130 PD=3.335+15 HANSEN



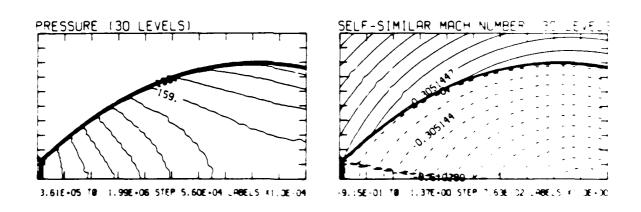
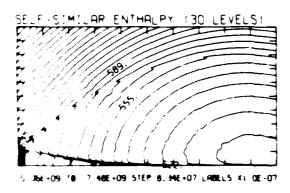
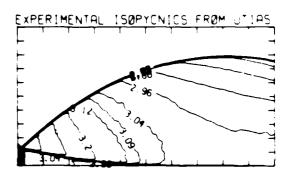


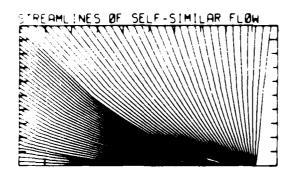
Figure 11e_H. 3low-up frame plots; Hansen.

Figure 11. Case 8, M_s = 1.66, θ_w = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.

MS= 1.66 ALP=40.00 [L=150 [R=387 UT=130 PO=3.33E+05 HANSEN







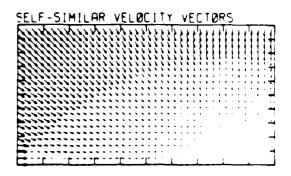


Figure 11eH. Blow-up frame plots; Hansen - continued.

Figure 11. Case 8, M_S = 1.66, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, SMR - continued.



Region	3/3;
0	1.00
1	3.73
2	6.06
3	4.59
а	6.21
Ь	6.37
С	6.53
d	6.69
е	6.85
f	7.00
g	7.16
h	7.32
i	6.06
j	5.90

Figure 12a. Interferogram.

0 3 f g h
XBB 859-7201

Figure 1255. Calculated isopycnics (v = 1.4) using the experimental fringes.

Figure 12. Case 9, $M_s=2.37$, $\epsilon_w=40^\circ$, Air, v=1.4 and Hansen EOS, DMR.

MSE 0.87 A/PE40.00 NR=510 NZ=110 /BE0= 90 P0=1.87E+18 HH/J8E1

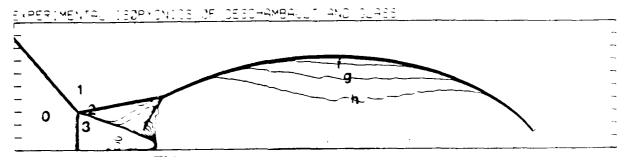


Figure 12b_H. Calculated isopycnics (Hansen) using the experimental fringes.

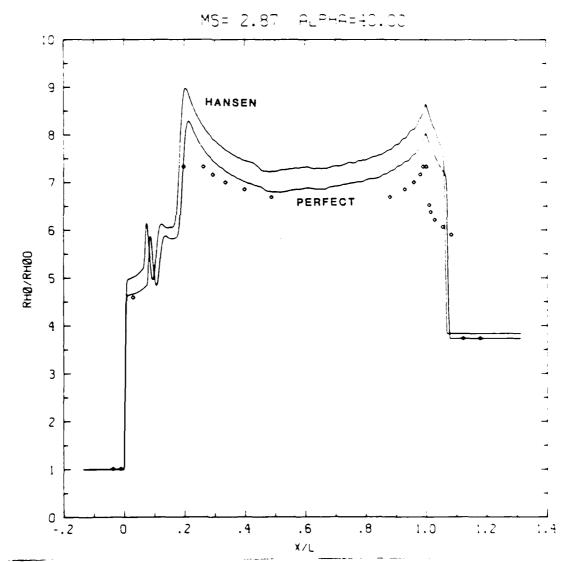


Figure 12c. Wall plots for p/p_0 , ρ/ρ_0 , $\gamma=1.4$ and Hansen calculations, with experimental data.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

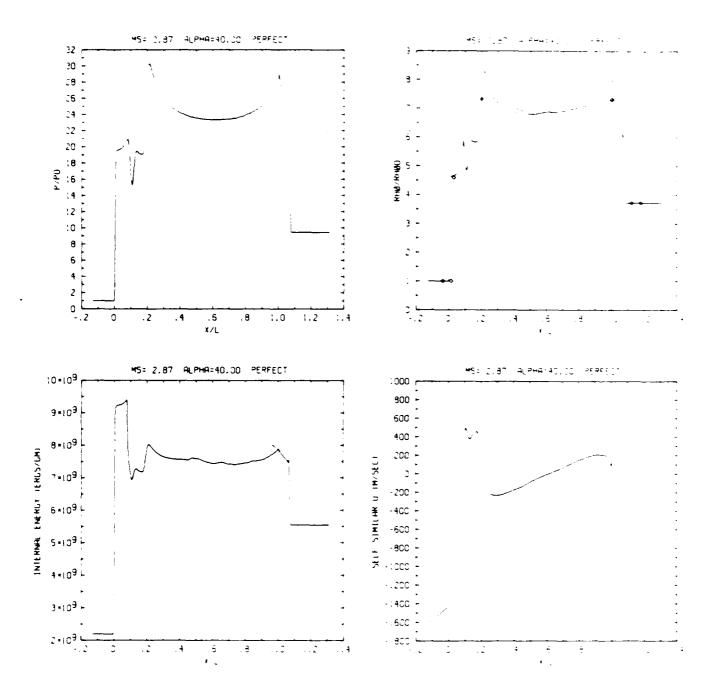


Figure 12c_p. Wall plots for p/p_0 , p/p_0 with experimental data included, e, u; $\gamma = 1.4$.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

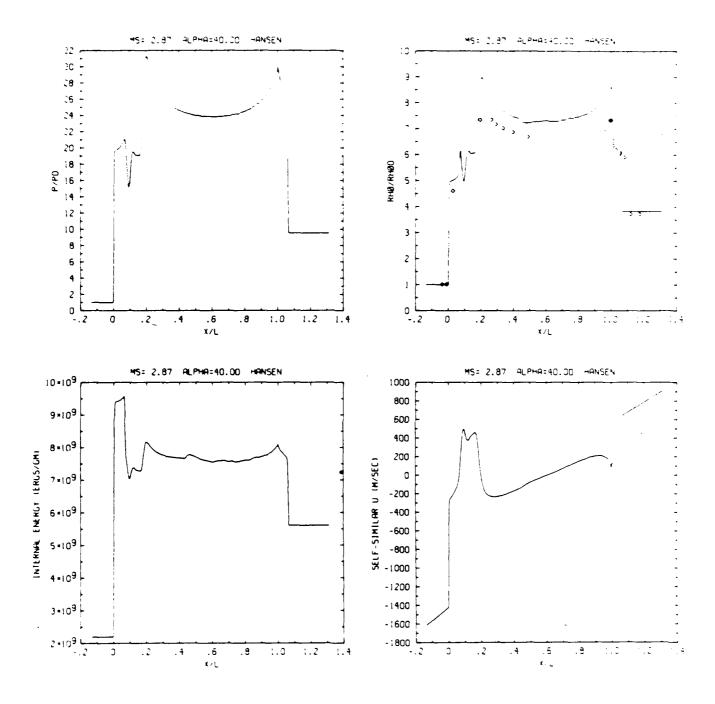
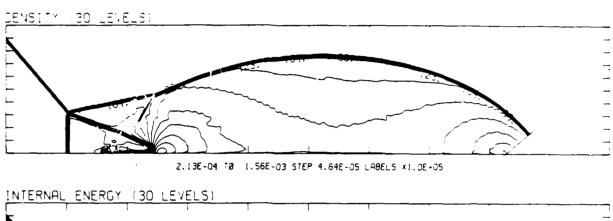
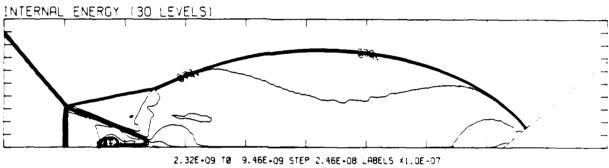


Figure 12cH. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 2.87 ALP=40.00 NR=510 NZ=110 KBEG= 90 FD=1.67E+05 FEFFED1





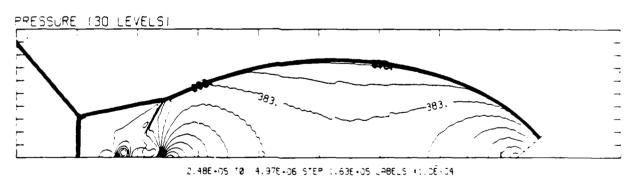
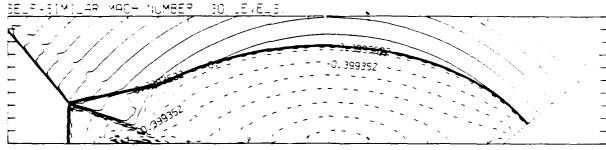


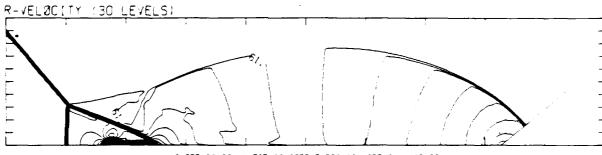
Figure 12dp. Whole-flowfield contour-plots; $\gamma = 1.4$.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

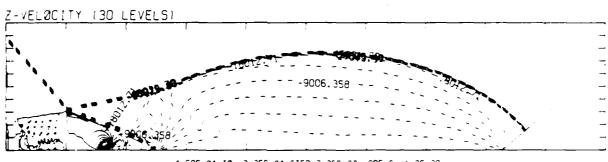
MS= 0.87 ALP=40.00 NR=510 NZ=110 MBE0= 90 PD=1.87E-08 PERFE



-8.99E-01 T0 2.10E+00 STEP 9.98E-02 LABELS <1.0E+00



2.95E+03 TØ 1.74E+05 STEP 5.90E+03 LABELS 41.0E-03

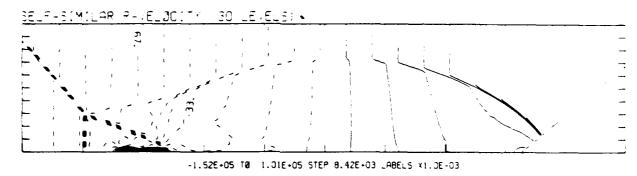


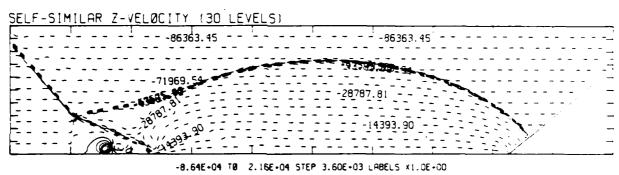
-4.50E+04 10 2.25E+04 STEP 2.25E+03 LRBELS <1.0E+00

Figure 12dp. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 2.87 ALP=40.00 MR=510 NZ=110 KBEG= 90 PG=1.67E+05 HANGEN





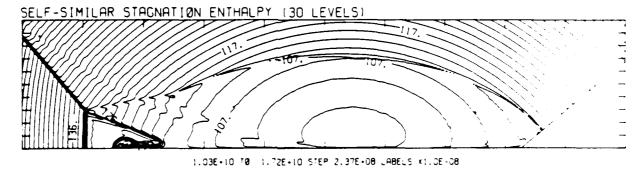
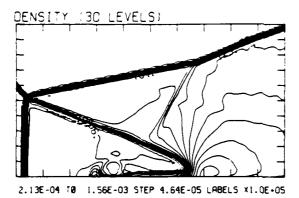
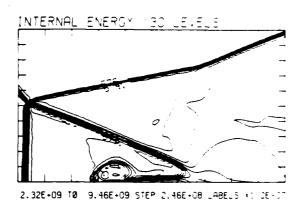


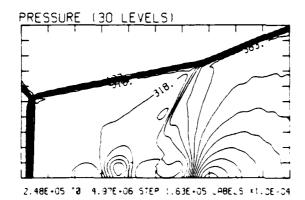
Figure 12dp. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 0.87 ALP=40.00 [L=347 [R=462 UT= 87 PD=1.87E+]5 REARE)*







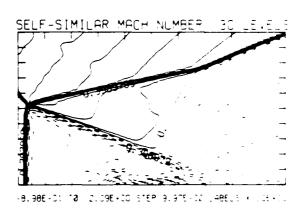
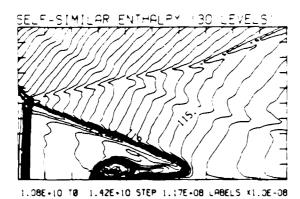
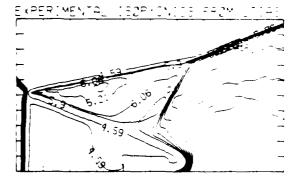


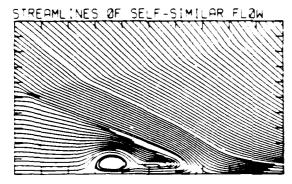
Figure $12e_p$. Blowup-frame plots; $\gamma = 1.4$.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 0.87 9_P=40.00 (L=347 (F=480 UT= 87 P0=),875+05 85545







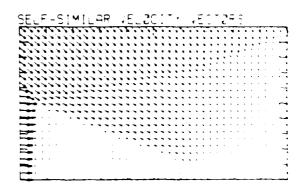


Figure 12e_p. Blowup-frame plots; y = 1.4 - continued.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 2.87 ALP=40.00 NR=510 NZ=110 MBEG= 30 F0=1.67E+35 HANSEN

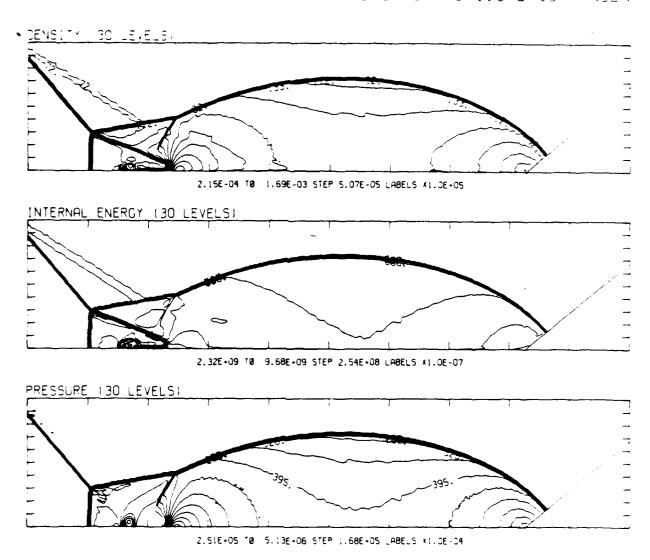


Figure 12d $_{\mathrm{H}}$. Whole-flowfield contour-plots; Hansen.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 0.87 ALP=40.00 NP=510 NZ=110 KBEG= 90 F0=1.875+05 HANGEY

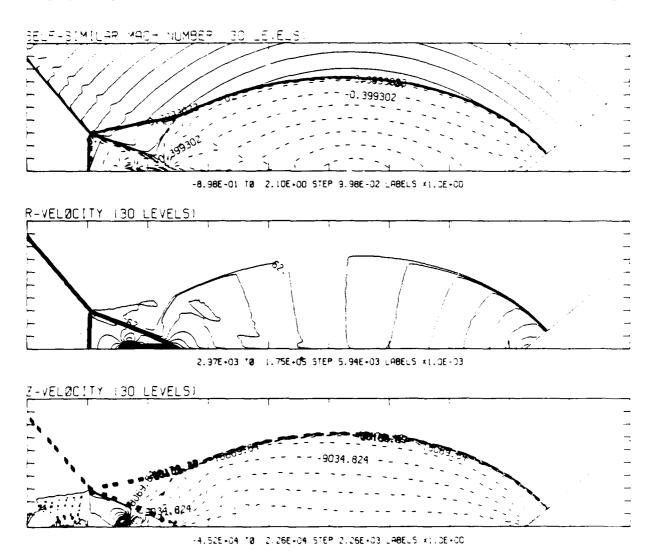
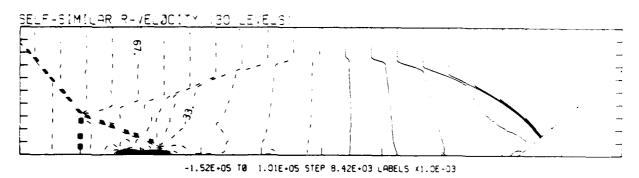


Figure 12dy. Whole-flowfield contour-plots; Hansen - continued.

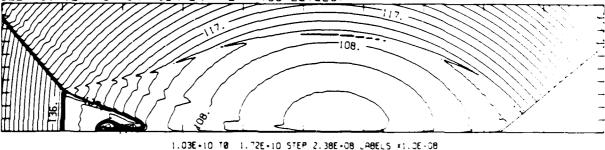
Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 2.87 ALP=40.00 NR=510 NZ=110 KBEG= 90 F0=1.67E+08 FE4FE



MILAR Z-VELØCITY (30 LEVELS)



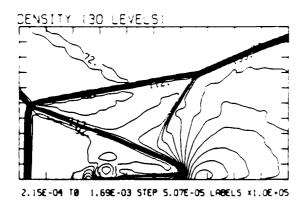


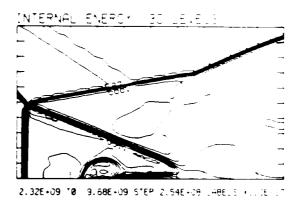
1.03E+10 T0 1.72E+10 STEP 2.38E+08 LABELS x1.3E+08

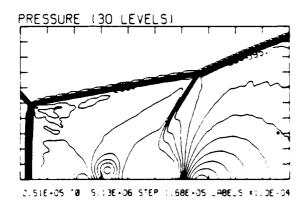
Figure 12d_H. Whole-flowfield contour-plots; Hansen - continued.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 0.87 ALP=40.00 (L=346 (R=46) UT= 87 PD=1.875+15 HANGE .







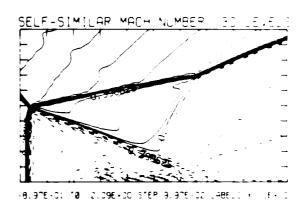
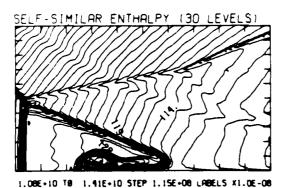
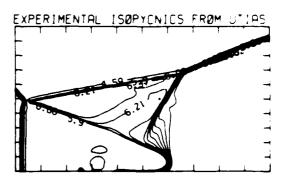


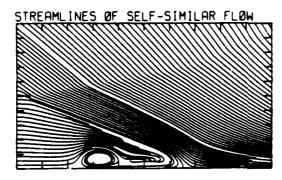
Figure 12e_H. Blowup-frame plots; Hansen.

Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 2.87 ALP=40.00 IL=346 IR=461 JT= 67 PO=1.67E+05 HANSEN







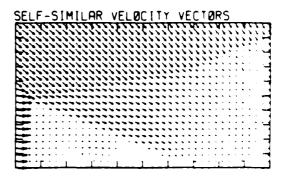
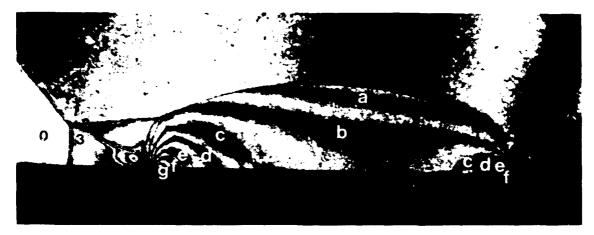


Figure 12e_H. Blowup-frame plots; Hansen - continued.

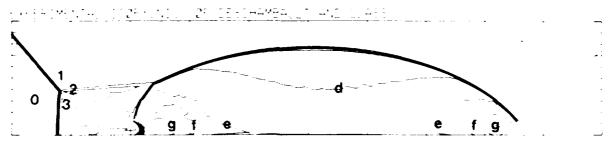
Figure 12. Case 9, M_S = 2.87, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.



Region	و ۱۵ء
0	1.00
1	4.41
2	7.16
3	5.08
a	7.60
ь	8.03
С	8.47
ď	8.91
e	9.35
f	9.78
g	10.22

Figure 13a. Interferogram.

BUILD HE FOREIL DOMESTID NESTOD MEEDS BUILDINGS. TO



XBB 859-7202

Figure 135p., Calculated isopycnics ($\gamma = 1.4$) using the experimental fringes.

Figure 13. Case 10, $M_S = 3.72$, $\theta_W = 40^\circ$, Air, $\gamma = 1.4$ and Hansen EOS, DMR.

MS= 3.72 ALR=40.00 NR=510 NZ=100 MBEG= 90 R0=4.005+14 HANDET

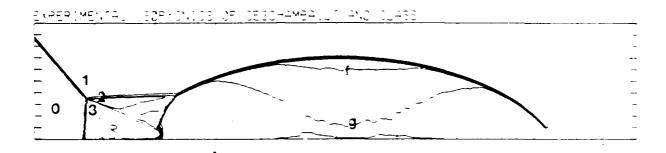


Figure 13b $_{\mathrm{H}}$. Calculated isopycnics (Hansen) using the experimental fringes.

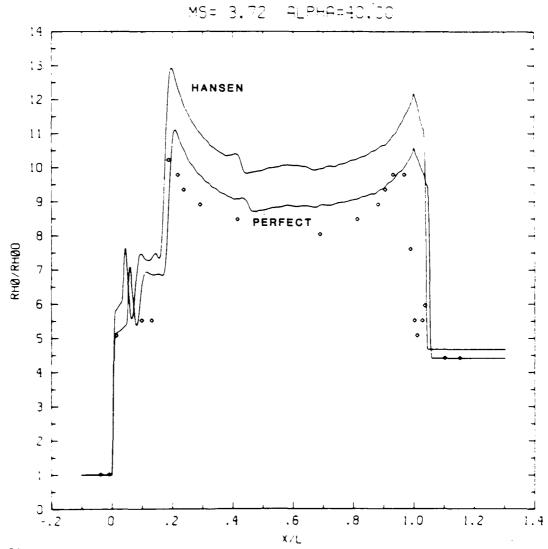


Figure 13c. Wall plot for ρ/ρ_0 , γ = 1.4 and Hansen calculations, with experimental data.

Figure 13. Case 10, $M_s = 3.72$, $9_w = 40^\circ$, Air, y = 1.4 and Hansen EOS, DMR - continued. 143

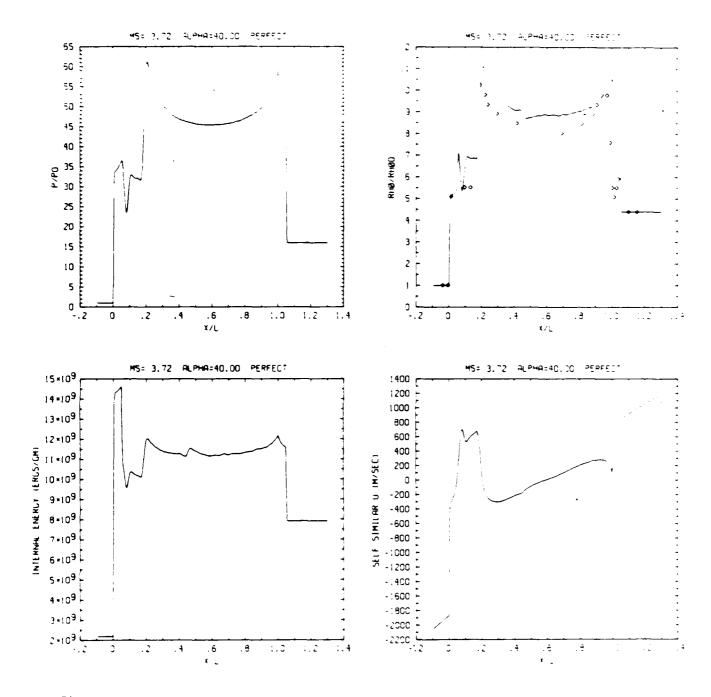


Figure 13c_p. Wall plots for p/p_0 , p/p_0 with experimental data included, e, u; $\gamma = 1.4$.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

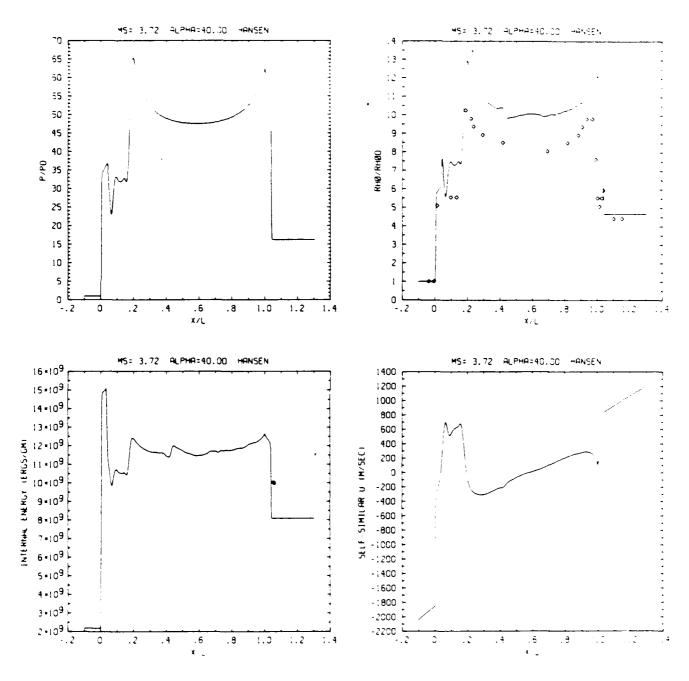
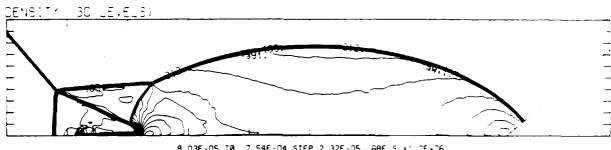


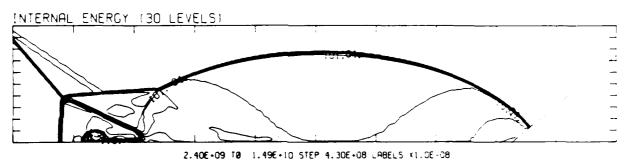
Figure 13c_H. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.70 ALP=40.00 NR=510 NZ=100 ABEG# 30 F0#8.005404 #8#F801



8.03E-05 TØ 7.54E-04 STEP 2.32E-05 LABELS x1.0E+06



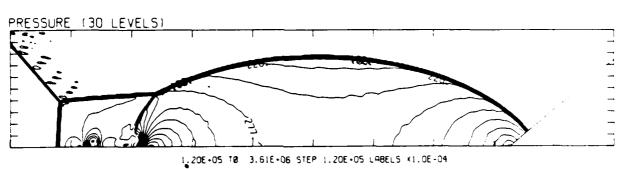
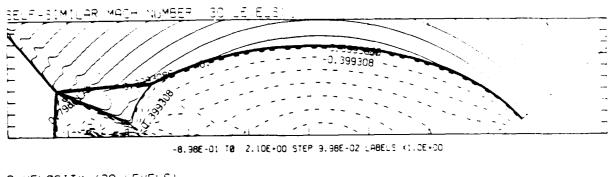
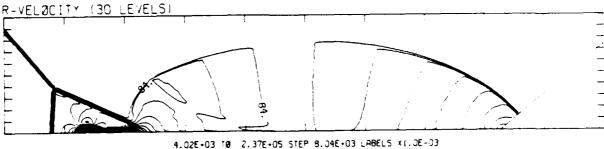


Figure 13d_p. Whole-flowfield contour-plots; $\gamma = 1.4$.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.72 ALP=40.00 NR=510 NZ=100 KBEG= 90 F0=8.00E+04 FEFFE17





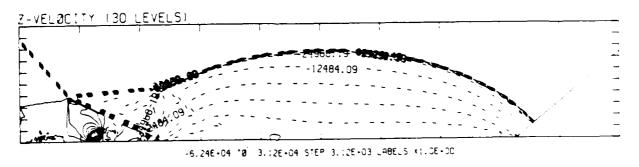
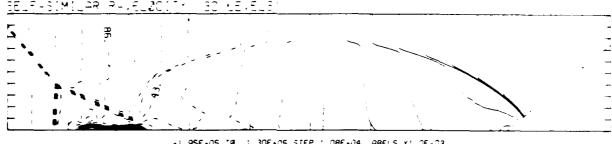


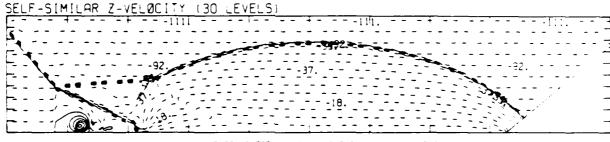
Figure 13d_p. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

1<mark>0 ALP=40.00 NR=5</mark>10 NZ=100 KBE0= 90 F0=8.005+04 P89F5



-1.95E+05 TO 1.30E+05 STEP 1.08E+04 LABELS X1.0E-03



-1.11E+05 TØ 2.79E+04 STEP 4.64E+03 LABELS X1.0E-03

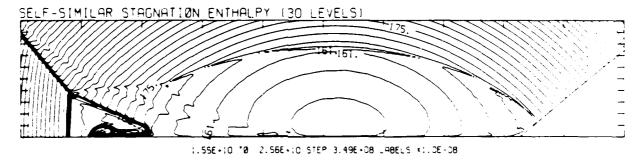
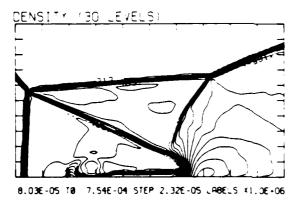
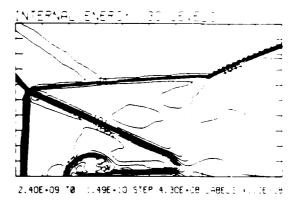


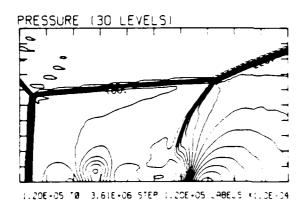
Figure 13dp. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 3.70 ALP=40.00 [L=356 [P=474 UT= 69 P0=6.005+14 PP=F=







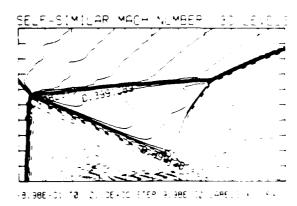
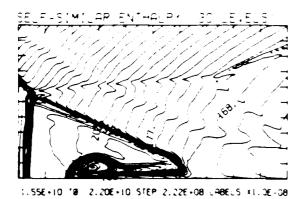
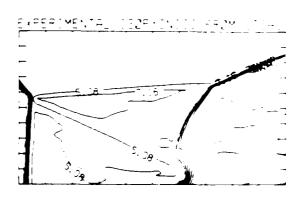


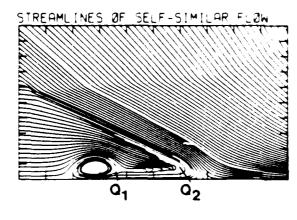
Figure 13ep. Blowup-frame plots; y = 1.4.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MB= 3.70 ALR=40.00 (L=358 (R=474 UT= 89 F0+8.00F+ 4 HHH







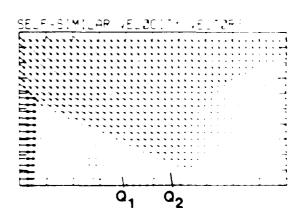


Figure 13ep. Blowup-frame plots; y = 1.4 - continued.

Figure 13. Case 10, M_S = 3.72, 9_M = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MBR 3.70 ALREAD.30 NAMESONER OD WEEDE GOVERNER, PROCESSION

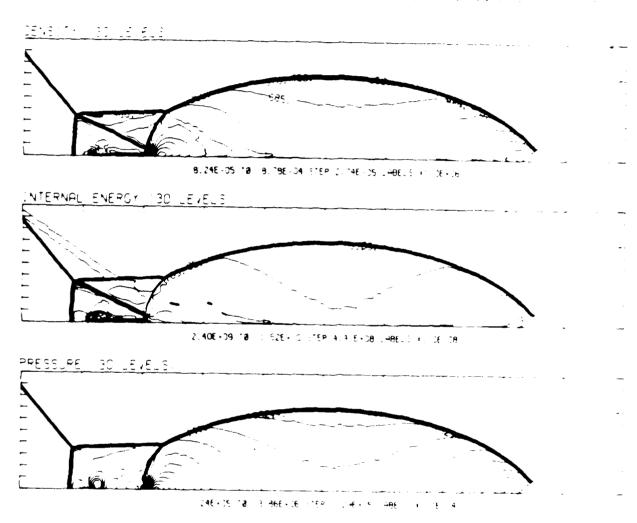


Figure 13d_H. Whole-flowfield contour-plots; Hansen.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

SAL SASSES CHARACTERISTICS

M84 3.70 ALP440.00 NR4510 NZ4100 KBE04 30 F048.018404 HAM

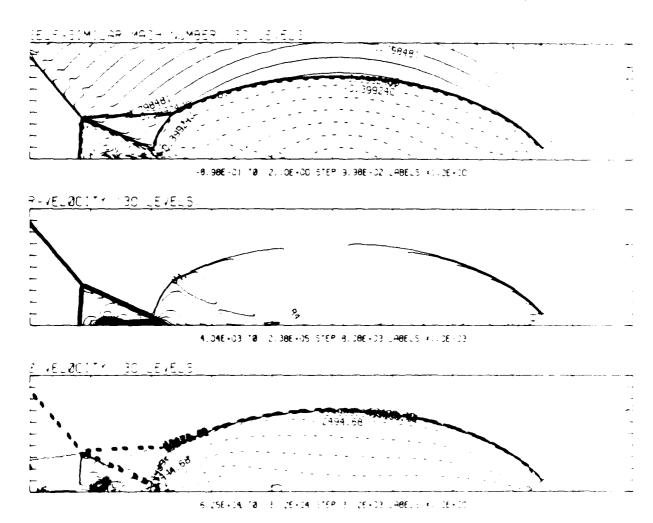
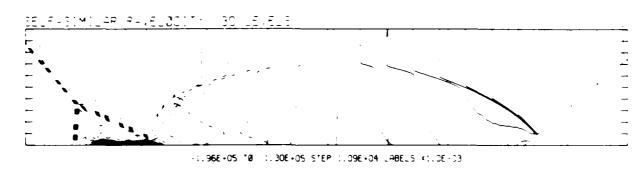
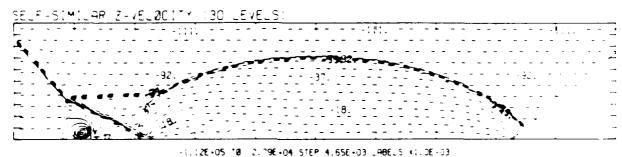


Figure 13dy. Whole-flowfield contour-plots; Hansen - continued.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EDS, DMR - continued.

MS= 3.70 ALP=40.00 NR=510 NZ=100 XBE0= 90 F0=8.005+04 HANSEN





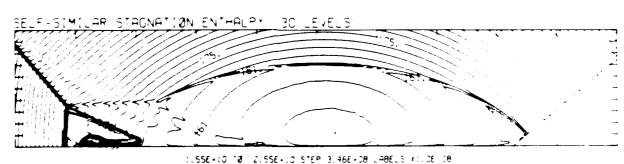
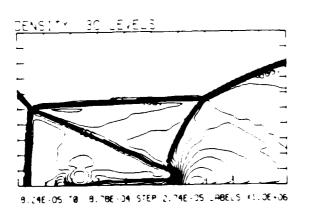
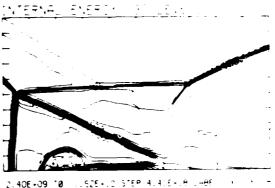


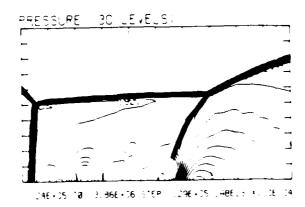
Figure 13d $_{\mathrm{H}}$. Whole-flowfield contour-plots; Hansen - continued.

Figure 13. Case 10, M_S = 3.72, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MSE 3.72 ALREAD.00 (14354 194470 UTA 89 FORE.005+14 M/V







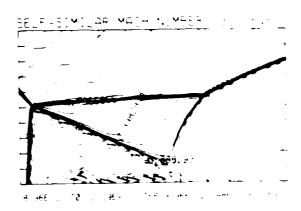
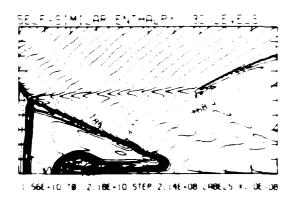
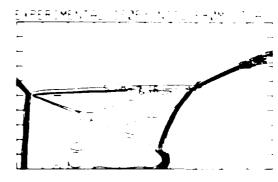


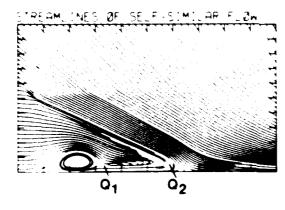
Figure 13e_H. Blowup-frame plots; Hansen.

Figure 13. Case 10, $M_s = 3.72$, $\theta_w = 40^\circ$, Air, $\gamma = 1.4$ and Hansen EOS, DMR - continued.

MS= 3.72 ALP=40.00 [L=354 [R=472 LT= 69 F0=6.006+0, HAY/F/







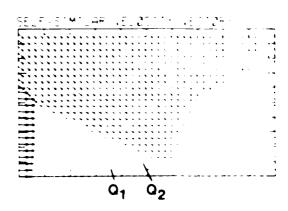


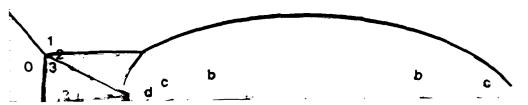
Figure 13e4. Blowup-frame plots; Hansen - continued.

Figure 13. Case 10, $M_S = 3.72$, $\rho_W = 40^\circ$, Air, $\gamma = 1.4$ and Hansen EOS, DMR - continued.



Region	5/50
0	1.00
1	4.86
2	7.90
3	5.37
a	9.78
ь	10.72
С	11.67
d	12.61
e	8.84
f	7.90
g	6.95
h	6.01
i	5.07
j	4.13

Figure 14a. Interferogram.



XBB 859-7203

Figure 14bp. Calculated isopycnics (v=1.4) using the experimental fringes.

Figure 14. Case 11, $M_s = 4.62$, $n_w = 400$, Air, v = 1.4 and Hansen 600, DMR.

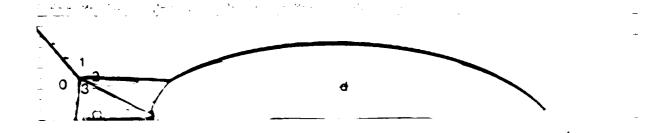


Figure 14: 1 in a lifeter of the control of the equation of the equation of the equation of the engage.

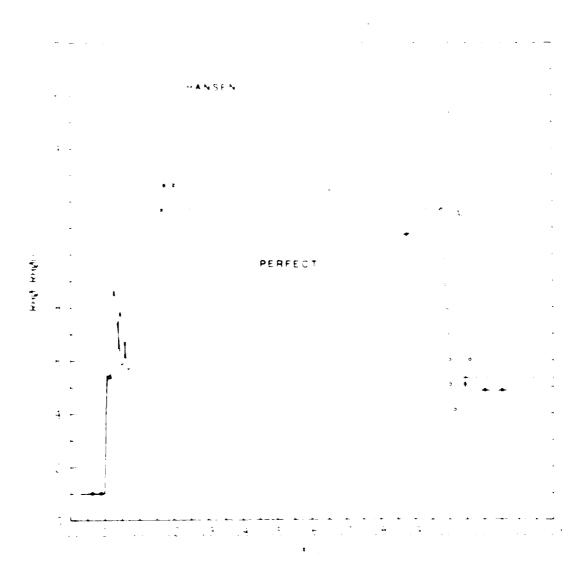


Figure 14c. Wall plot for ρ/ρ_0 , $\nu=1.4$ and Hansen calculations, with experiment data.

Figure 14. Case 11, $M_s = 4.62$, $A_s = 40^\circ$, Air, v = 1.4 and Hansen FM. DMR = continued.

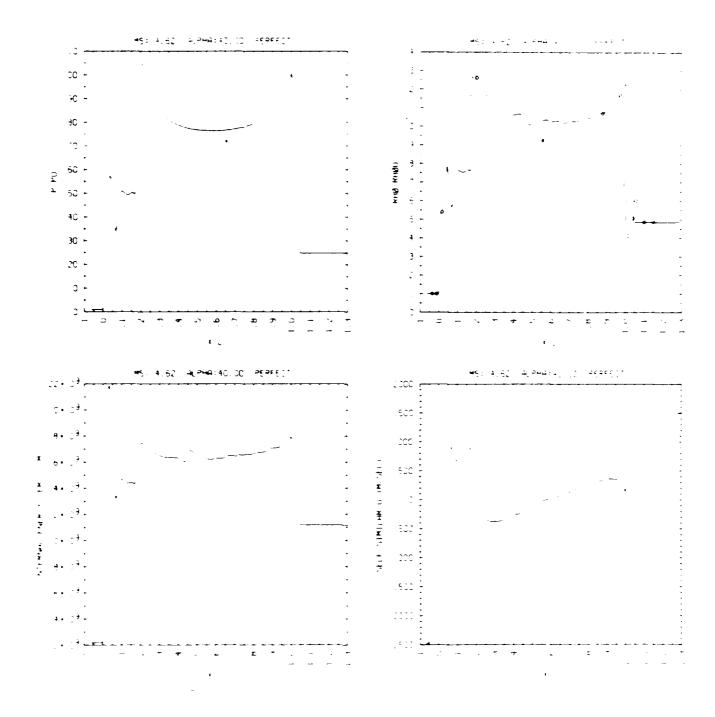


Figure 14c_p. Wall plots for p/p_0 , p/p_0 with experimental data included, e, u; $\gamma = 1.4$.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

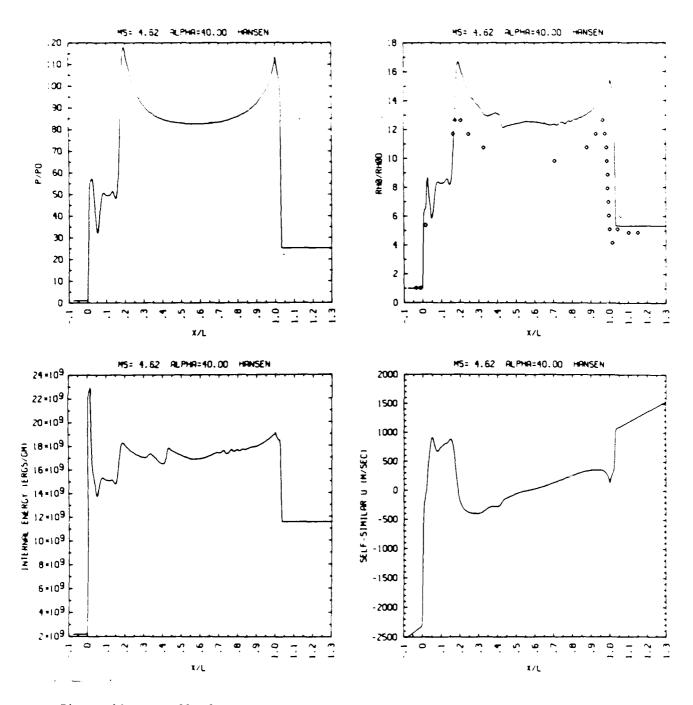


Figure 14c_H. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, u; Hansen.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 4.62 ALP=40.00 NR=510 NZ= 30 48E3= 30 F0=1.80E+34 FEFFE37

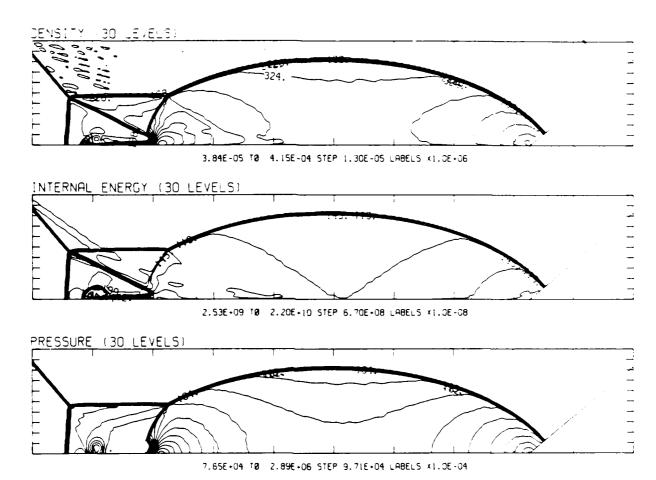
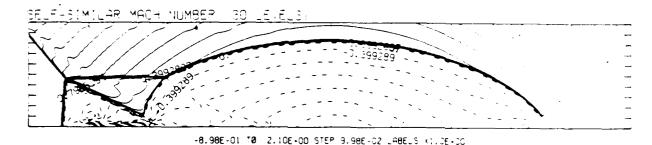
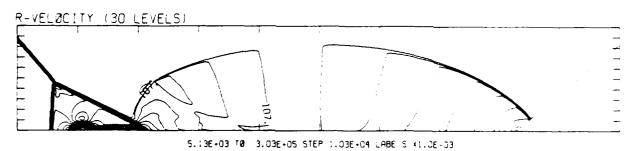


Figure $14d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$.

Figure 14. Case 11, $M_s = 4.62$, $\theta_w = 40^\circ$, Air, $\gamma = 1.4$ and Hansen EOS, DMR - continued.

MS= 4.82 ALP=40.00 NP=510 NZ= 80 MEE3= 80 PD=0.808-14 REFFED





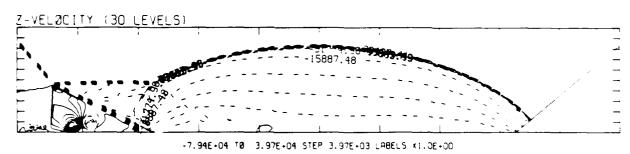
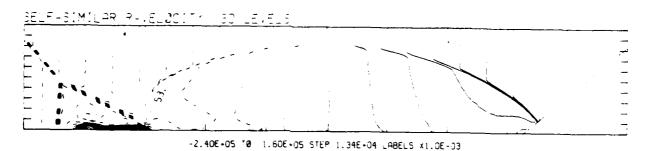
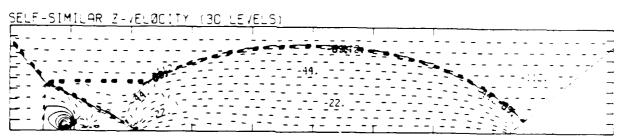


Figure $14d_p$. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

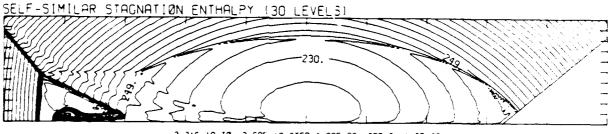
Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 4.62 ALP=40.00 NR=510 NZ= 30 48E0= 30 PO=2.80E+04 REFREDI





-1.29E+05 T0 3.93E+04 STEP 5.61E+03 LABELS X1.0E-03



2.21E+10 T0 3.60E+10 STEP 4.80E+08 LABELS X1.3E-08

Figure 14d_p. Whole-flowfield contour-plots; $\gamma = 1.4$ - continued.

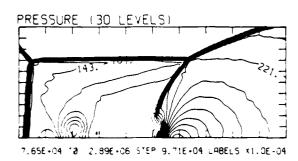
Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 4.60 ALP=40.00 (L=344 (R=485 UT= 80 PO=2.805-14 Againg)



2.53E+09 TØ 2.20E+10 STEP 5.70E+08 L-86LS (1.0E+08

3.84E-05 TØ 4.15E-04 STEP 1.30E-05 LABELS X1.0E+06



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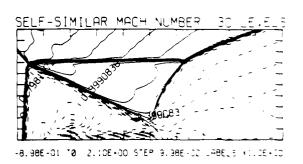
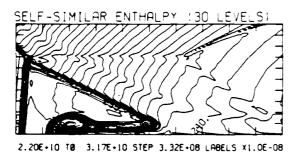
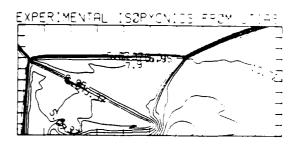


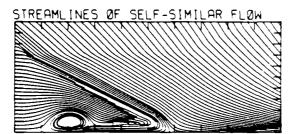
Figure 14e_p. Blowup-frame plots; $\gamma = 1.4$.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 4.60 ALP=40.00 (L=344 [P=485 JT= 60 PD=0.805+j4 P5445])







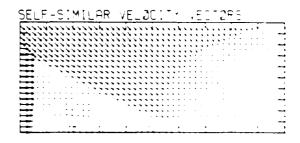


Figure 14e_p. Blowup-frame plots; $\gamma = 1.4$ - continued.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

48 | | 4.82 | ALF#40.20 | NA#8.3 | NZ# | 90 | 4880# | 90 | 60#3.438# | 4.60

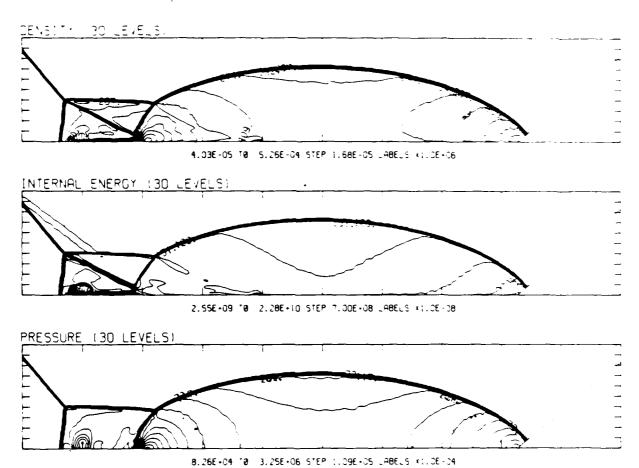
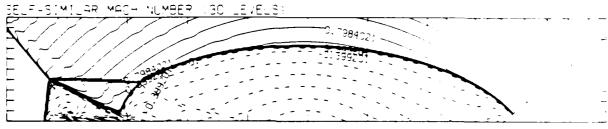


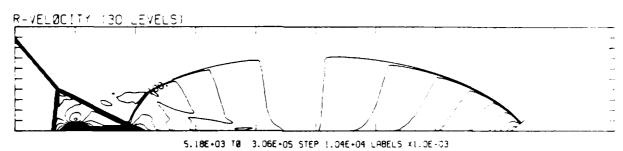
Figure 14d_H. Whole-flowfield contour-plots; Hansen.

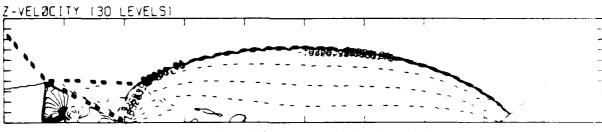
Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

MS= 4.60 ALF=40.00 MR=510 ME= 90 MBE0= 90 F0=0.800.800.800



-8.98E-01 10 2.10E+00 STEP 9.98E-02 LABELS 41.0E+00



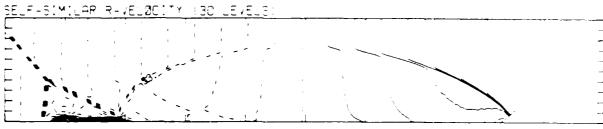


-8.:6E+04 T0 4.38E+04 STEP 4.08E+03 LABELS X1.3E+00

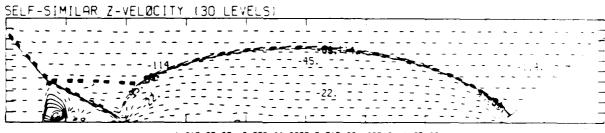
Figure $14d_{\mathrm{H}}$. Whole-flowfield contour-plots; Hansen - continued.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

M8= 4.82 FLF=40.00 NR=810 NZ= 90 M8E0= 90 F0=1.8054 & HAVE



-2.42E+05 T0 1.61E+05 STEP 1.34E+04 LABELS <1.0E+03



-1.31E+05 18 3.99E+04 STEP 5.71E+03 LABELS X1.0E-03

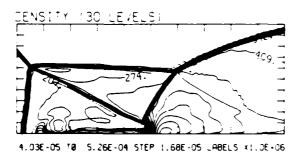


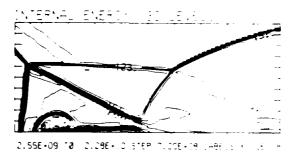
2.23E+10 10 3.67E+10 STEP 4.97E+08 LABELS X1.0E-08

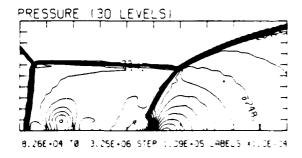
Figure $14d_{H}$. Whole-flowfield contour-plots; Hansen - continued.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

M59 4.62 ALP#40.00 (L#34) (F#46) (T# 59 8) 4(.5) 4 4







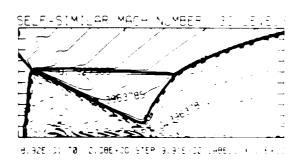
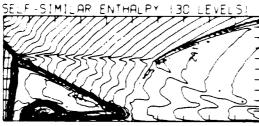


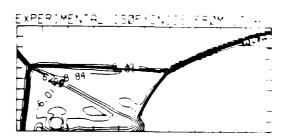
Figure 14e_H. Blowup-frame plots; Hansen.

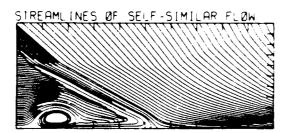
Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.

M3= 4.60 ALP=40.00 IL=341 IR=481 UT= 59 F0=0.805+34 H4K/+>



2.22E+10 18 3.16E+10 STEP 3.24E+08 LABELS X1.0E-08





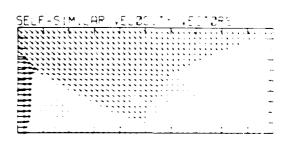
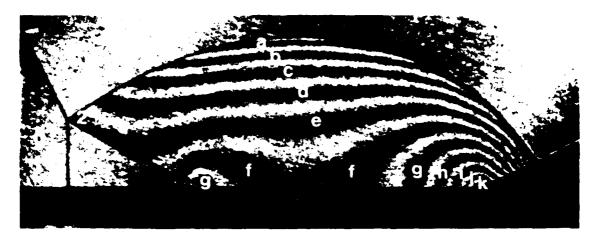
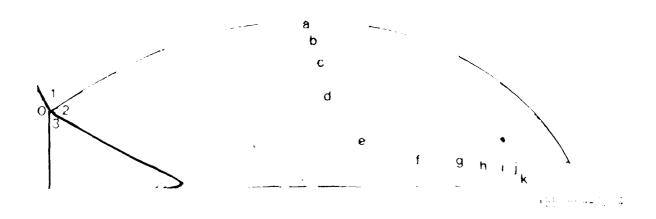


Figure 14e_H. Blowup-frame plots; Hansen - continued.

Figure 14. Case 11, M_S = 4.62, θ_W = 40°, Air, γ = 1.4 and Hansen EOS, DMR - continued.



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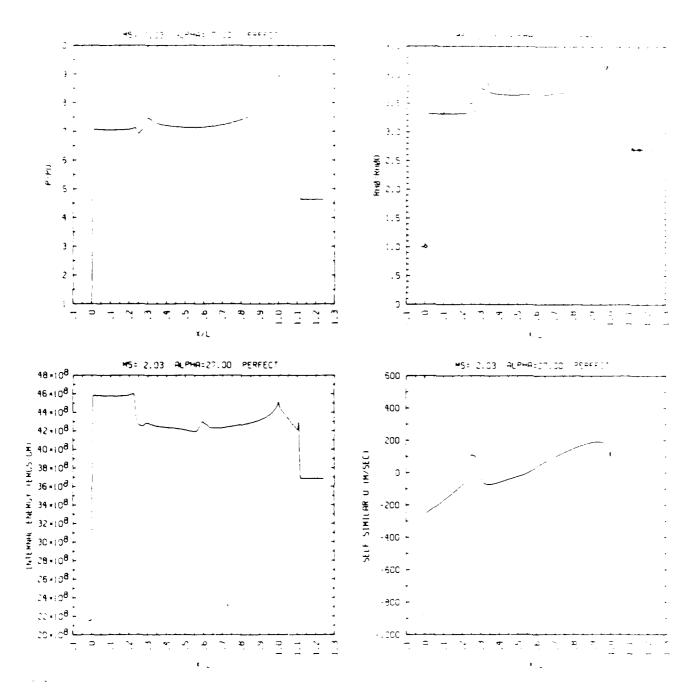


Figure 15c. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \bar{u} .

Figure 15. Case 12, M_S = 2.03, θ_W = 27°, Air, γ = 1.4, SMR - continued.

MOE 2.03 ALRELT.00 (ARESE NZELED HABOVIT ELL.)

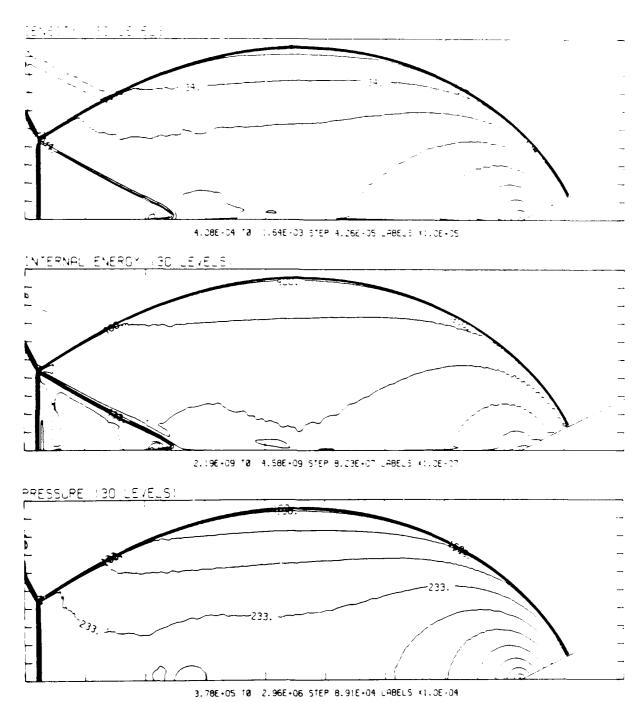
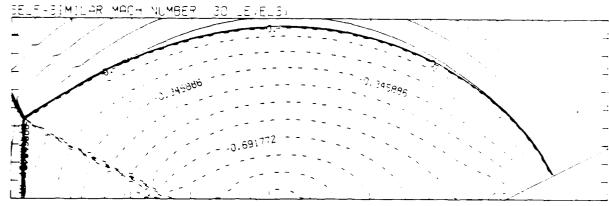


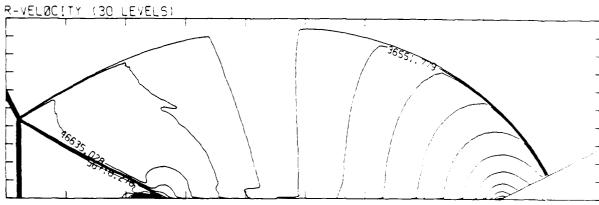
Figure 15d. Whole-flowfield contour-plots.

Figure 15. Case 12, M_S = 2.03, θ_W = 27°, Air, γ = 1.4, SMR - continued.

MB# 0.08 ALF#07.00 NF#405 NZ#180 KBE3# 75 F0#8.88E406 RE##807



-9.51E-01 70 1.64E+00 STEP 8.65E-02 LABELS *1.0E+00



1.26E+03 TØ 7.44E+04 STEP 2.52E+03 LABELS X1.3E+00

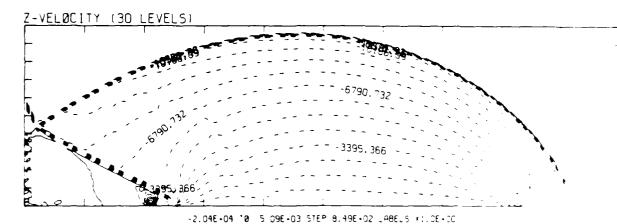
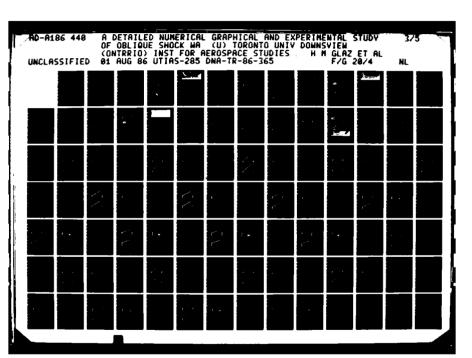
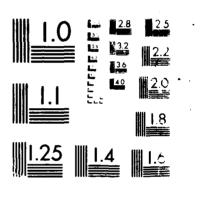


Figure 15d. Whole-flowfield contour-plots - continue:.

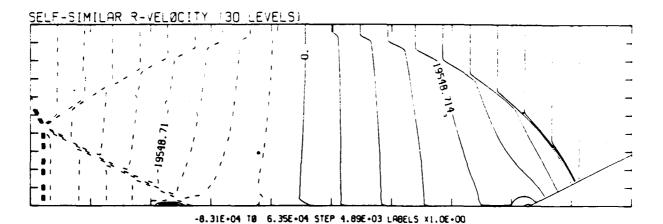
Figure 15. Case 12, $M_S = 2.03$, $\theta_W = 27^\circ$, Air, $\gamma = 1.4$, $M_S = 1.4$





MICROCHER RESOLUTION TEST FHAT

MS= 2.03 ALP=27.00 NR=425 NZ=130 KBEG= 75 PO=3.33E+05 PEPFEDT



SELF-SIMILAR Z-VELØCITY (3C LEVELS)

3772-8036

-30218.24

-30218.24

-22663.68

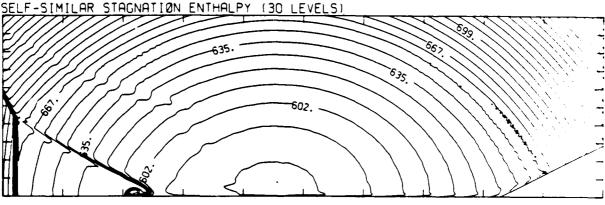
-22663.68

-15109.12

-30218.24

-7554.561



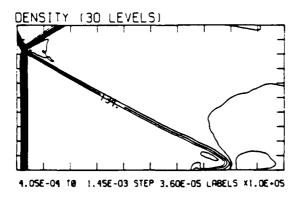


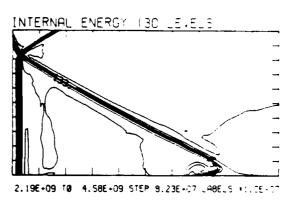
5.87E+09 TØ 8.21E+09 STEP 8.08E+07 LABELS x1.0E-07

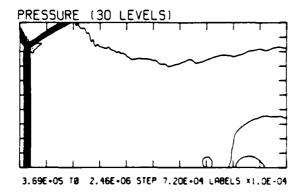
Figure 15d. Whole-flowfield contour-plots - continued.

Figure 15. Case 12, M_S = 2.03, θ_W = 27°, Air, γ = 1.4, SMR - continued.

MS= 2.03 ALP=27.00 [L=295 [R=419 JT= 69 PC=3.33E+05 FEFFE]]







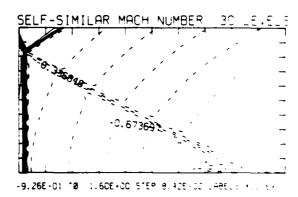
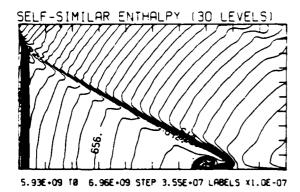
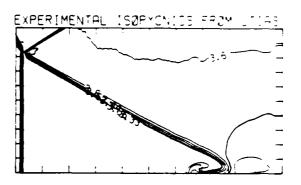


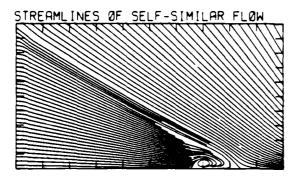
Figure 15e. Blowup-frame plots.

Figure 15. Case 12, $M_S = 2.03$, $\theta_W = 27^{\circ}$, Air, $\gamma = 1.4$, SMR.

MS= 2.03 ALP=27.00 [L=295 [R=419 UT= 69 PC=3.33E+35 PERFE]







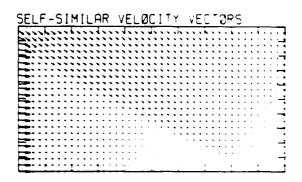


Figure 15e. Blowup-frame plots - continued.

Figure 15. Case 12, M_S = 2.03, θ_W = 27°, Air, γ = 1.4, SMR.



Region	2/2;	Region	z/z:
0	1.00	g	13.32
1	5.63	ħ	13.95
1'	6.89	i	14.58
2	7.44	j	15.21
3	5.74	k	15.84
a	9,53	1	16.47
5	10.16	31	17.10
С	10.79	n	17.73
d	11,42	0	18.36
e	12.05	p	6.37
Í	12.68	Р	8.07

Figure 16a. Interferogram.

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XBB 859-7205

Figure 16b. Calculated isopycnics using the experimental fringes.

Figure 16. Case 13, Mg = 3.70, Ag = 270, Air, Mansen, EOS, OMR.

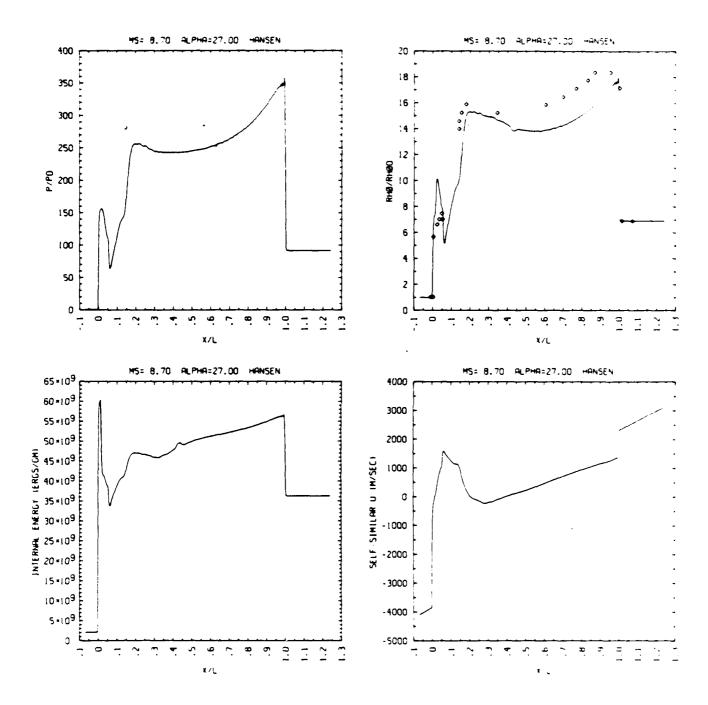


Figure 16c. Wall plots for p/p_0 , p/ρ_0 with experimental data included, e, \bar{u} .

Figure 16. Case 13, M_S = 8.70, θ_W = 27°, Air, Hansen, EOS, DMR - continued.

MS= 8.70 ALP=27.00 NR=530 NZ= 85 KBES= 30 PO=4.10E-04 H4M3EM

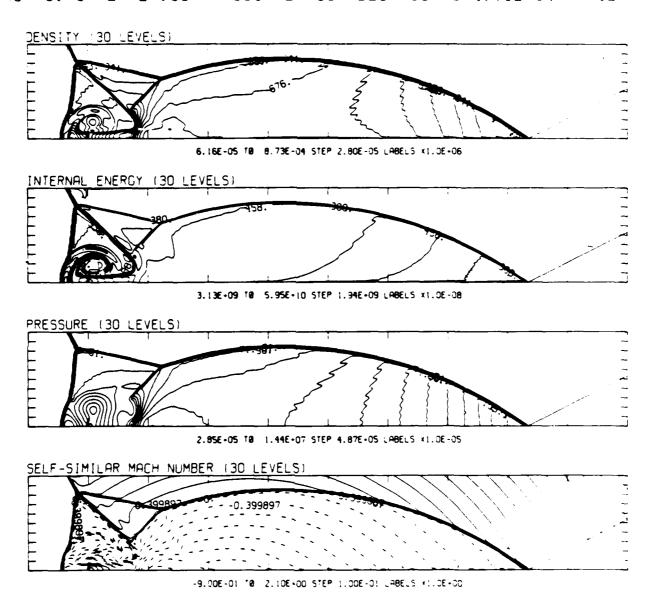


Figure 16d. Whole-flowfield contour-plots.

Figure 16. Case 13, M_S = 8.70, θ_W = 27°, Air, Hansen, EOS, DMR - continued.

MS= 8.70 ALP=27.00 NR=530 NZ= 85 KBEG= 80 PC=4.10E+04 H4NSEN

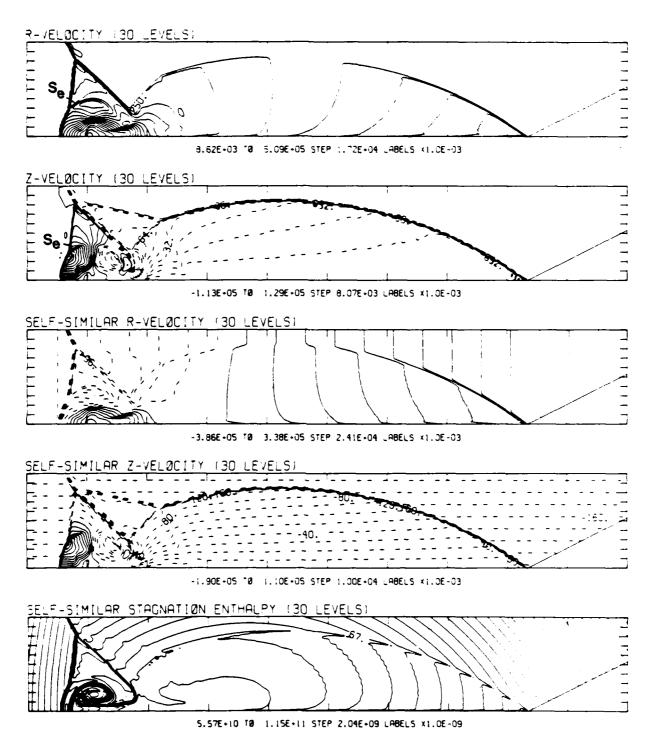
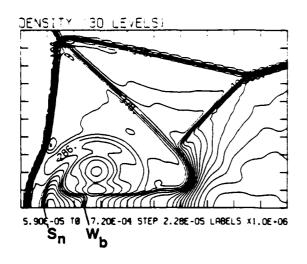
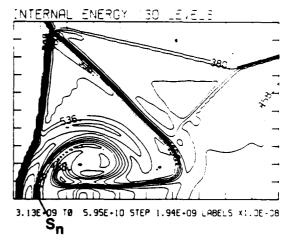


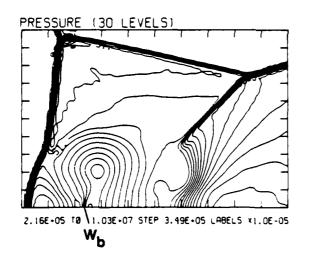
Figure 16d. Whole-flowfield contour-plots - continued.

Figure 16. Case 13, M_c = 8.70, θ_u = 27°, Air, Hansen, EOS, DMR - continued.

MS= 8.70 ALP=27.30 IL=396 [R=503 UT= T3 F0=4.10E+34 HANSEN







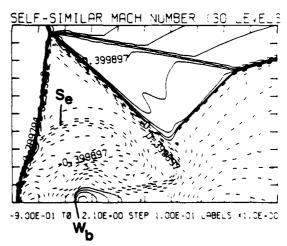
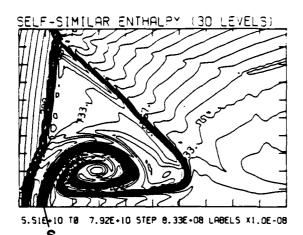
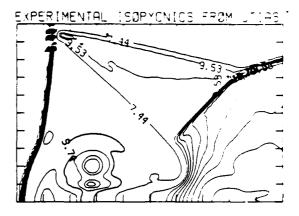


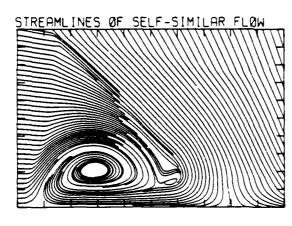
Figure 16e. Blowup-frame plots.

Figure 16. Case 13, M_S = 8.70, θ_W = 27°, Air, Hansen, EOS, DMR - continued.

MS= 8.70 ALP=27.00 IL=396 IR=503 UT= 73 PO=4.10E+04 HANSE'.







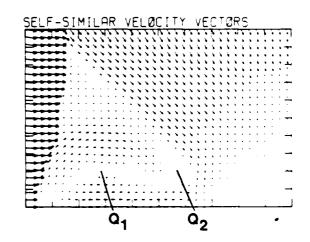


Figure 16e. Blowup-frame plots - continued.

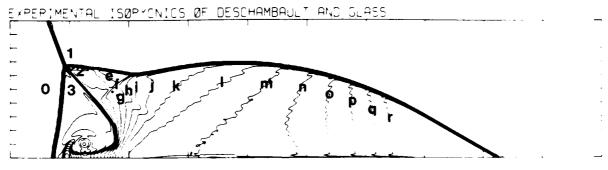
Figure 16. Case 13, M_S = 8.70, θ_W = 27°, Air, Hansen, EOS, DMR - continued.



Region	٥/٥	Region	۵/۵
0	1.00	h	9.24
1	5.47	i	9.56
1'	6.13	j	9.88
2	6.65	k	10.21
3	5.59	1	10.53
a	6.97	m	10.85
Ъ	7.30	n	11.18
С	7.62	0	11.50
d	7.94	p	11.82
e	8.27	q	12.15
f	8.59	r	12.47
g	8.91	s	5.91

Figure 17a. Interferogram.

MS= 7.19 ALP=20.00 NP=510 NZ=120 KBEG= 90 F0=8.018+ 1 Heb. F0



XBB 859-7206

Figure 17b. Calculated isopycnics using the experimental fringes.

Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR.

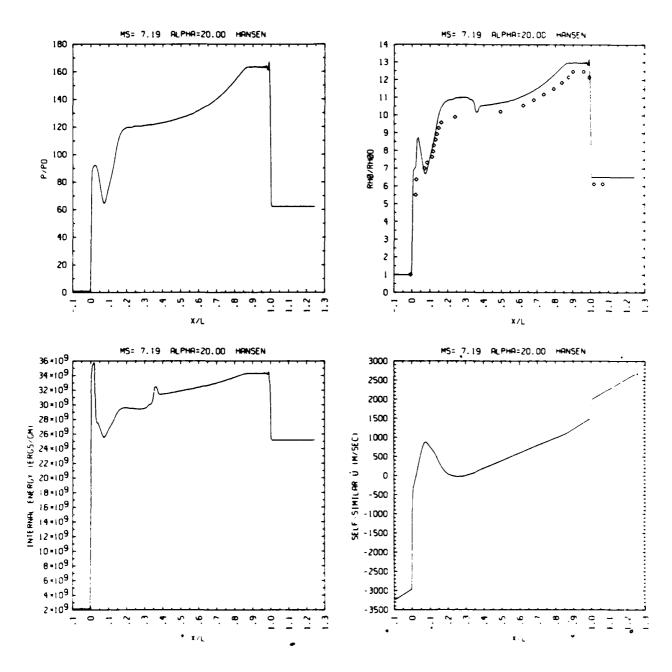


Figure 17c. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \bar{u}_*

Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR - continued.

MS= 7.19 ALP=20.00 NR=510 NZ=120 KBEG= 90 PC=8.00E+04 H4NSEN

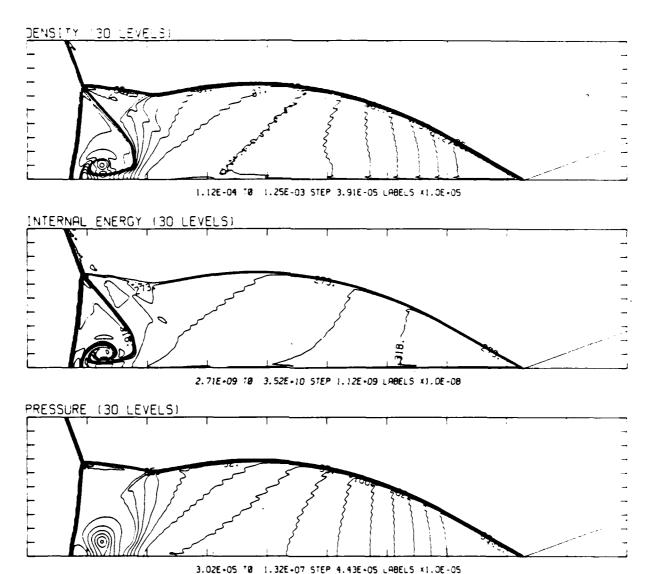
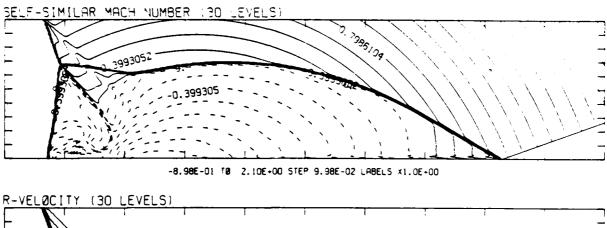
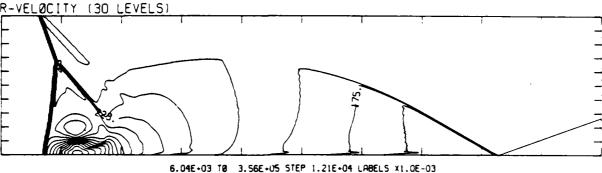


Figure 17d. Whole-flowfield contour-plots.

Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR - continued.

MS= 7.19 ALP=20.00 NR=510 NZ=120 KBEG= 90 PC=8.00E+04 HANSEN





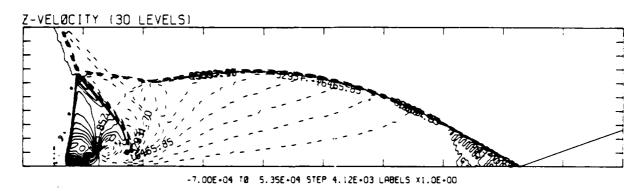
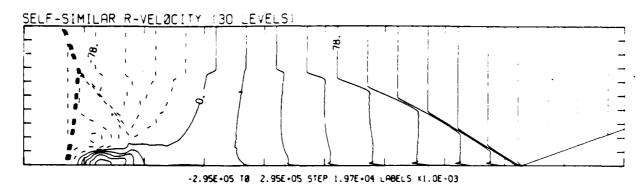
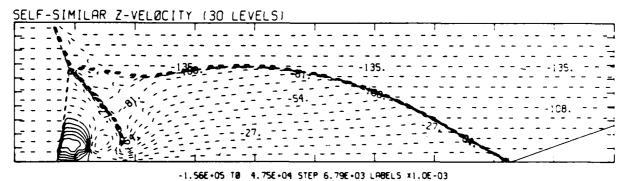


Figure 17d. Whole-flowfield contour-plots - continued.

Figure 17. Case 14, $\rm M_S$ = 7.19, $\rm \theta_W$ = 20°, Air, Hansen EOS, C/DMR - continued.

MS= 7.19 ALP=20.00 NR=510 NZ=120 KBEG= 90 PD=8.00E+04 HAMSEN





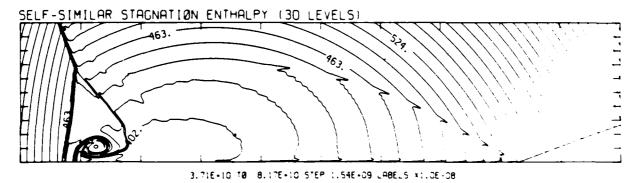


Figure 17d. Whole-flowfield contour-plots - continued.

Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR - continued.

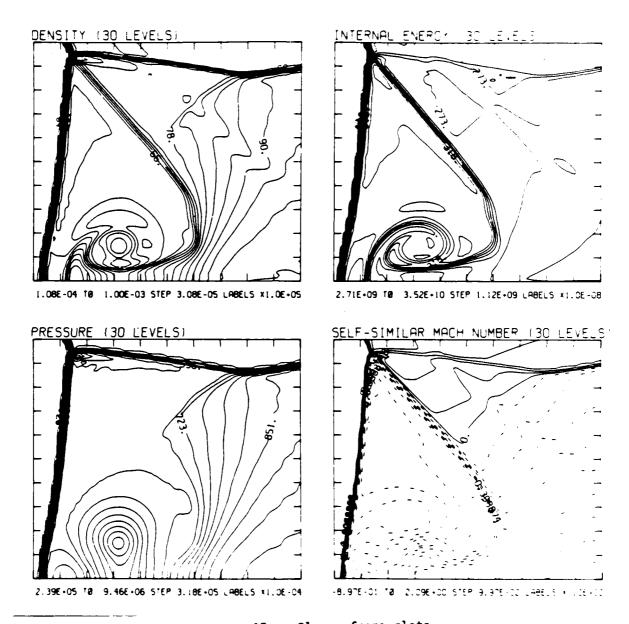
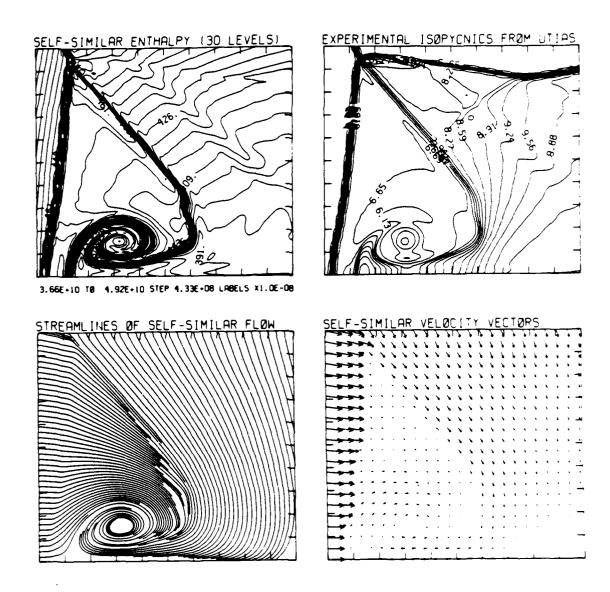


Figure 17e. Blowup-frame plots.

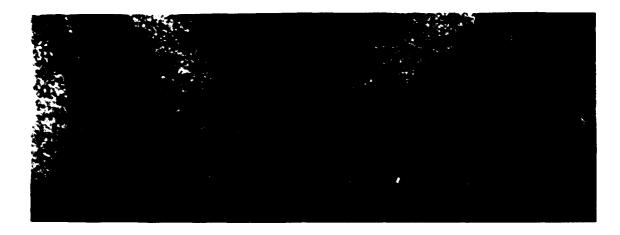
Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR - continued.

MS= 7.19 ALP=20.00 IL=383 IR=476 JT= 85 PO=8.00E+04 HANSEN



· Figure 17e. Blowup-frame plots - continued.

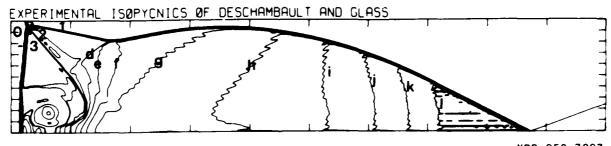
Figure 17. Case 14, M_S = 7.19, θ_W = 20°, Air, Hansen EOS, C/DMR - continued.



Region	۵/۵	Region	۵/۵
0	1.00	f	10.73
1	5.64	g	11.37
1'	6.93	h	12.02
2	6.85	i	12.66
3	5.72	j	13.31
а	7.50	k	13.95
Ь	8.14	1	14.60
С	8.79	m	15.25
d	9.43	n	15.89
e	10.08	p	6.85

Figure 18a. Interferogram.

MS= 8.86 ALP=20.00 NR=575 NZ=110 KBEG= 75 PO=4.10E+04 HANSEN



XBB 859-7207

Figure 18b. Calculated isopycnics using the experimental fringes.

Figure 18. Case 15, M_S = 8.86, A_W = 20°, Air, Hansen EOS, DMR.

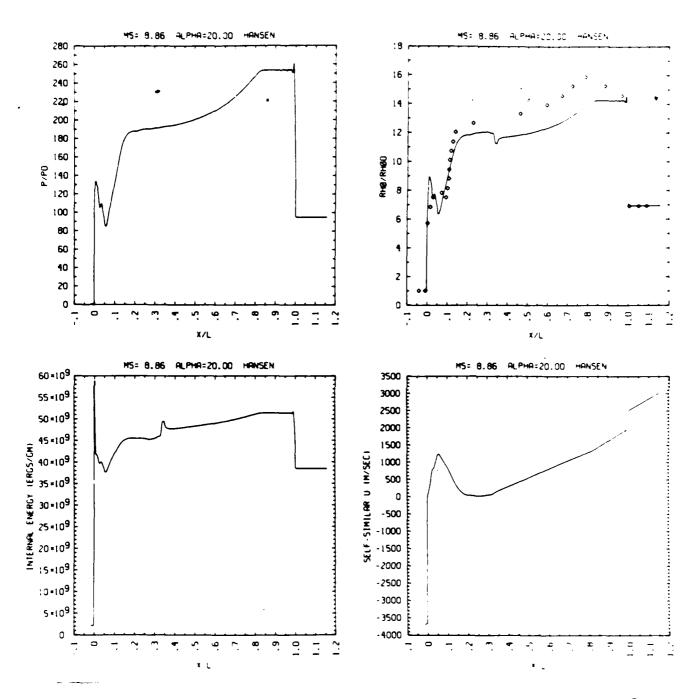


Figure 18c. Wall plots for p/p_0 , ρ/ρ_0 with experimental data included, e, \bar{u} .

Figure 18. Case 15, M_S = 8.86, θ_W = 20°, Air, Hansen EOS, DMR - continued.

MS= 8.86 ALP=20.00 NR=575 NZ=110 KBEG= 75 PO=4.13E+34 H44.8E4.

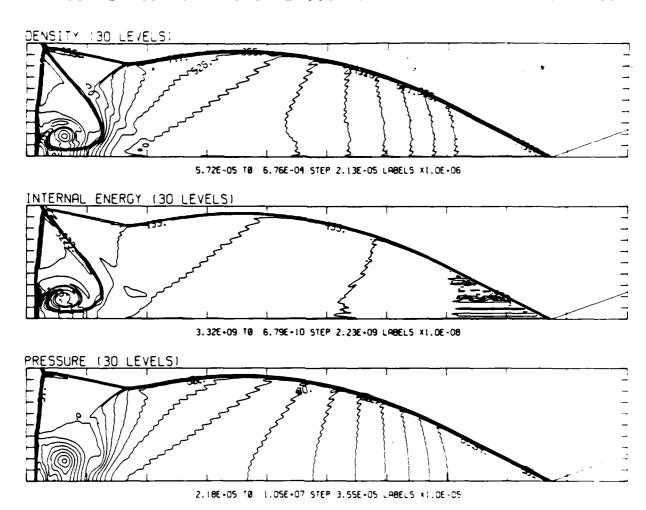


Figure 18d. Whole-flowfield contour-plots.

Figure 18. Case 15, M_S = 8.86, θ_W = 20°, Air, Hansen EOS, DMR - continued.

MS= 8.86 ALP=20.00 NR=575 NZ=110 KBEG= 75 PD=4.105+14 HAY.EEY.

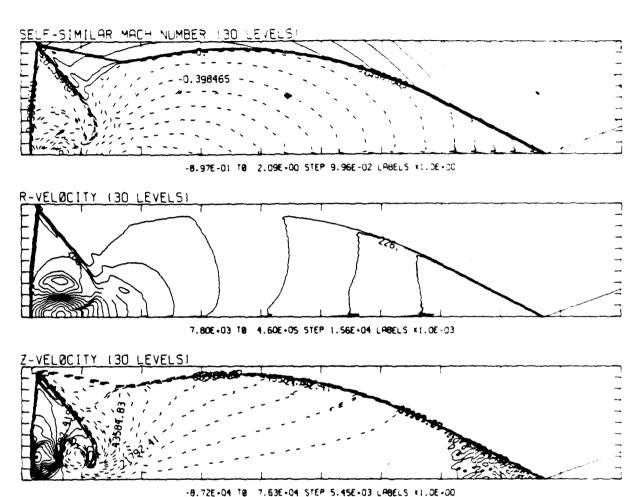
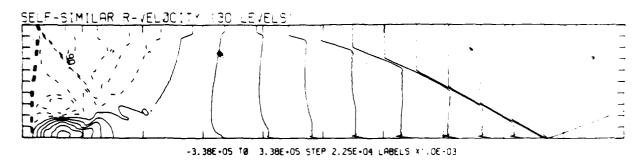
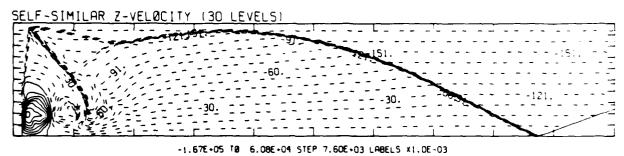


Figure 18d. Whole-flowfield contour-plots - continued.

Figure 18. Case 15, M_S = 8.86, q_W = 20°, Air, Hansen EOS, DMR - continued.

MS= 8.86 ALP=20.00 NR=575 NZ=110 KBEG= 75 PC=4.10E+04 HANSEN





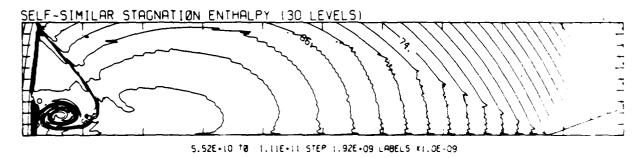


Figure 18d. Whole-flowfield contour-plots - continued.

Figure 18. Case 15, M_S = 8.86, θ_W = 20°, Air, Hansen EOS, DMR - continued.

MS= 8.86 ALP=20.00 IL=466 IR=568 JT= 99 PO=4.10E+04 HANSEN

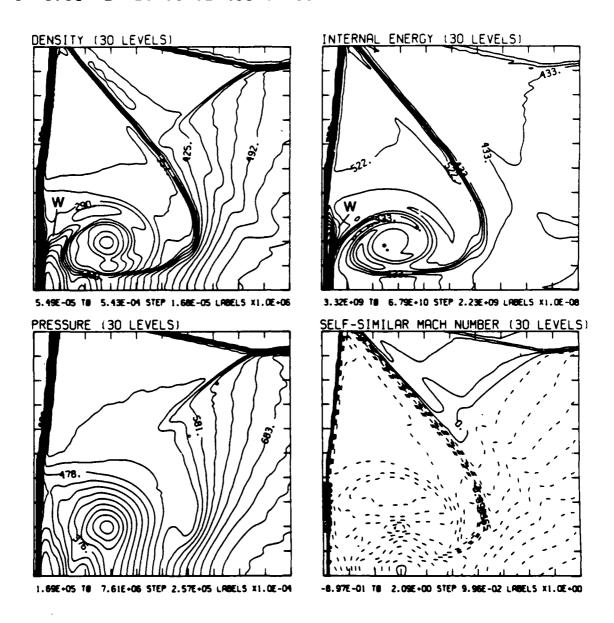


Figure 18e. Blowup-frame plots.

Figure 18. Case 15, M_S = 8.86, θ_W = 20°, Air, Hansen EOS, DMR - continued.

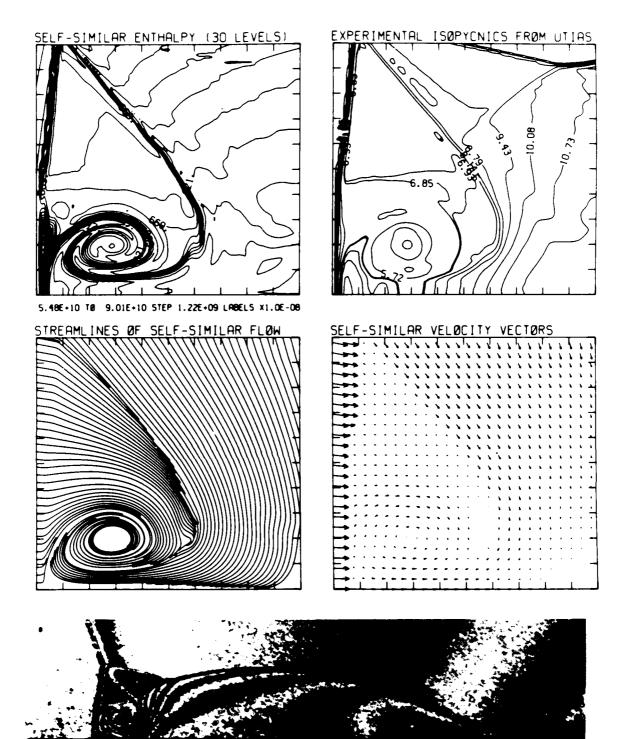
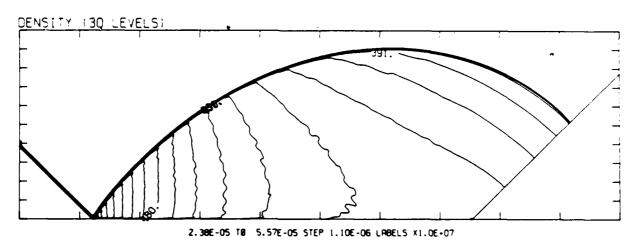


Figure 18f. Reproduction of the interferogram of Exp. 974, Ref. [14]; $M_S = 10.18$, $\theta_W = 20^\circ$, Air. XBB 859-7208

Figure 18. Case 15, M_s = 8.86, θ_w = 20°, Air, Hansen EOS, DMR - continued.

MS= 1.30 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+34 PERFECT



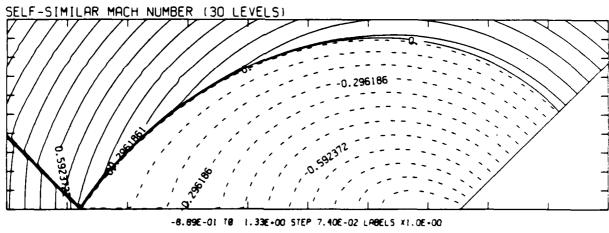


Figure 19.1a. $M_s = 1.30$, whole-flowfield contour plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4.

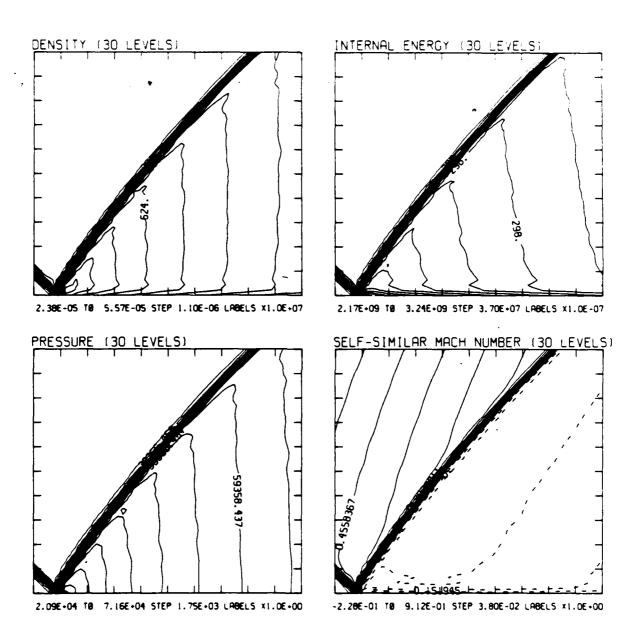


Figure 19.1b. $M_s = 1.30$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

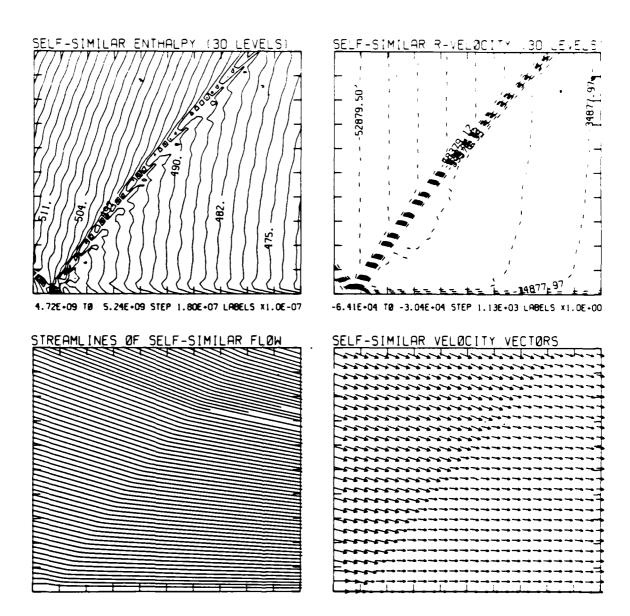
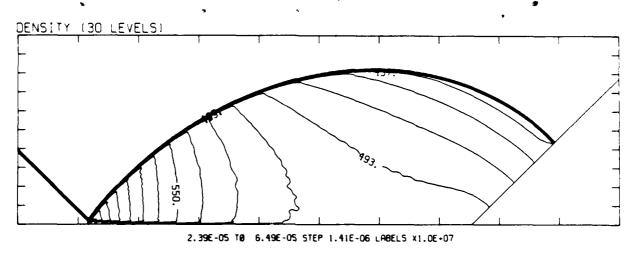


Figure 19.1b. $M_S = 1.30$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 1.40 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+04 PERFECT



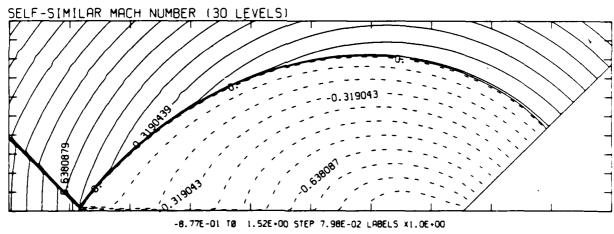


Figure 19.2a. $M_s = 1.40$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 1.40 ALP=45.00 IL=393 IR=444 UT= 48 PC=2.30E+64 PERFECT

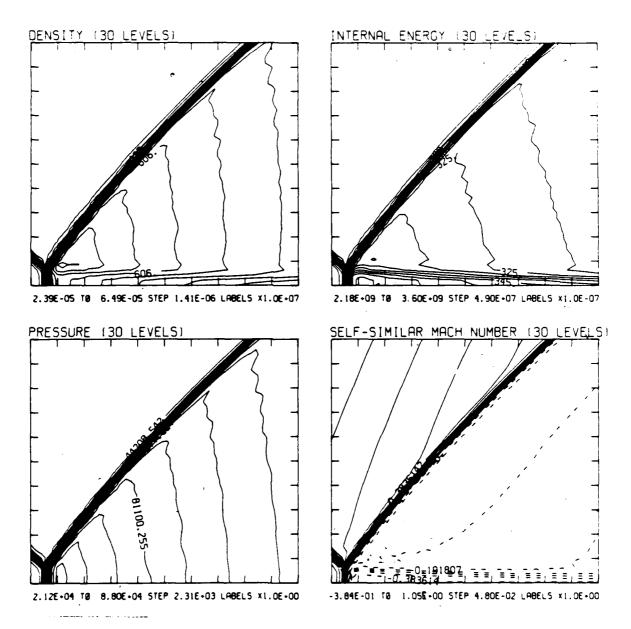


Figure 19.2b. $M_s = 1.40$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.40 ALP=45.00 IL=393 IR=444 JT= 48 PO=2.00E+04 PERFECT

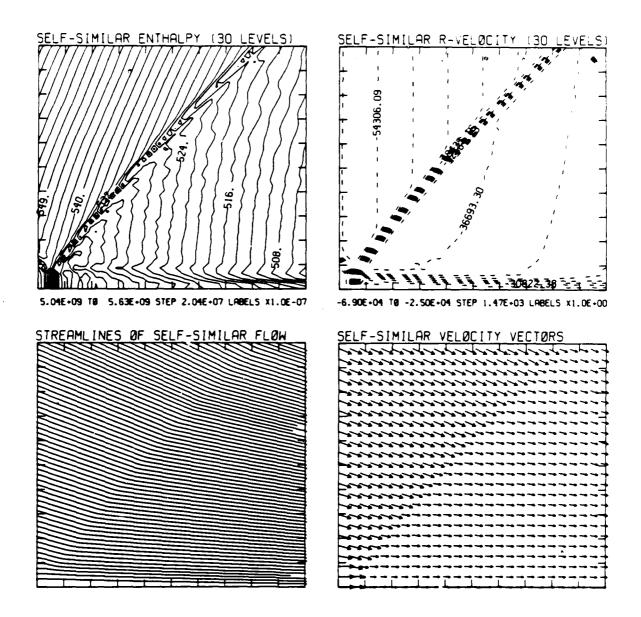
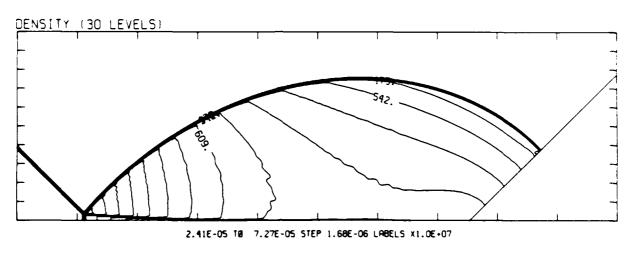


Figure 19.2b. $M_s = 1.40$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 1.50 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 PERFECT



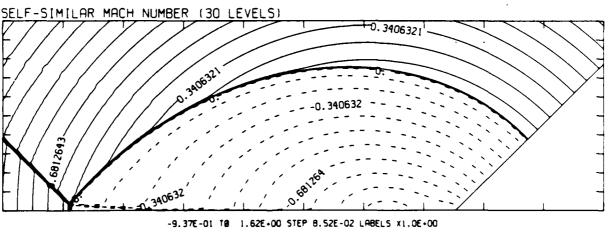


Figure 19.3a. $M_s = 1.50$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.50 ALP=45.00 IL=395 IR=447 UT= 49 PC=2.00E+04 PERFECT

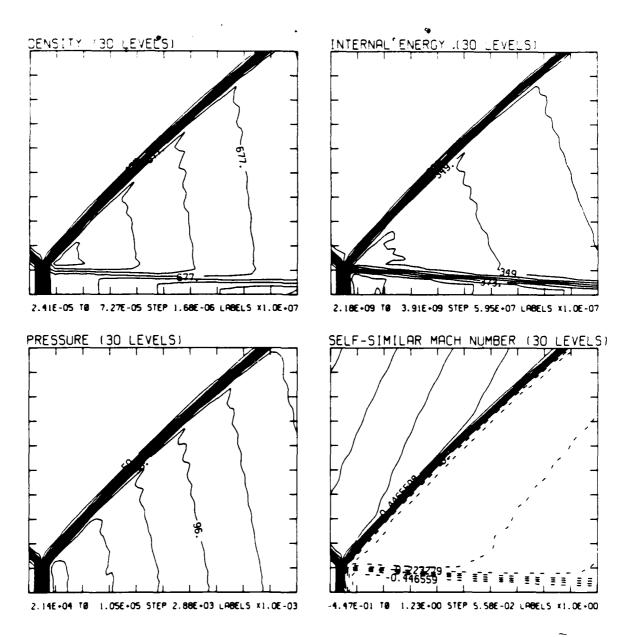


Figure 19.3b. $M_s = 1.50$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.50 ALP=45.00 IL=395 IR=447 UT= 49 PC=2.00E+34 PERFECT

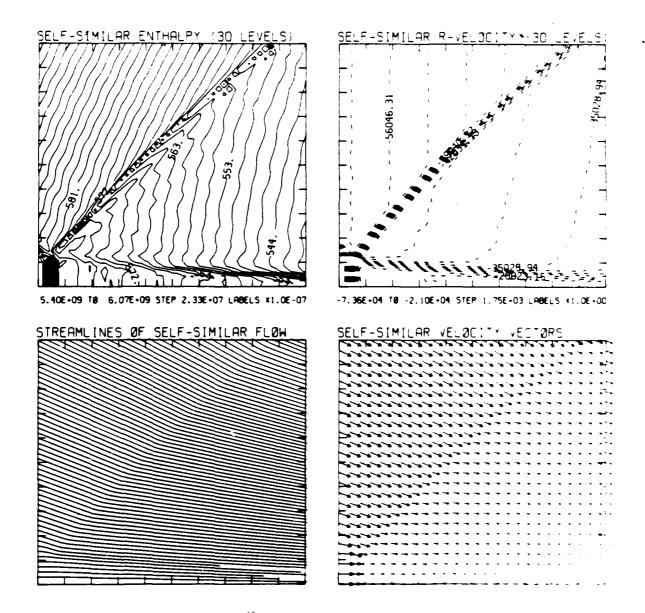
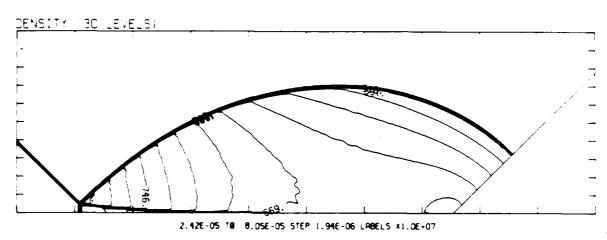
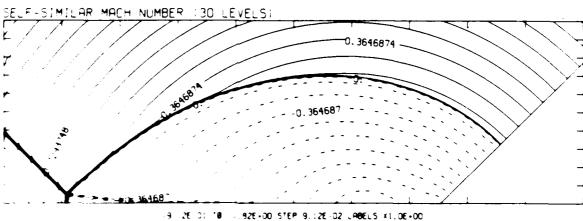


Figure 19.3b. $M_c = 1.50$, blowup-frame plots - continues.

Figure 19. Transition set 1, $\theta_a = 45^{\circ}$, y = 1.4 - continues.

MS= 1.60 ALP=45.00 NR=500 NZ=160 KBEG=125 PS=2.00E+04 PE₹FE0T





wire 13.4a. M_s = 1.60, whole-flowfield contour-plots.

where the continued is $\hat{x}_{ij} = 45^{\circ}$, $\gamma = 1.4$ - continued.

MS= 1.60 ALP=45.00 IL=397 IR=449 UT= 49 PC=2.00E+04 PERFECT

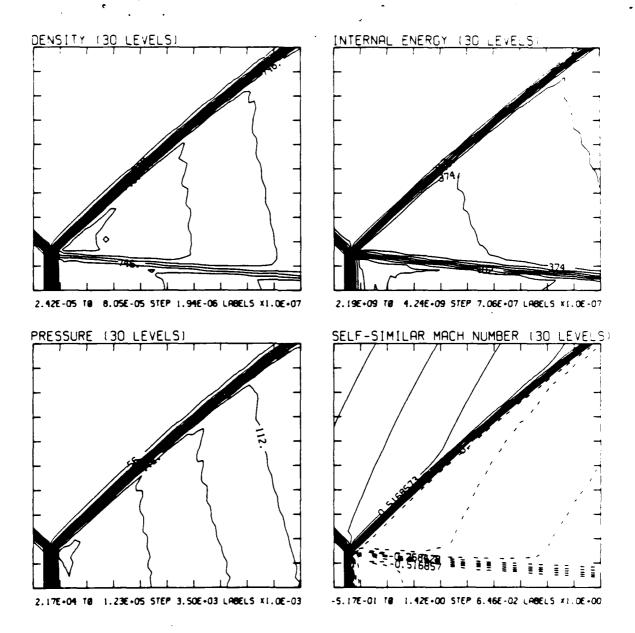


Figure 19.4b. $M_S = 1.60$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.60 ALP=45.00 [L=397 [R=449 UT= 49 PC=2.00E+04 PFRFF07

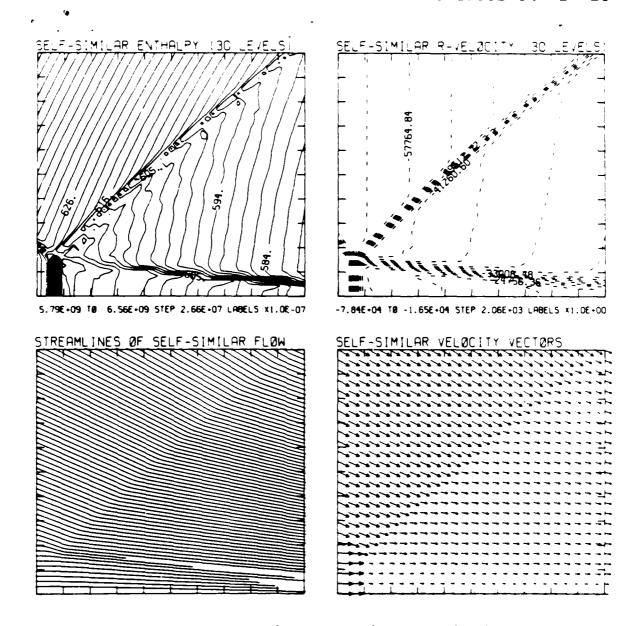
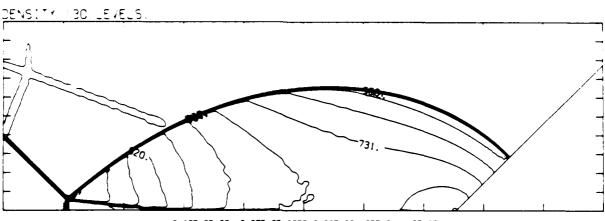


Figure 19.4b. $M_s = 1.60$, blowup-frame plots - continued.

Figure 19. Transition set 1, $q_{\rm w} = 45^{\circ}$, $\gamma = 1.4$ - continued.

MS= 1.70 ALP=15.00 NP=500 NZ=160 KBEG=125 PG=2.00E+04 PERFECT



2.43E-05 TB 8.87E-05 STEP 2.22E-06 LABELS X1.0E+07

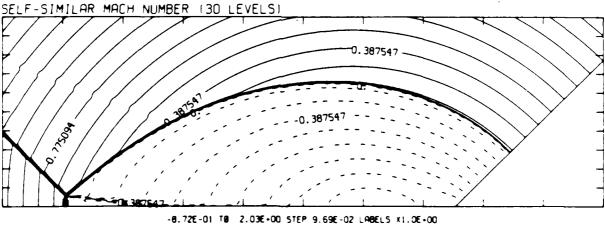


Figure 19.5a. $M_S = 1.70$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.70 ALP=45.00 IL=398 IR=450 UT= 49 P0=2.00E+04 PERFECT

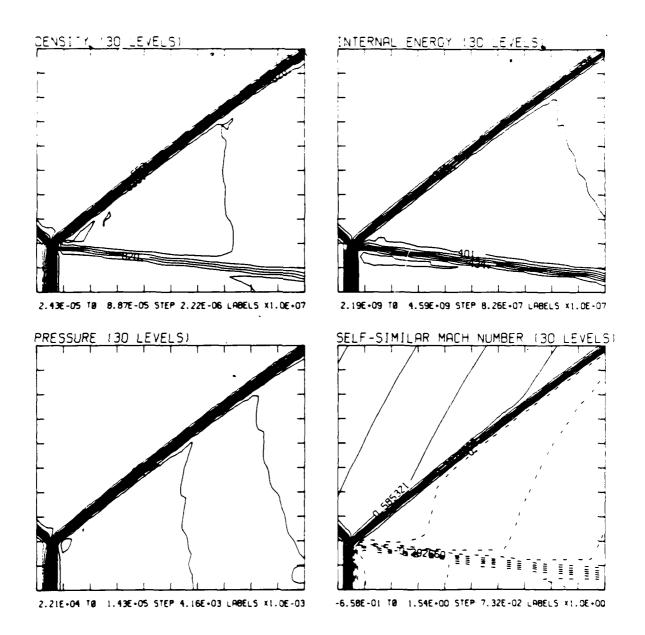


Figure 19.5b. $M_S = 1.70$, blowup-frame plots.

Figure 19. Transition set 1, θ_w = 45°, γ = 1.4 - continued.

MS= 1.70 ALP=45.00 IL=398 IR=450 UT= 49 PC=2.00E+04 PERFECT

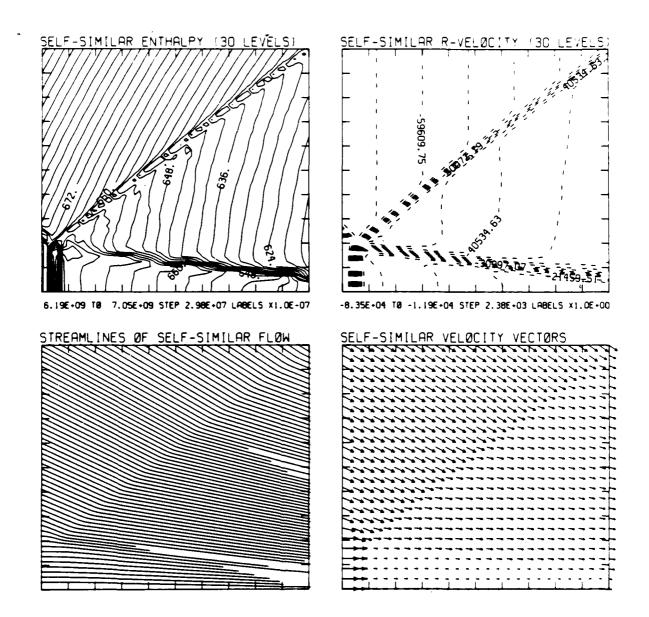
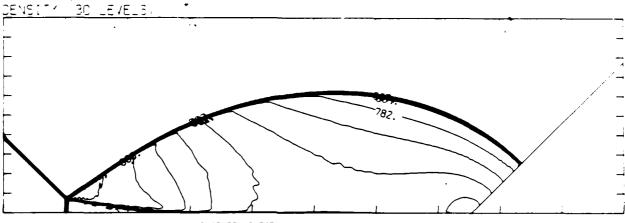


Figure 19.5b. $M_s = 1.70$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.80 ALP=45.00 NF=500 NZ=160 KBEG=105 PC=2.00E+34 PERFECT



2.44E-05 TØ 9.54E-05 STEP 2.45E-06 LABELS X1.0E+07

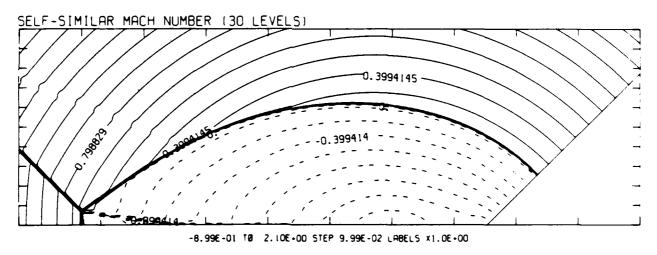


Figure 19.6a. M_s = 1.80, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.80 ALP=45.00 [L=400 [R=453 UT= 50 PO=2.00E+04 PERFECT

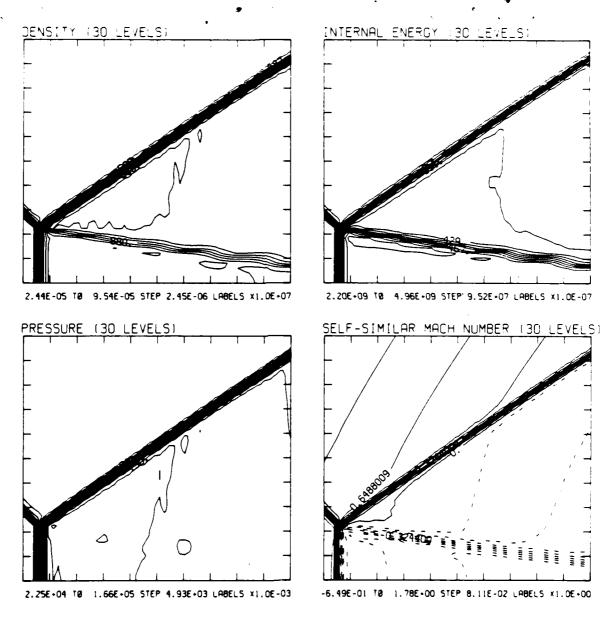


Figure 19.6b. $M_S = 1.80$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.80 ALP=45.00 IL=400 IR=453 UT= 50 PC=2.00E+24 PERFECT

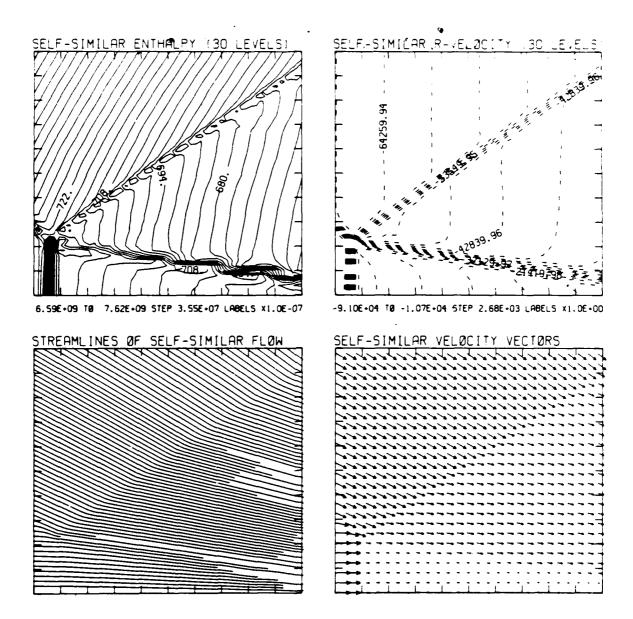
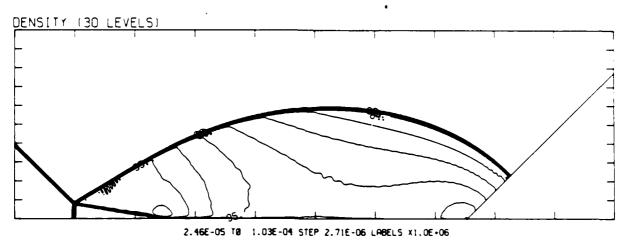


Figure 19.6b. $M_S = 1.80$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.90 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+34 PERFECT



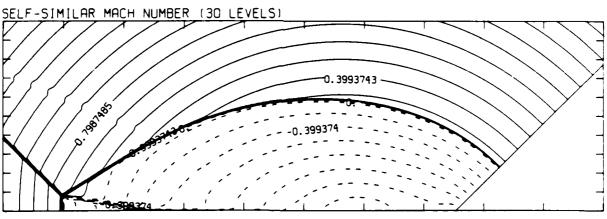


Figure 19.7a. $M_s = 1.90$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.90 ALP=45.00 IL=401 IR=454 UT= 50 PC=2.00E+04 PERFECT

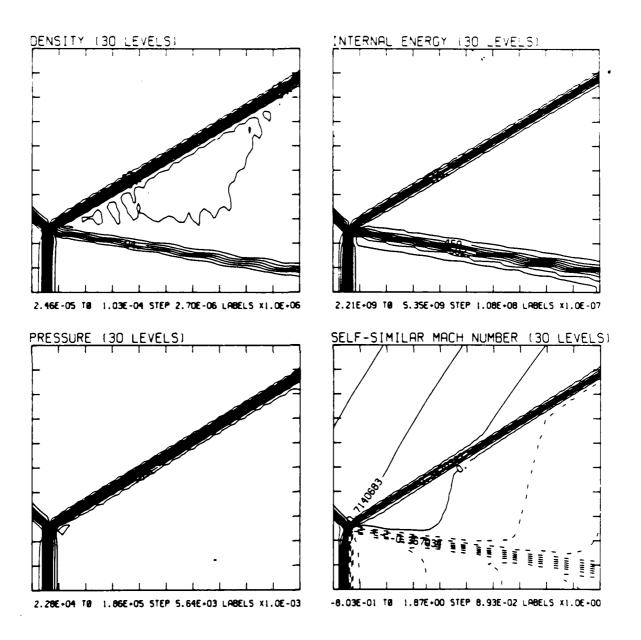


Figure 19.7b. $M_s = 1.90$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 1.90 ALP=45.00 [L=401 [R=454]UT= 50 PG=2.00E+04 PERFECT

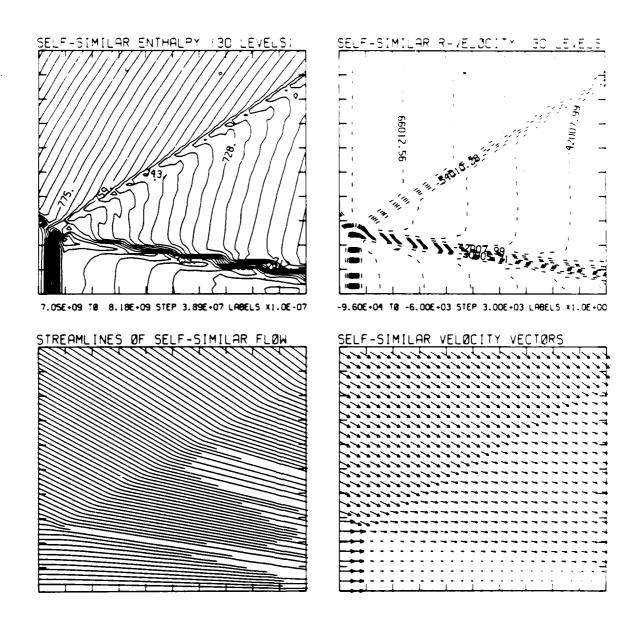
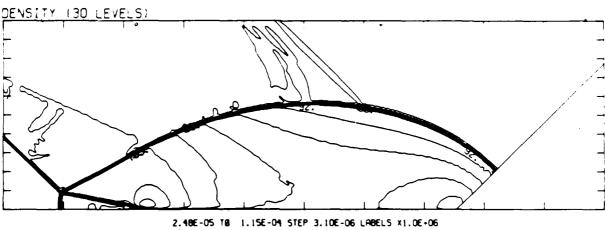


Figure 19.7b. $M_s = 1.90$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.00 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 PERFECT



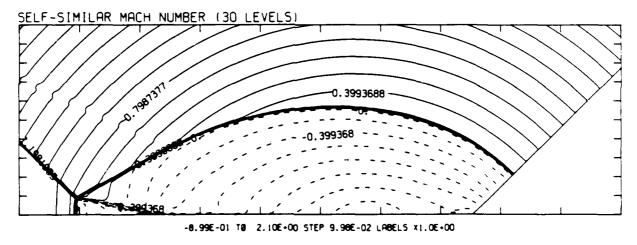


Figure 19.8a. $M_S = 2.00$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

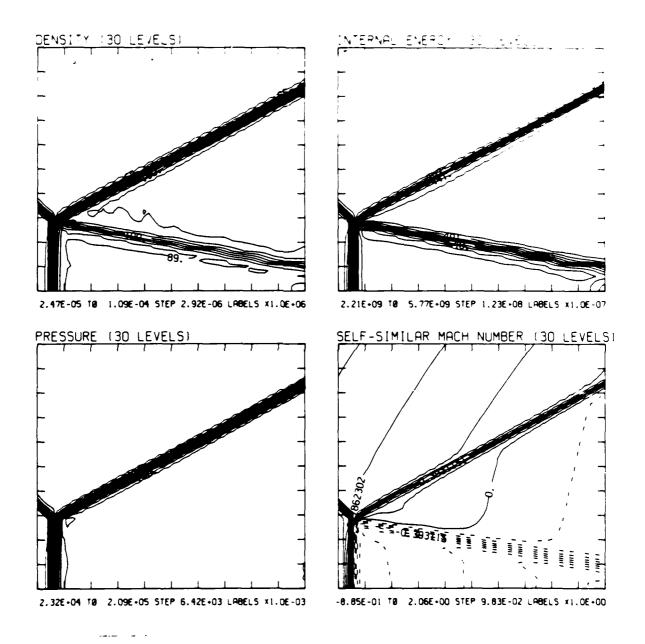


Figure 19.8b. $M_S = 2.00$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

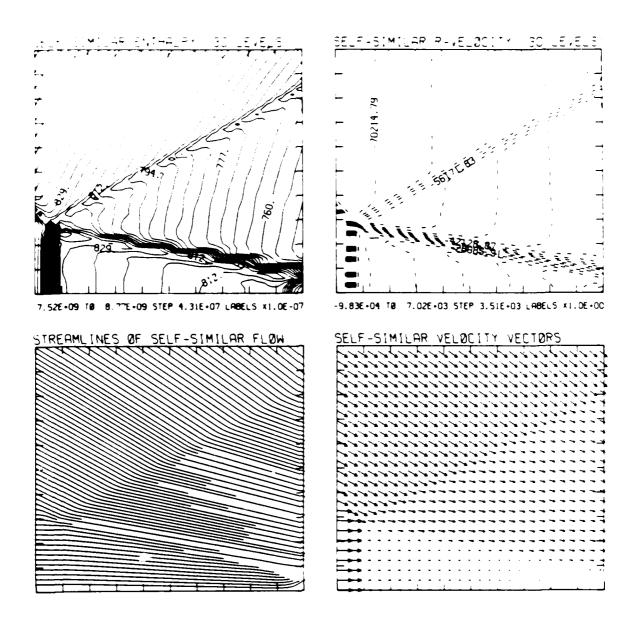
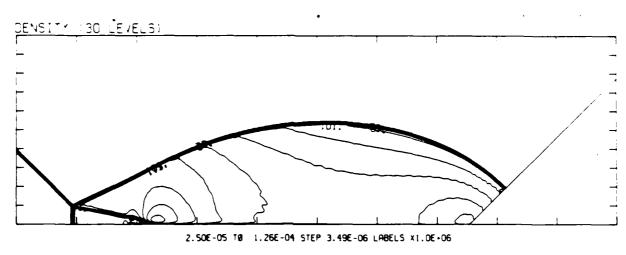


Figure 19.8b. $M_s = 2.00$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 2.10 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+04 PERFECT



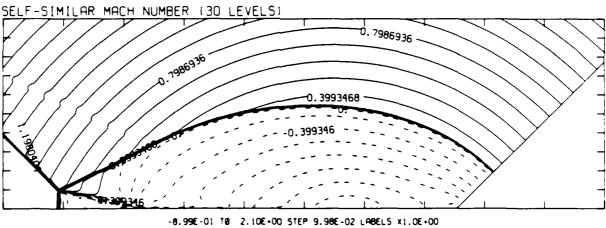


Figure 19.9a. $M_s = 2.10$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.10 ALP=45.00 [L=404 [P=457 UT= 50 PO=2.00E+04 REREEDT

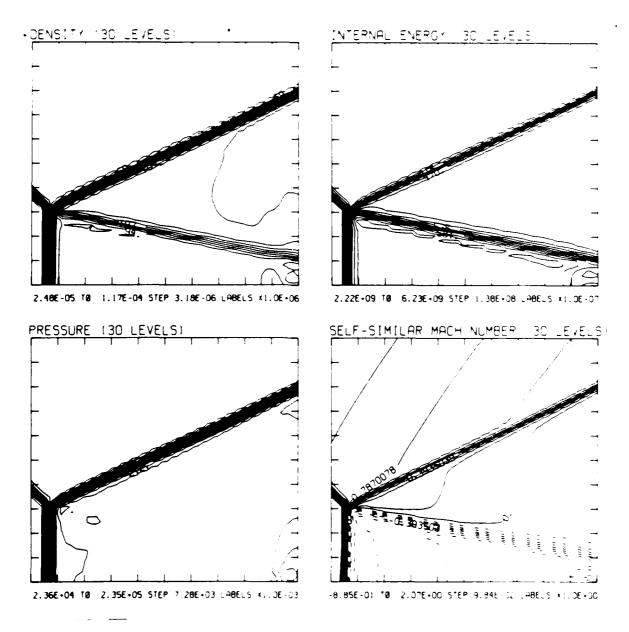


Figure 19.9b. $M_s = 2.10$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_w = 45^{\circ}$, $\gamma = 1.4$ - continued.

MS= 2.10 ALP=45.00 [L=404 [P=457 UT= 50 PO=2.00E+04 PERFECT

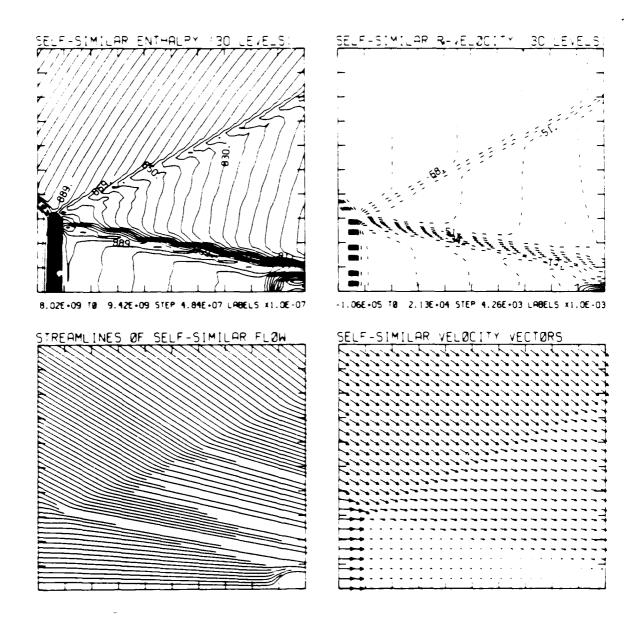
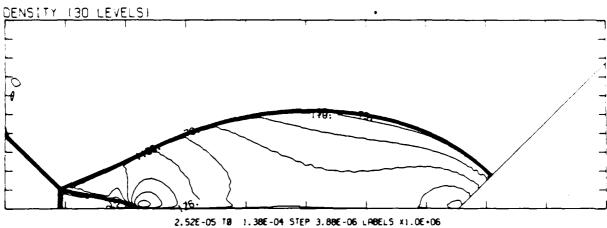


Figure 19.9b. $M_s = 2.10$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.20 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 PERFECT



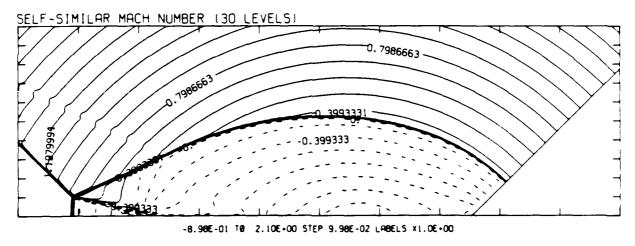


Figure 19.10a. $M_s = 2.20$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.20 ALP=45.00 [L=405 [R=458 UT= 50 PC=2.00E+04 PERFECT

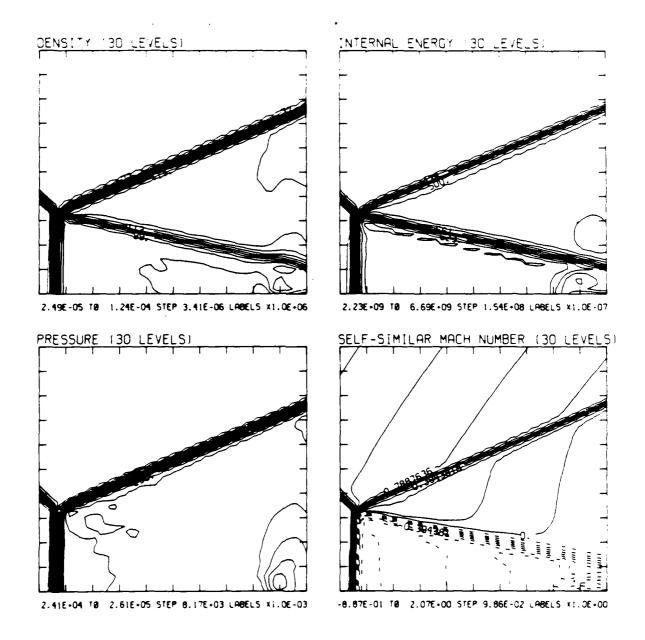


Figure 19.10b. $M_S = 2.20$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.20 ALP=45.00 [L=405 [R=458 UT= 50 P0=2.00E+04 PERFECT

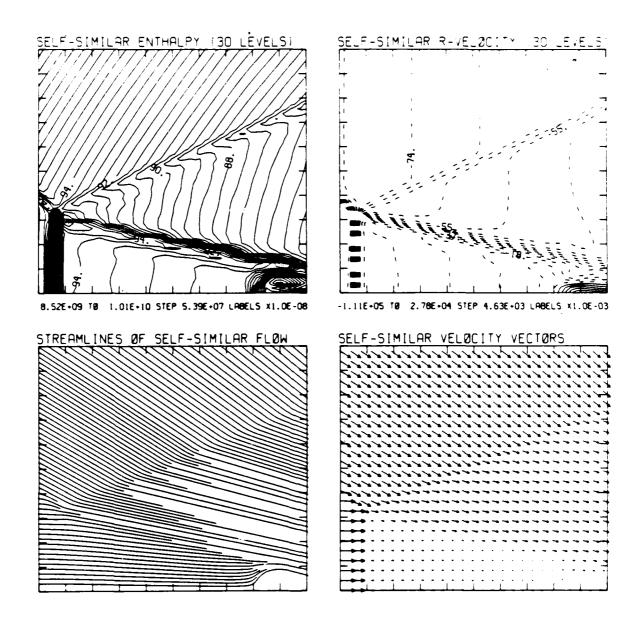
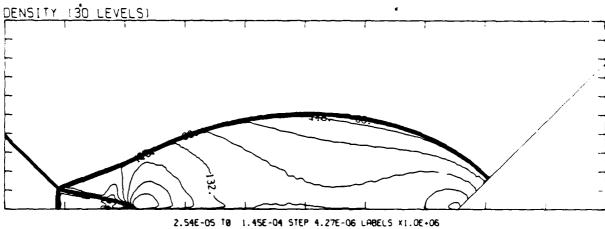


Figure 19.10b. $M_s = 2.20$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.30 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 PERFECT



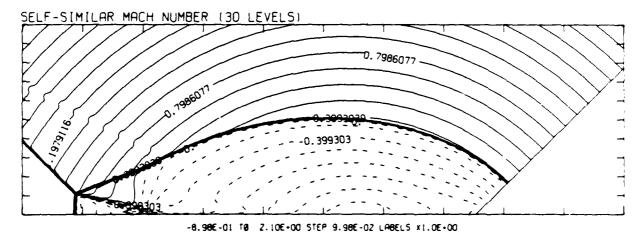


Figure 19.11a. $M_S = 2.30$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.30 ALP=45.00 IL=406 IR=459 UT= 50 PC=2.00E+04 PERFECT

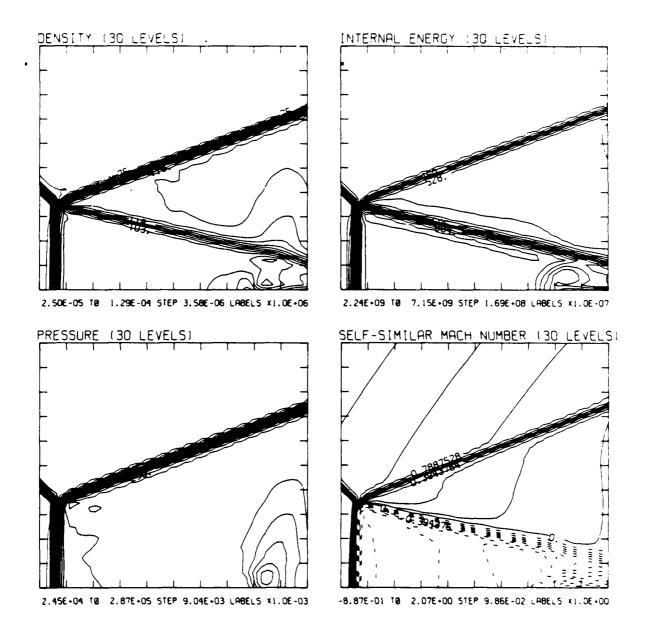


Figure 19.11b. $M_s = 2.30$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.30 ALP=45.00 IL=406 [R=459 UT= 50 PO=2.00E+04 PERFECT

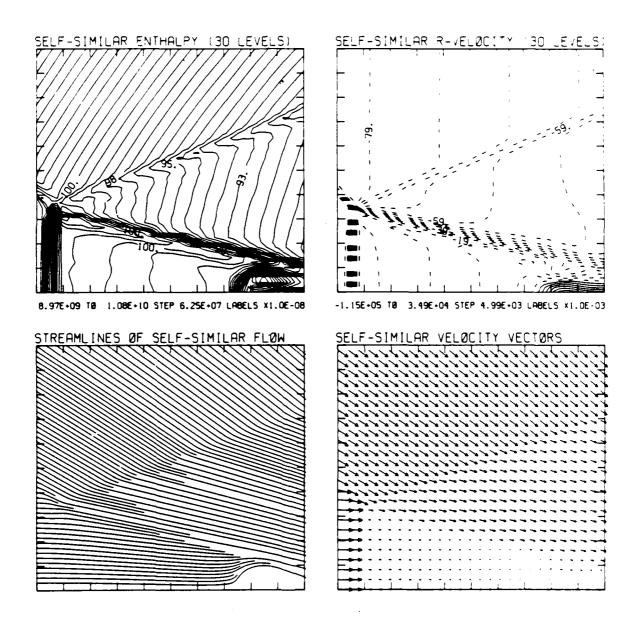
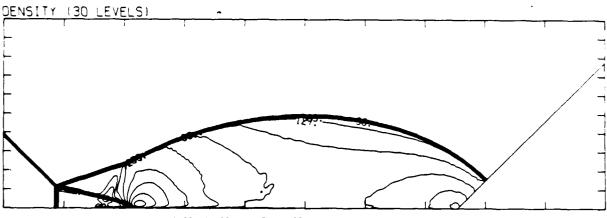


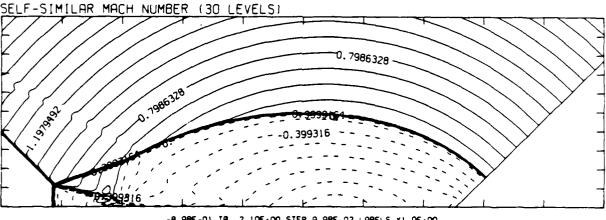
Figure 19.11b. $M_s = 2.30$, blowup-frame plots - continued.

Figure 19. Transition set 1, θ_w = 45°, γ = 1.4 - continued.

MS= 2.40 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 PERFECT



2.56E-05 T0 1.60E-04 STEP 4.65E-06 LRBELS X1.0E+06



-8.98E-01 T0 2.10E+00 STEP 9.98E-02 LABELS X1.0E+00

Figure 19.12a. $M_S = 2.40$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.40 ALP=45.00 IL=406 IR=460 JT= 51 PO=2.00E+04 PERFECT

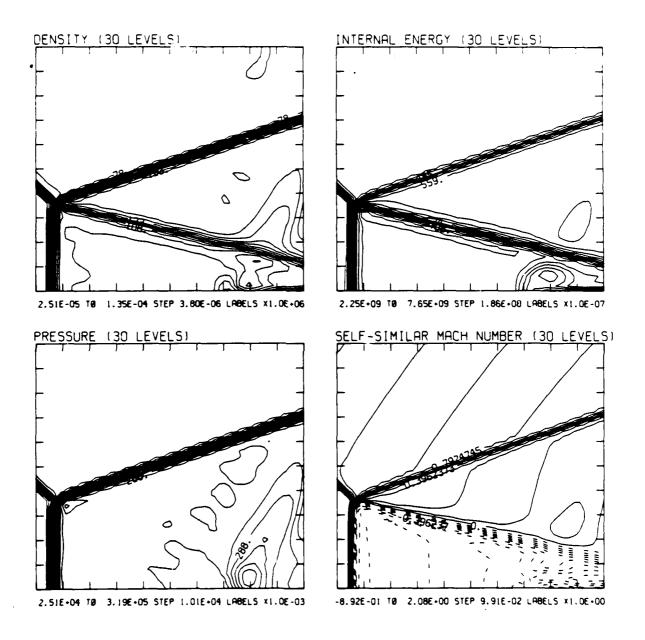


Figure 19.12b. $M_S = 2.40$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

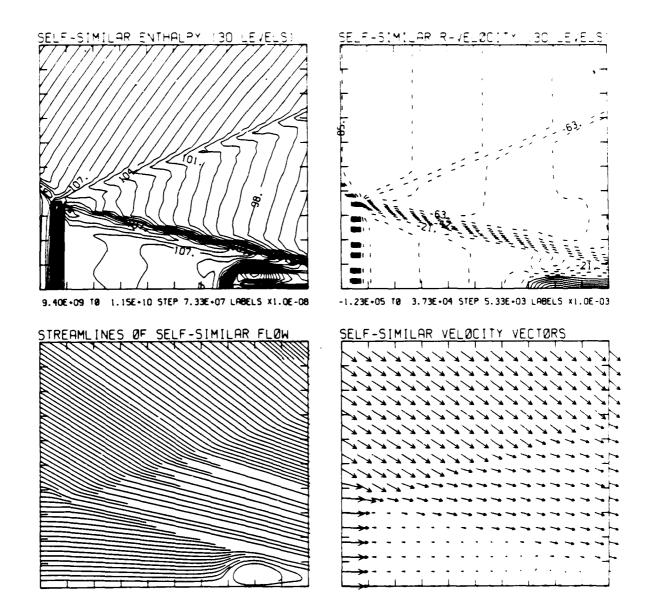
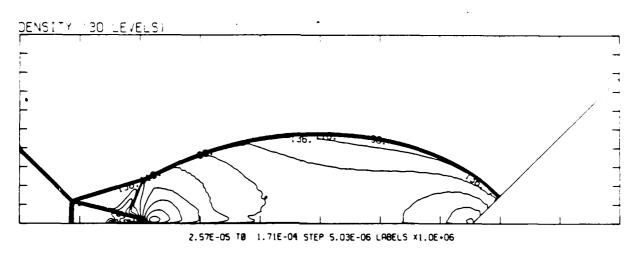


Figure 19.12b. $M_s = 2.40$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 2.50 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+04 PERFECT



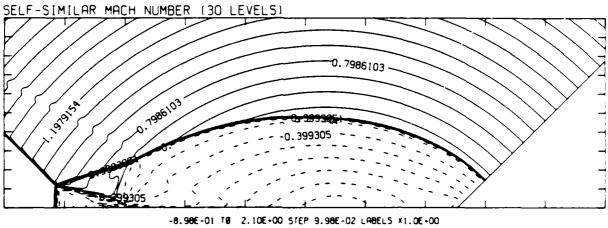


Figure 19.13a. $M_s = 2.50$, whole-flowfield contour-plots.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 2.50 ALP=45.00 [L=406 [R=460 UT= 5] P0=2.00E+04 PERFECT

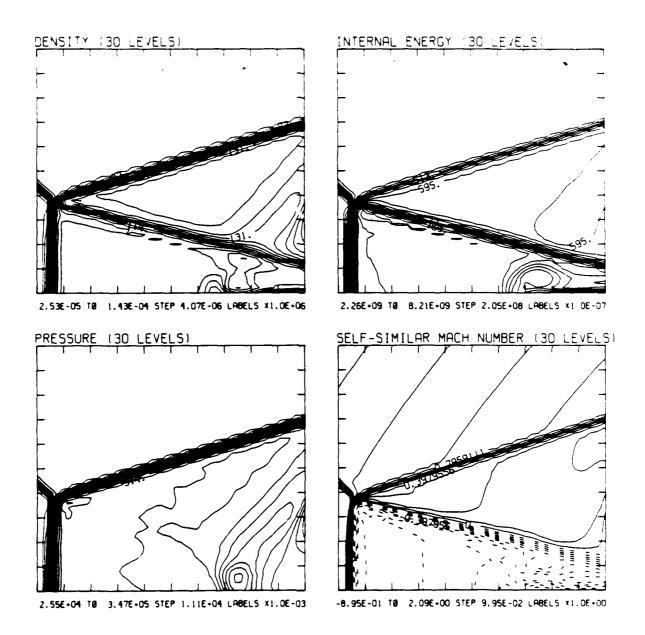


Figure 19.13b. $M_S = 2.50$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm W}$ = 45°, γ = 1.4 - continued.

MS= 2.50 ALP=45.00 IL=406 IR=460 JT= 51 PC=2.00E+04 PERFECT

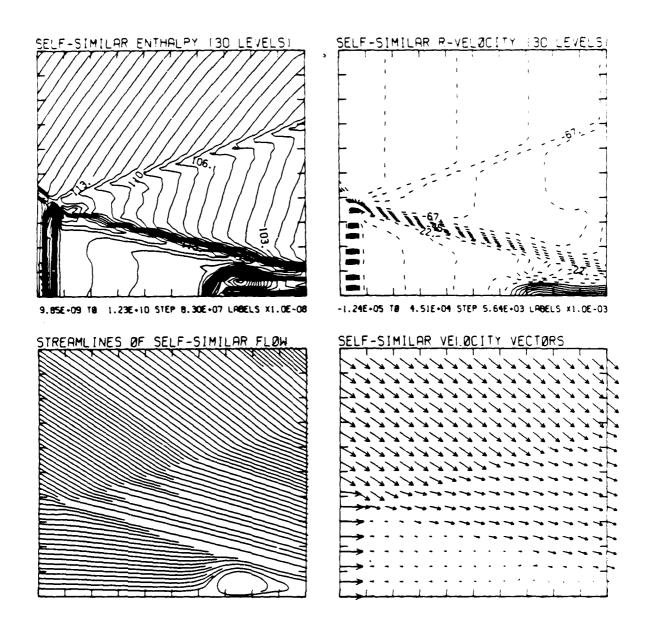
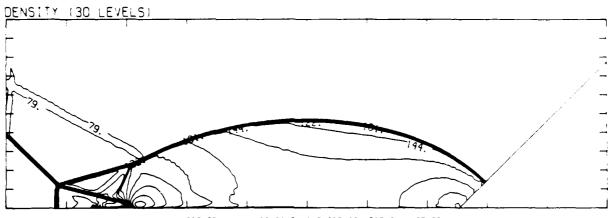


Figure 19.13b. $M_s = 2.50$, blowup-frame plots - continued.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 2.60 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.005+14 FERFEDT



2.59E-05 TØ 1.82E-04 STEP 5.39E-06 LABELS X1.0E+06

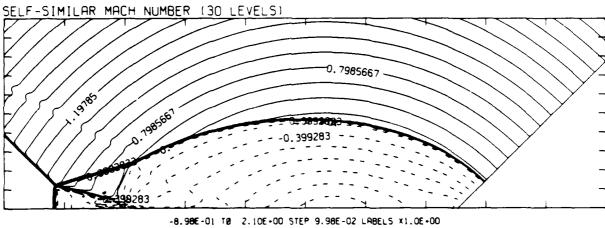


Figure 19.14a. $M_S = 2.60$, whole-flowfield contour-plots.

Figure 19. Transition set 1, θ_{W} = 45°, γ = 1.4 - continued.

MS= 0.60 ALR445.00 (L=407 (R=46) UT= 5) PD=0.00E+04 PERFECT

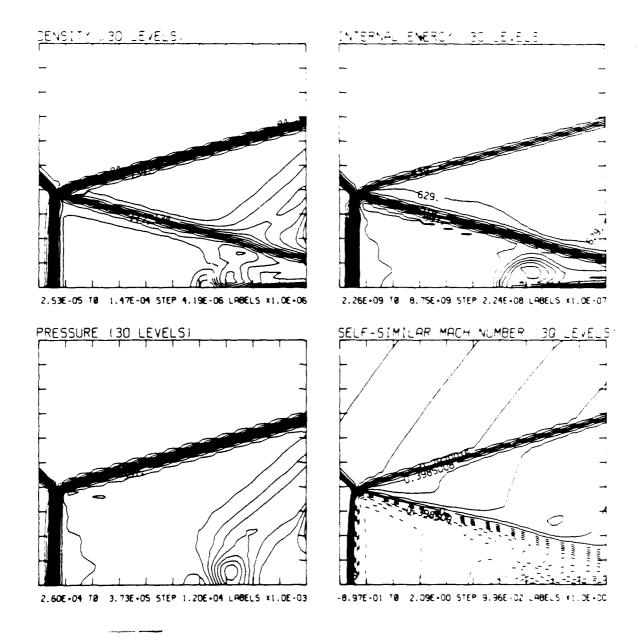


Figure 19.14b. $M_s = 2.60$, blowup-frame plots.

Figure 19. Transition set 1, $\theta_{\rm w}$ = 45°, γ = 1.4 - continued.

MS= 2.60 ALP=45.00 [L=407 [P=46] UT= 5] PS=2.00E+34 RERREST

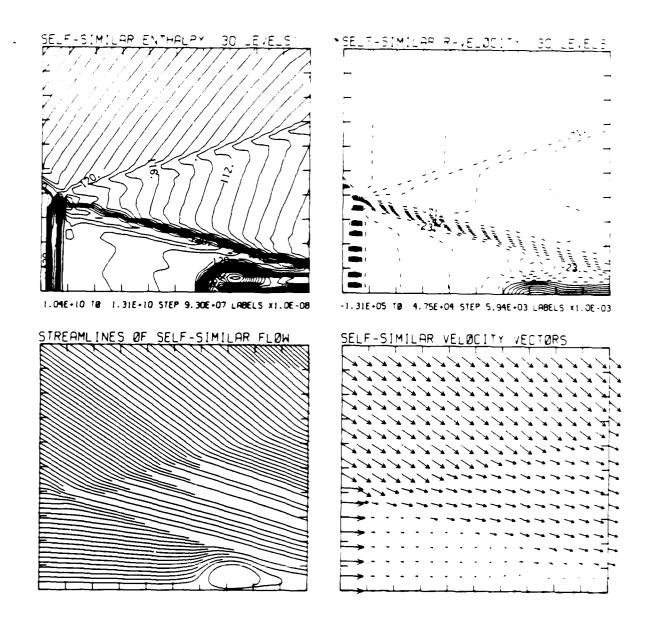
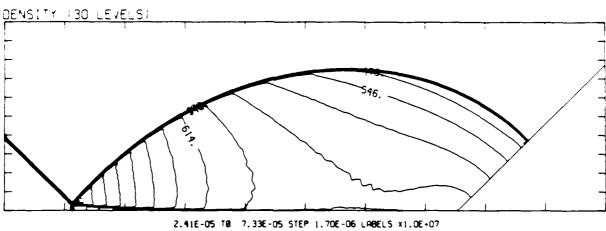


Figure 19.14b. $M_S = 2.60$, blowup-frame plots - continued.

Figure 19. Transition set 1, $9_W = 45^\circ$, $\gamma = 1.4$ ~ continued.

MS= 1.50 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 HANSEN



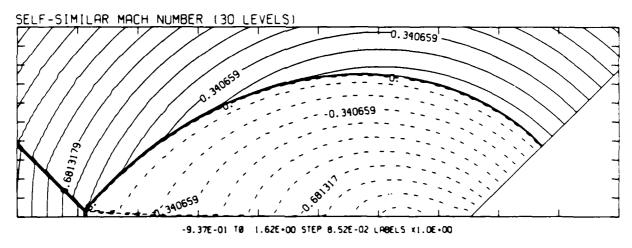


Figure 20.1a. $M_S = 1.50$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen.

MS= 1.50 ALP=45.00 IL=395 IR=447 UT= 49 PG=2.00E+04 HANSEN

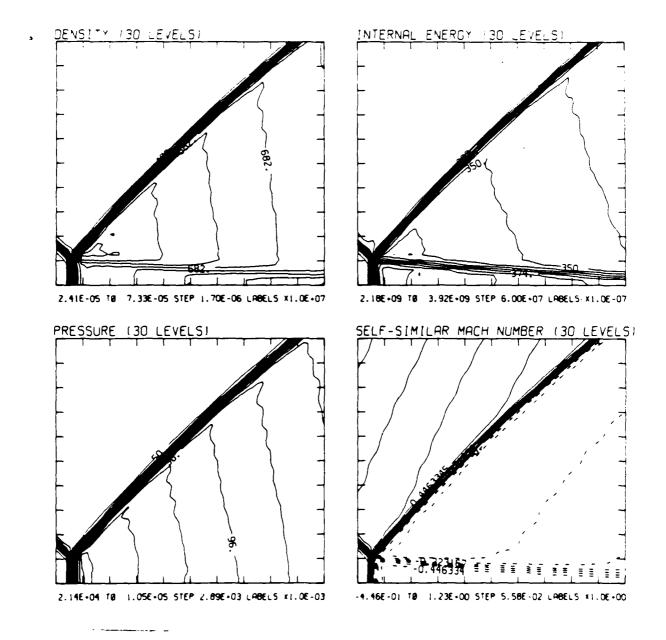


Figure 20.1b. $M_S = 1.50$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm w}$ = 45°, Hansen - continued.

MS= 1.50 ALP=45.00 [L=395 [R=447 UT= 49 PO=2.00E+04 HANSEN

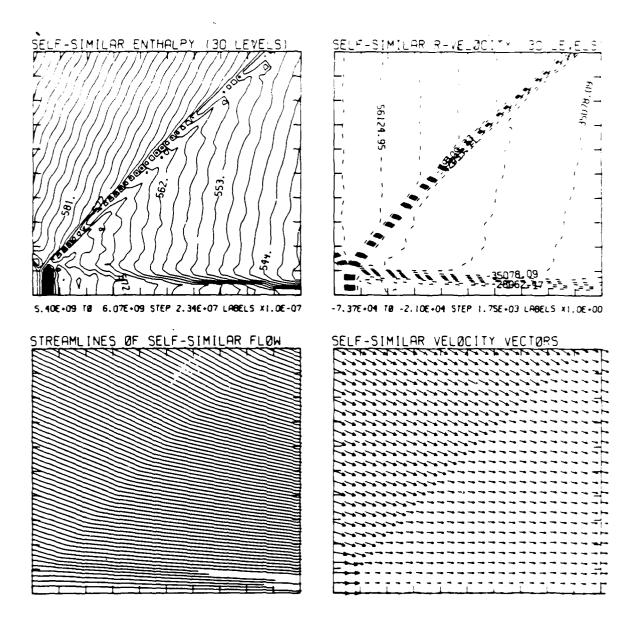
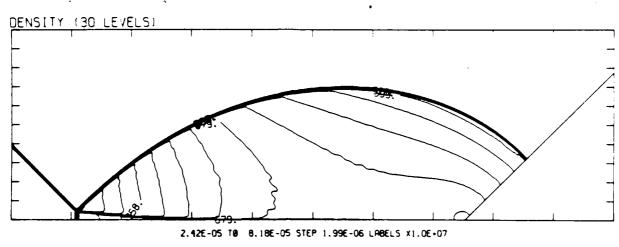


Figure 20.1b. $M_s = 1.50$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.60 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+04 HANSEN



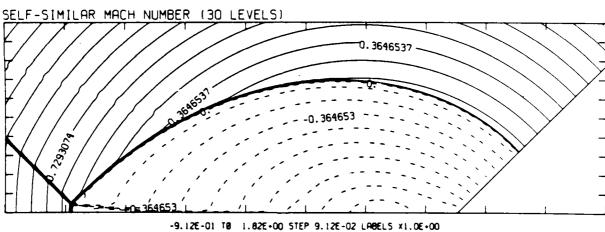


Figure 20.2a. $M_S = 1.60$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

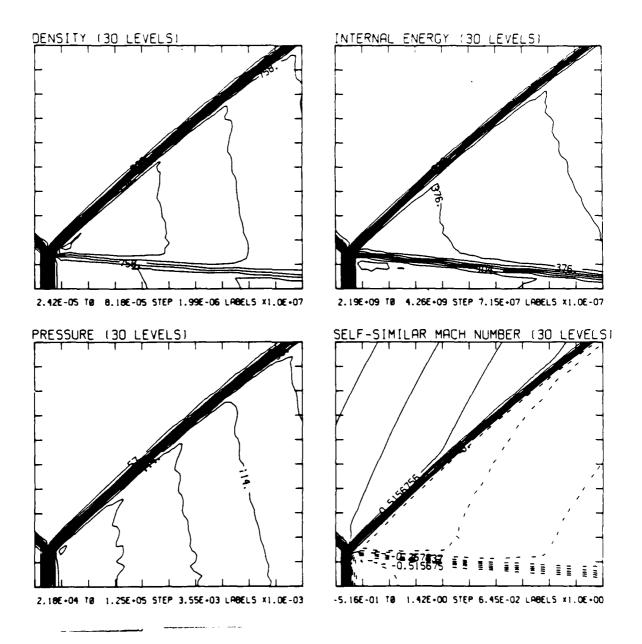


Figure 20.2b. $M_s = 1.60$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm w}$ = 45°, Hansen - continued.

MS= 1.60 ALP=45.00 IL=396 IR=448 UT= 49 P0=2.00E+04 HANSEN

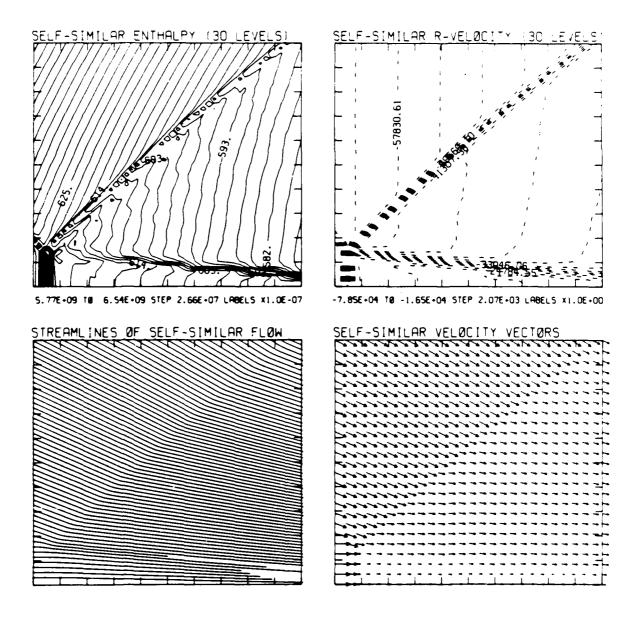
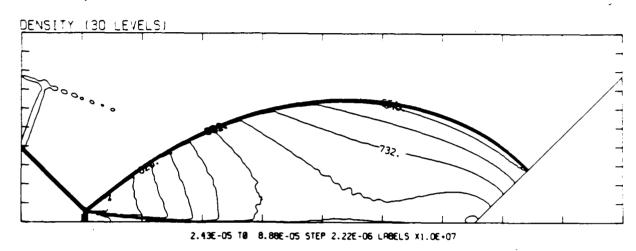


Figure 20.2b. $M_s = 1.60$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm w}$ = 45°, Hansen - continued.

MS= 1.70 ALP=45.00 NR=500 NZ=160 KBEG=125 PG=2.00E+04 HANSEN



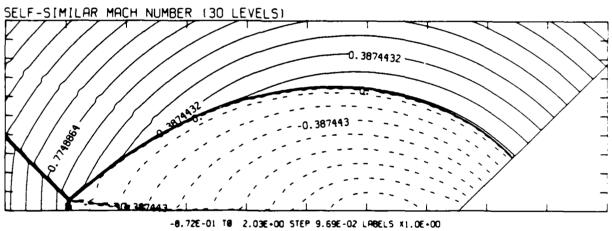


Figure 20.3a. $M_S = 1.70$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.70 ALP=45.00 IL=398 IR=450 JT= 49 PC=2.00E+04 HANSEN

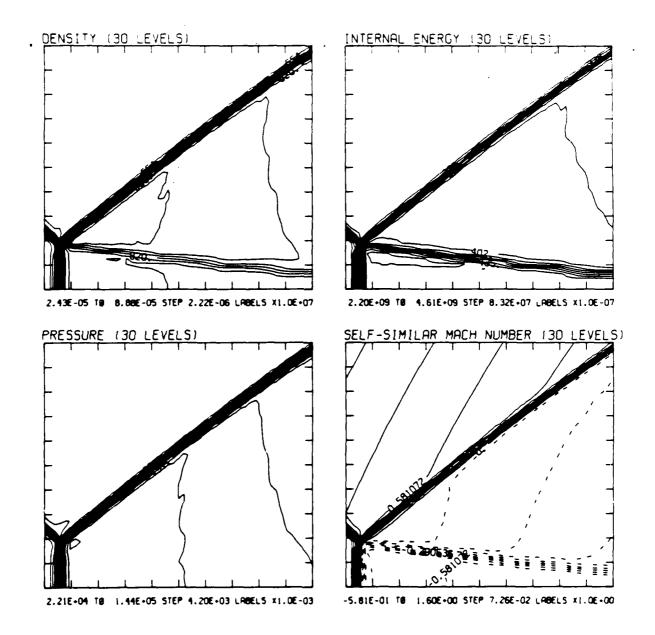


Figure 20.3b. $M_s = 1.70$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.70 ALP=45.00 IL=398 IR=450 UT= 49 PO=2.00E+04 HANSEN

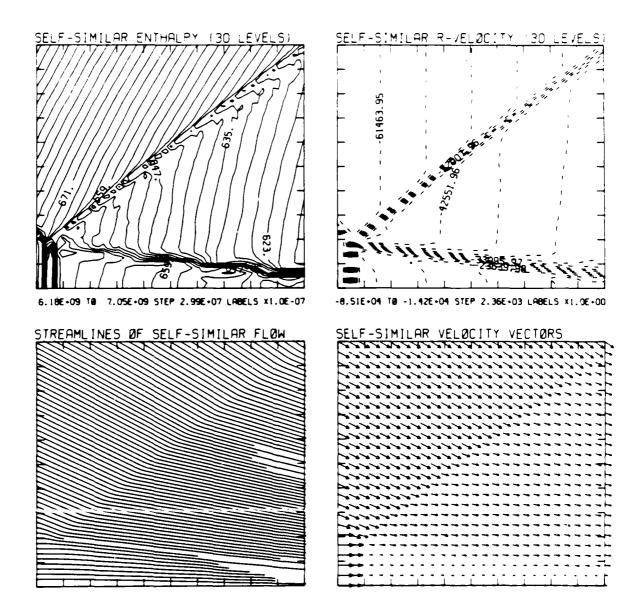
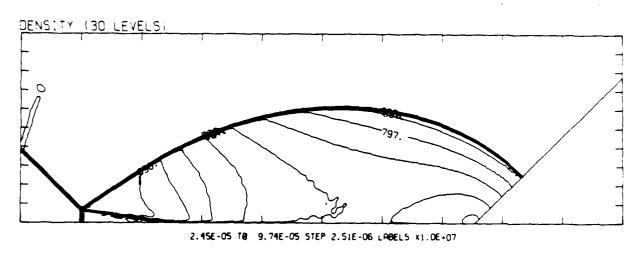


Figure 20.3b. $M_s = 1.70$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.80 ALP=45.00 NR=500 NZ=160 KBEG=125 P0=2.00E+04 HANSEN



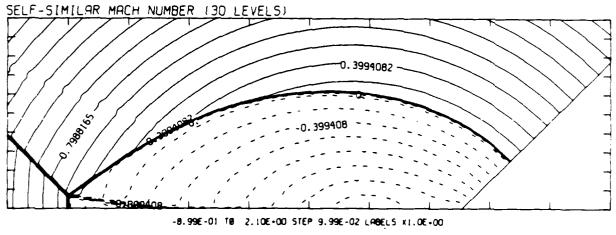


Figure 20.4a. $M_S = 1.80$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

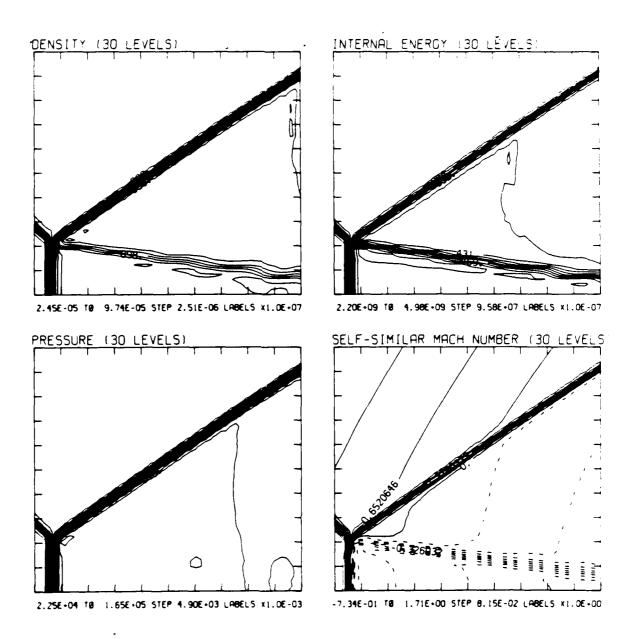


Figure 20.4b. $M_s = 1.80$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.80 ALP=45.00 [L=400 [R=453 JT= 50 PO=2.00E+04 HANSEN

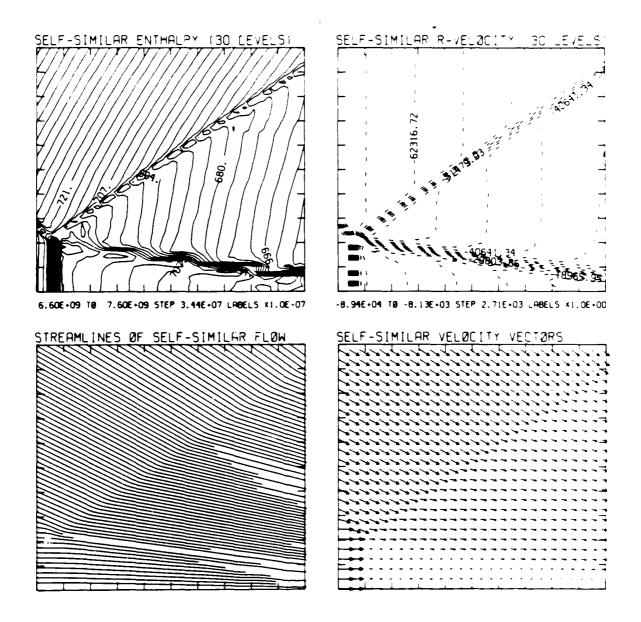
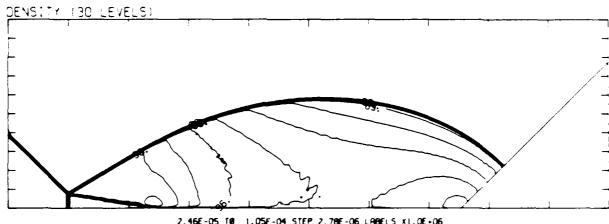


Figure 20.4b. $M_s = 1.80$, blowup-frame plots - continued.

Figure 20. Transition set 1, $q_{\rm w} = 45^{\circ}$, Hansen - continued.

MS= 1.90 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 HANSEN



2.46E-05 TO 1.05E-04 STEP 2.78E-06 LABELS XI.0E+06

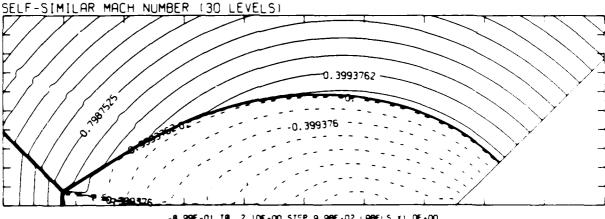


Figure 20.5a. $M_s = 1.90$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

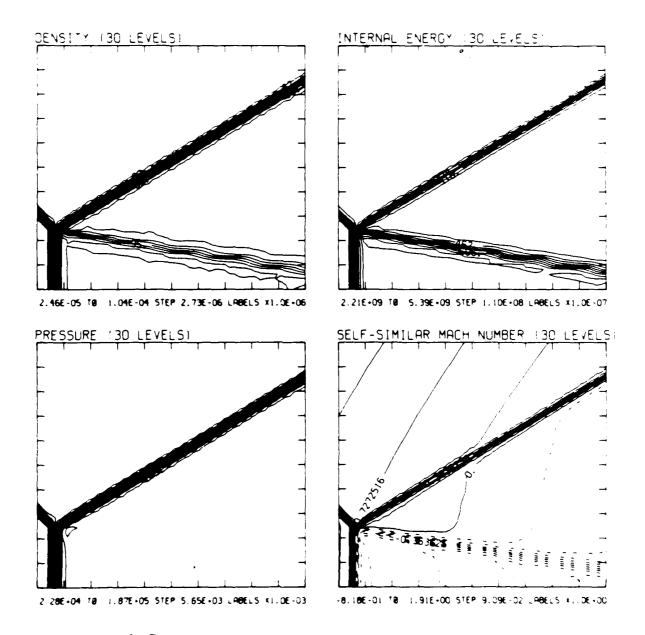


Figure 20.5b. $M_s = 1.90$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 1.90 ALP=45.00 [L=40] [R=454 UT= 50 PG=2.00E+04 HANSE'.

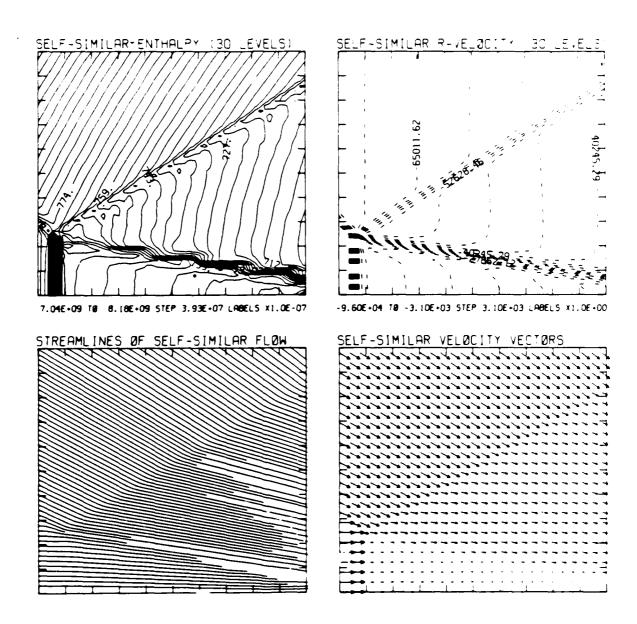
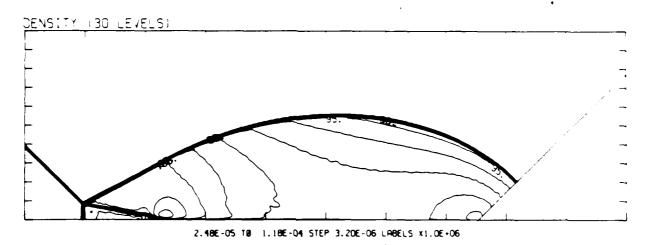


Figure 20.5b. $M_s = 1.90$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.00 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 HANSEN



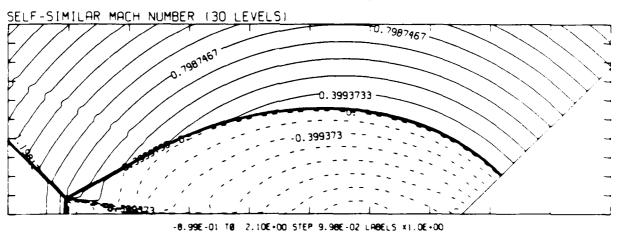


Figure 20.6a. $M_s = 2.00$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $A_{w} = 45^{\circ}$, Hansen - continued.

MS= 2.00 ALP=45.00 IL=402 IR=455 UT= 50 PO=2.00E+04 HPNSEN

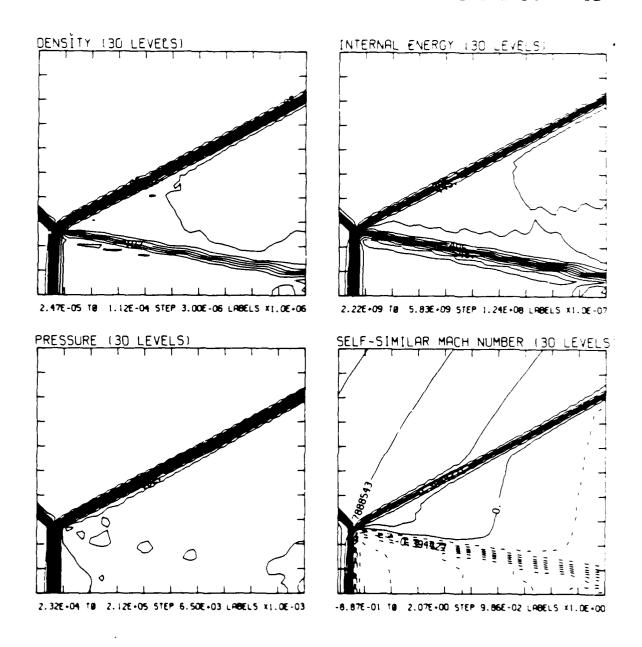


Figure 20.6b. $M_s = 2.00$, blowup-frame plots.

Figure 20. Transition set 1, $q_{\mu} = 45^{\circ}$, Hansen - continued.

MS= 2.00 ALP=45.00 IL=402 IR=455 UT= 50 PG=2.00E+34 HANSEN

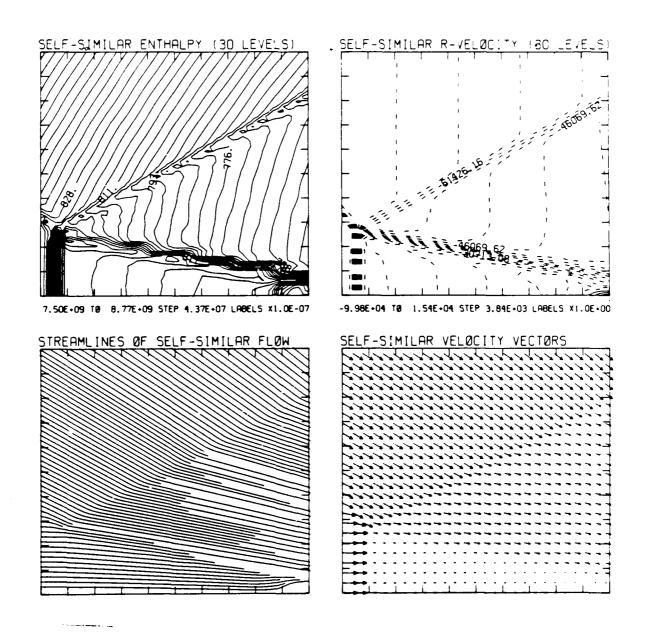
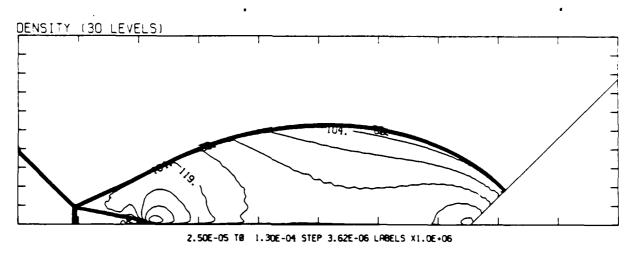


Figure 20.6b. $M_c = 2.00$, blowup-frame plots - continued.

Figure 20. Transition set 1, θ_{u} = 45°, Hansen - continued.

MS= 2.10 ALP=45.00 NR=500 NZ=160 KBEG=125.P0=2.00E+04 HANSEN



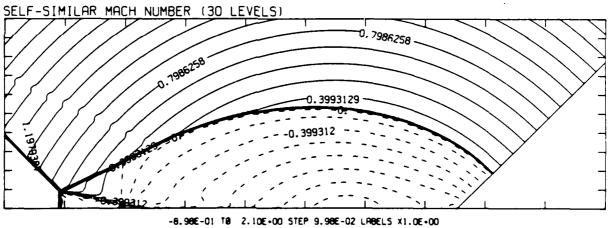


Figure 20.7a. $M_s = 2.10$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

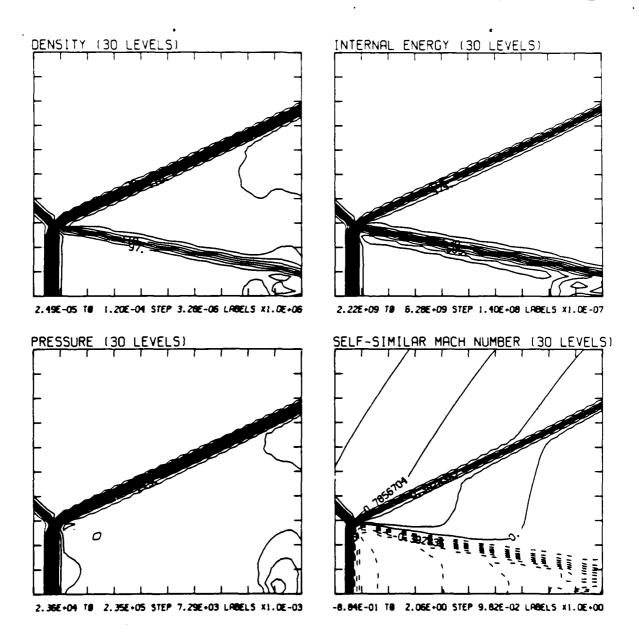


Figure 20.7b. $M_s = 2.10$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.10 ALP=45.00 IL=404 IR=457 JT= 50 PO=2.00E+04 HANSEN

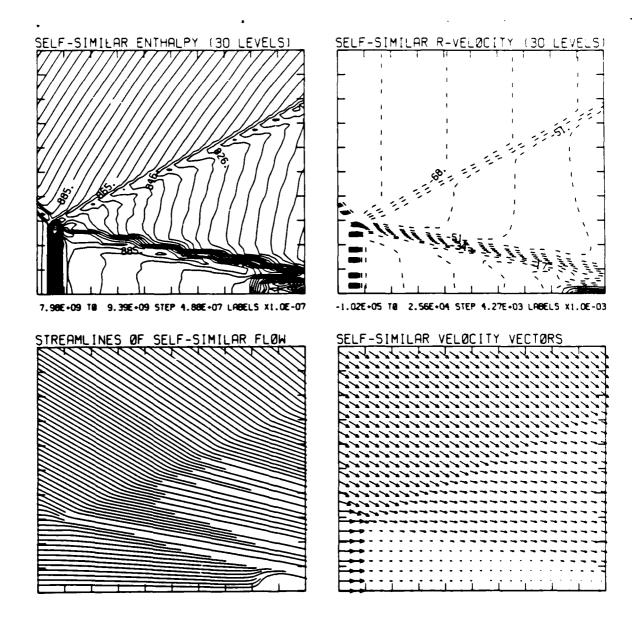
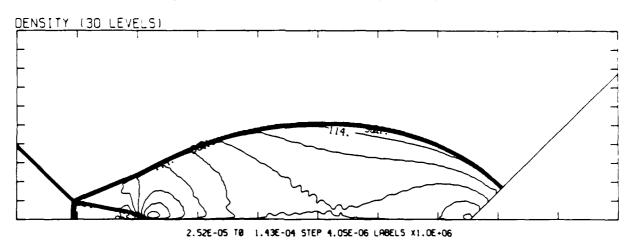
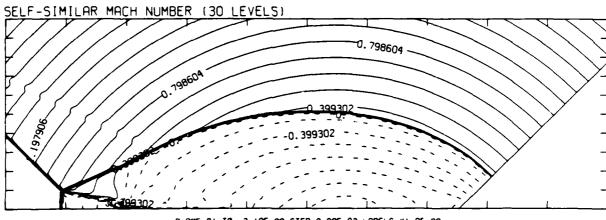


Figure 20.7b. $M_s \approx 2.10$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.20 ALP=45.00 NR=500 NZ=160 KBEG=125 PO=2.00E+04 HANSEN





-8.98E-01 T0 2.10E+00 STEP 9.98E-02 LRBELS X1.0E+00

Figure 20.8a. $M_s = 2.20$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.20 ALP=45.00 IL=404 IR=457 UT= 50 PG=2.00E+04 HANSE'.

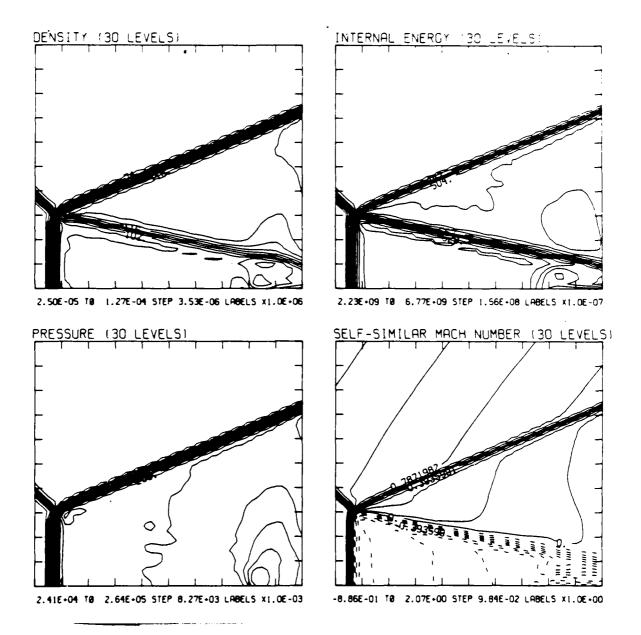


Figure 20.8b. $M_s = 2.20$, blowup-frame plots.

Figure 20. Transition set 1, $e_w = 45^{\circ}$, Hansen - continued.

MS= 2.20 ALP=45.00 IL=404 IR=457 JT= 50 PG=2.00E+64 HANSEN

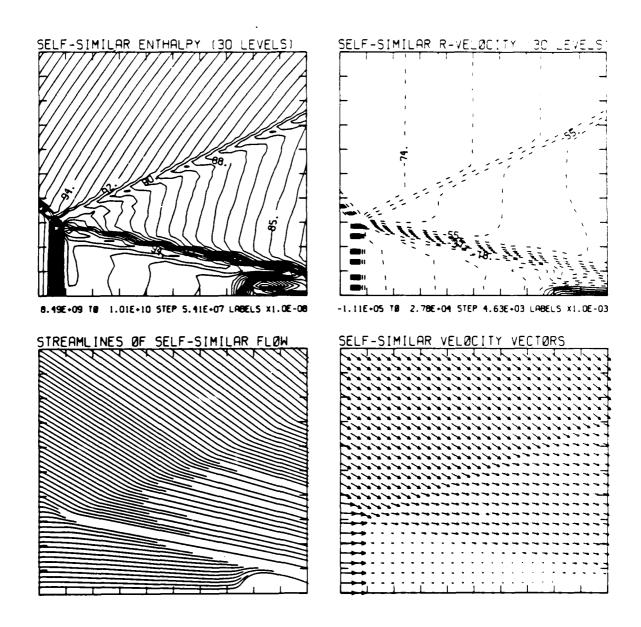
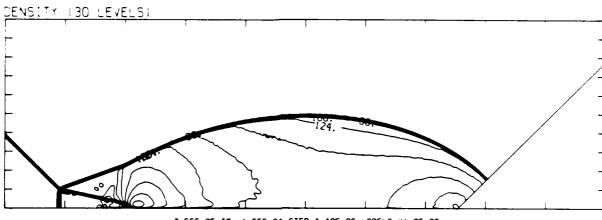


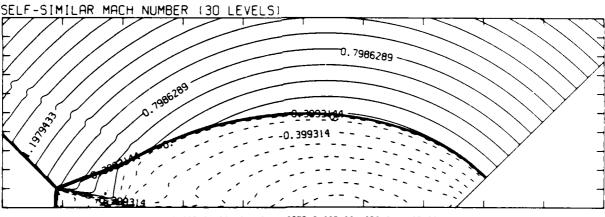
Figure 20.8b. $M_S = 2.20$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.30 ALP=45.00 NR=500 NZ=160 KBEG=125 PC=2.00E+04 HANSEN



2.55E-05 TØ 1.55E-04 STEP 4.48E-06 LABELS X1.0E+06



-8.98E-01 TØ 2.10E+00 STEP 9.98E-02 LABELS X1.0E+00

Figure 20.9a. $M_S = 2.30$, whole-flowfield contour-plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.30 ALP=45.00 IL=405 IR=458 UT= 50 P0=2.00E+04 HANSEN

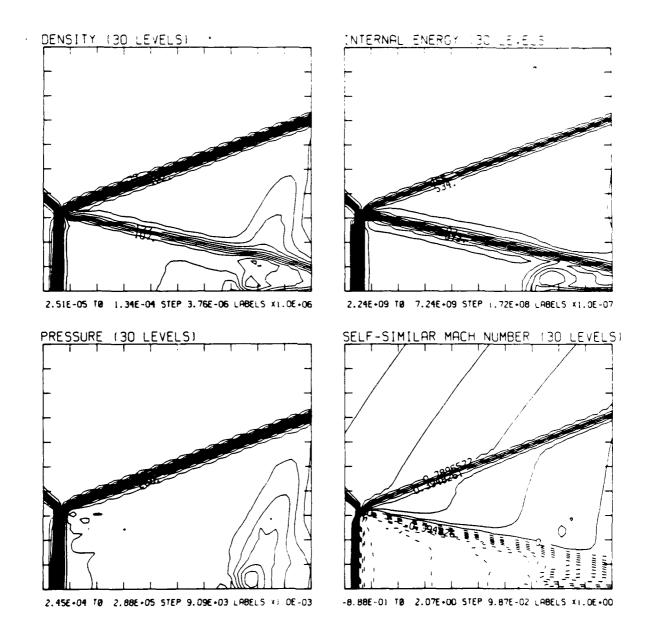


Figure 20.9b. $M_s = 2.30$, blowup-frame plots.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 2.80 ALP=45.00 [L=4]5 [P=458 UT= 50 P0=0.00E+04 HANSEN

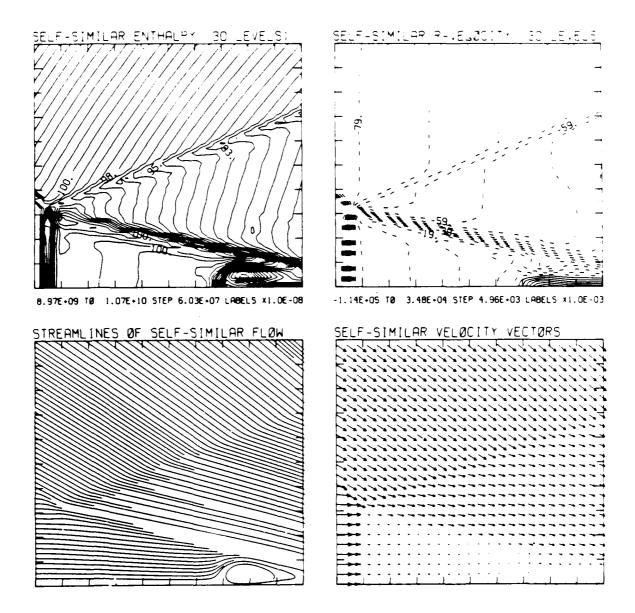
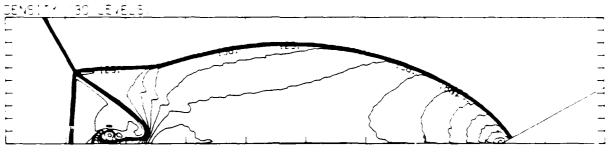


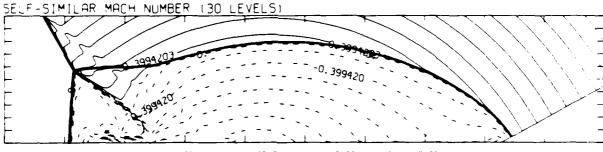
Figure 20.9b. $M_s = 2.30$, blowup-frame plots - continued.

Figure 20. Transition set 1, $\theta_{\rm W}$ = 45°, Hansen - continued.

MS= 4.00 PLP=29.00 NP=510 NZ=110 KBEG= BÖ PD=2.00E+04 FERFEDT



2.69E-05 TB 2.40E-04 STEP 7.33E-06 LABELS XI.0E+06



-8.99E-01 TØ 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 21.1a. $\theta_{\rm w}$ = 29°, whole-flowfield contour-plots.

Figure 21. Transition set 2, $M_S = 4.0$, $\gamma = 1.4$.

MS= 4.00 ALR=09.00 (L=38) [R=458 LT= T4 R0=0.00E+04 RERFED]

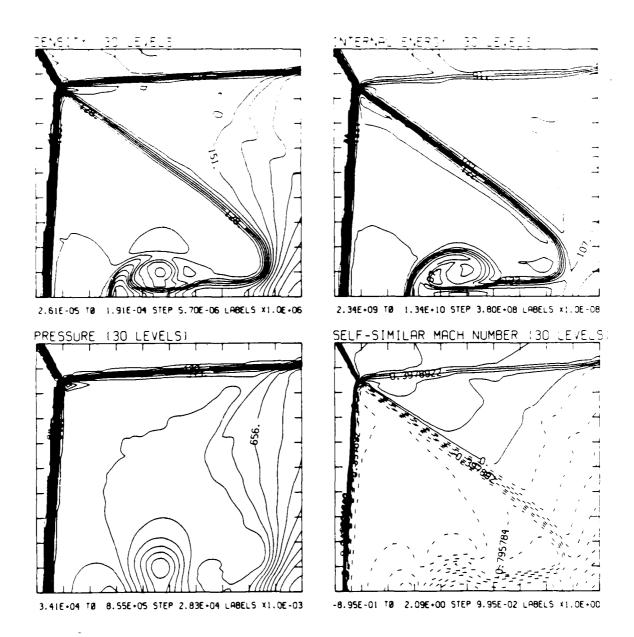


Figure 21.1b. $\theta_w = 29^{\circ}$, blowup-frame plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=29.00 [L=38] [R=458 UT= T4 P0=2.00E+34 REFFEDT

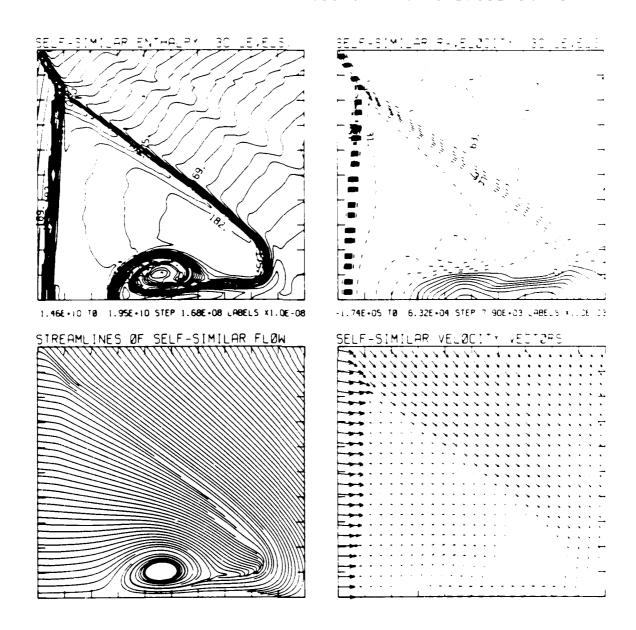
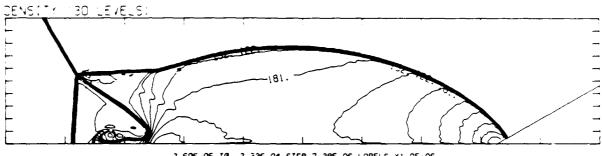


Figure 21.1b. $\theta_{u} = 29^{\circ}$, blowup-frame plots - continued.

Figure 21. Transition set 2, M_s = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=30.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 PERFECT



2.69E-05 TØ 2.33E-04 STEP 7.38E-06 LABELS X1.0E+06

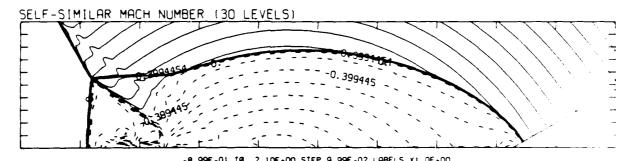
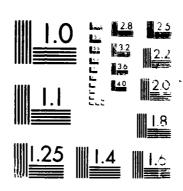


Figure 21.2a. $\theta_{\rm W} = 30^{\circ}$, whole-flowfield contour-size .

Figure 21. Transition set 2, $M_S = 4.0$, v = 1.4

A DETAILED NUMERICAL GRAPHICAL AND EXPERIMENTAL STUDY
OF OBLIQUE SHOCK MA (U) TORONTO UNIV DOMNSVIEW
(ONTRIO) INST FOR AEROSPACE STUDIES H MEDZ ET AL
81 AUG 86 UTINS-285 DNA-TR-86-365 F/G 20/4 740-A186 448 4/5 NL UNCLASSIFIED



MICROCORPORESOLUTION TEST CHARACTER NATIONAL MORES OF METAL ARE SUB-

MS= 4.00 ALP=30.00 IL=379 IR=456 UT= 74 PO=2.00E+04 PERFECT

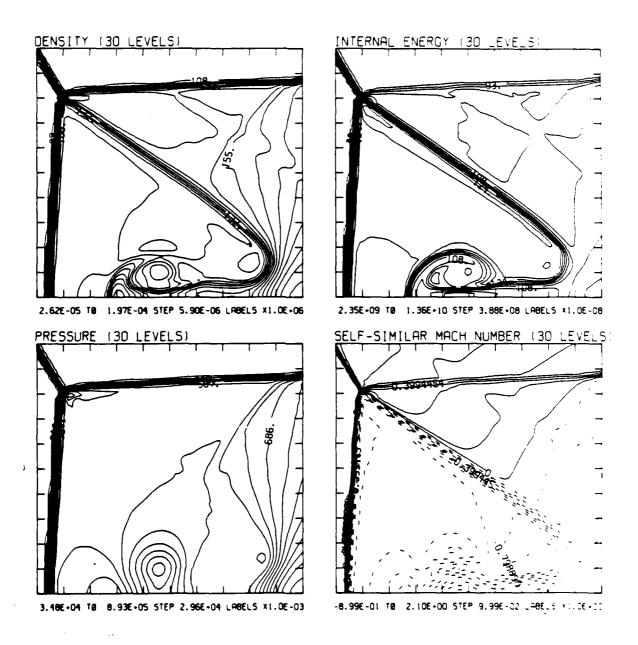


Figure 21.2b. $\theta_{\rm W}$ = 30°, blowup-frame plots.

Figure 21. Transition set 2, M_s = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=30.00 IL=379 IR=456 JT= 74 PO=2.00E+04 PERFECT

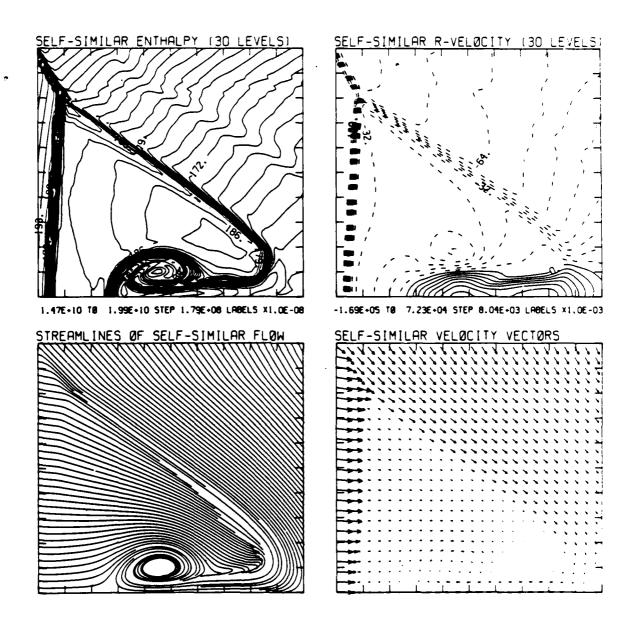
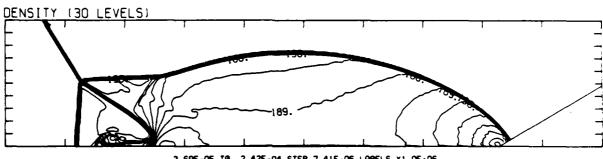


Figure 21.2b. $\theta_{w} = 30^{\circ}$, blowup-frame plots - continued.

Figure 21. Transition set 2, M_s = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=31.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 PERFECT



2.69E-05 T8 2.42E-04 STEP 7.41E-06 LABELS X1.0E+06

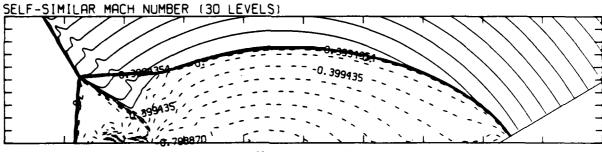


Figure 21.3a. $\theta_w = 31^\circ$, whole-flowfield contour-plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=31.00 IL=377 IR=453 JT= 73 PO=2.00E+04 PERFECT

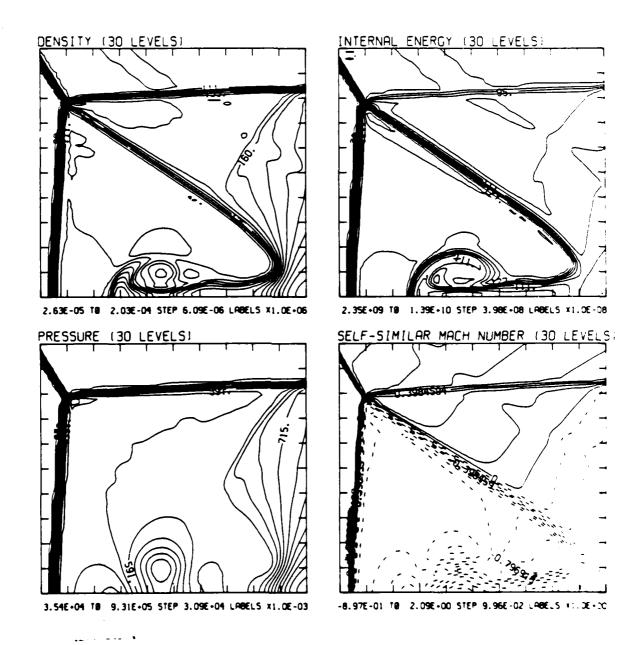


Figure 21.3b. $\theta_{\rm W}$ = 31°, blowup-frame plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=31.00 IL=377 IR=453 UT= 73 PO=2.00E+C4 PERFECT

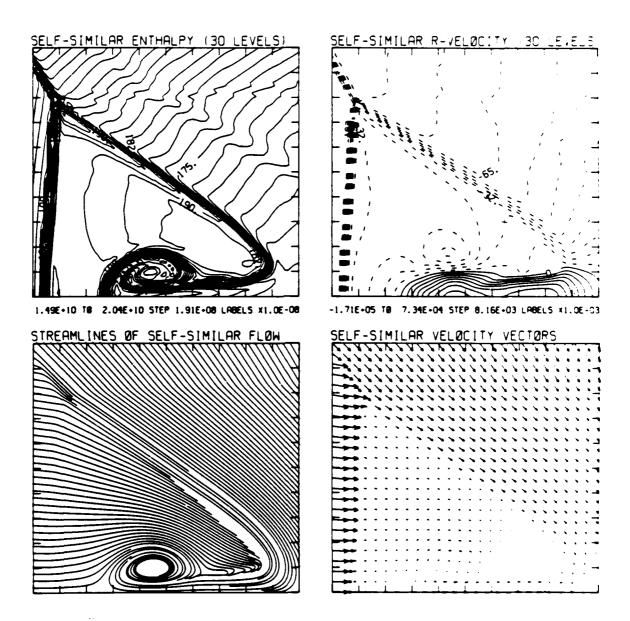
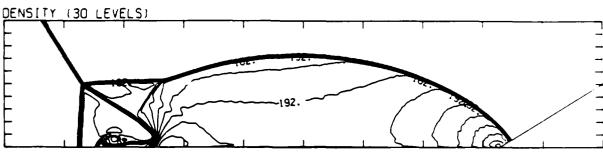


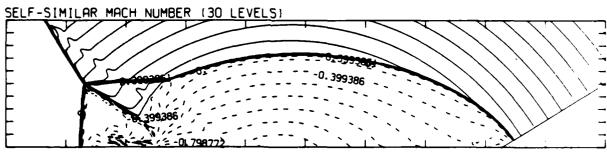
Figure 21.3b. $\theta_{\rm w}$ = 31°, blowup-frame plots - continued.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=32.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 PEFFEET



2.70E-05 TØ 2.45E-04 STEP 7.51E-06 LABELS ×1.0E+06



-8.99E-01 T8 2.10E+00 STEP 9.98E-02 LABELS X1.0E+00

Figure 21.4a. $\theta_{\rm W}$ = 32°, whole-flowfield contour-plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=32.00 IL=375 IR=450 UT= T2 P0=2.00E+04 PERFECT

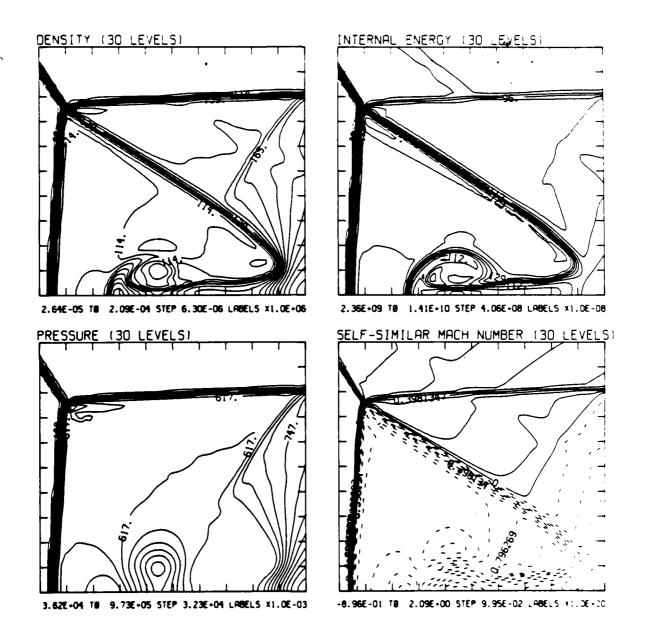


Figure 21.4b. $\theta_w = 32^{\circ}$, blowup-frame plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=32.00 [L=375 IR=450 UT= T2 PO=2.00E+04 PERFECT

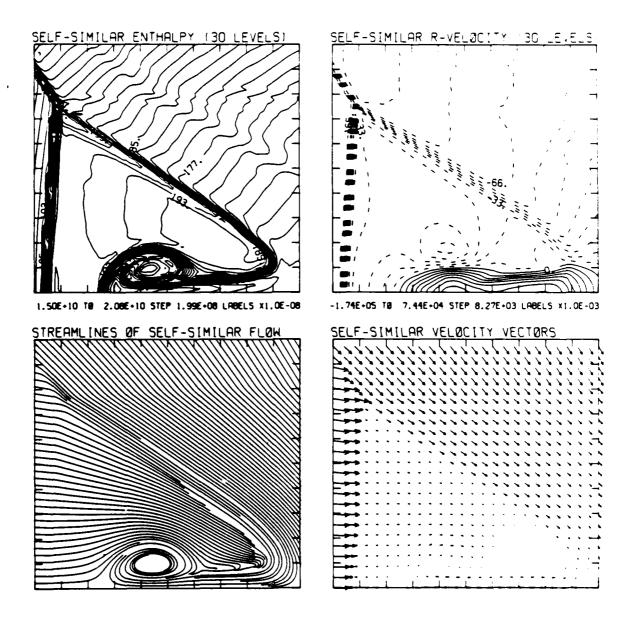
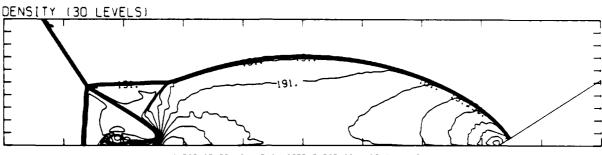


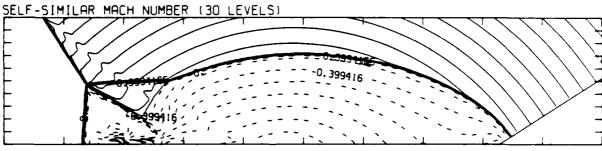
Figure 21.4b. $\theta_{\rm w}$ = 32°, blowup-frame plots - continued.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=33.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 PEFFECT



2.70E-05 TØ 2.44E-04 STEP 7.50E-06 LABELS X1.0E+06



-8.99E-01 TØ 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 21.5a. $\theta_{\rm w}$ = 33°, whole-flowfield contour-plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=33.00 IL=371 IR=446 UT= T2 P0=2.30E+34 PERFECT

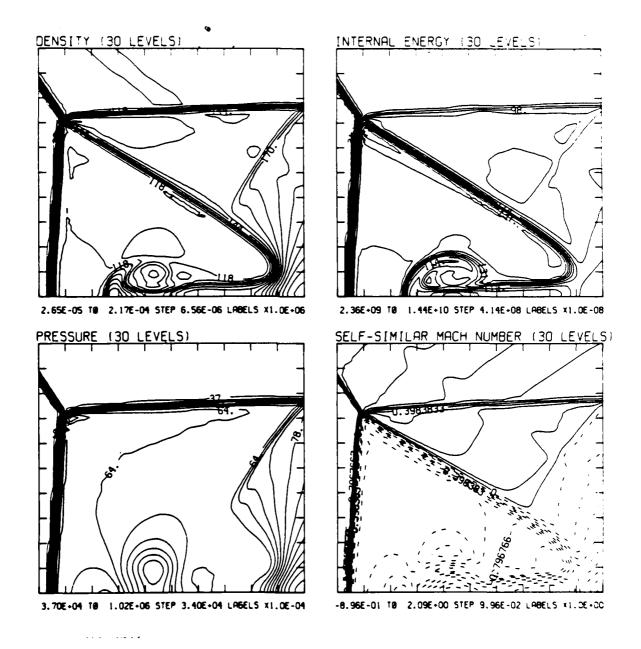


Figure 21.5b. $\theta_{\rm W}$ = 33°, blowup-frame plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=33.00 IL=37! IR=446 UT= 72 PG=2.00E+84 PERFEST

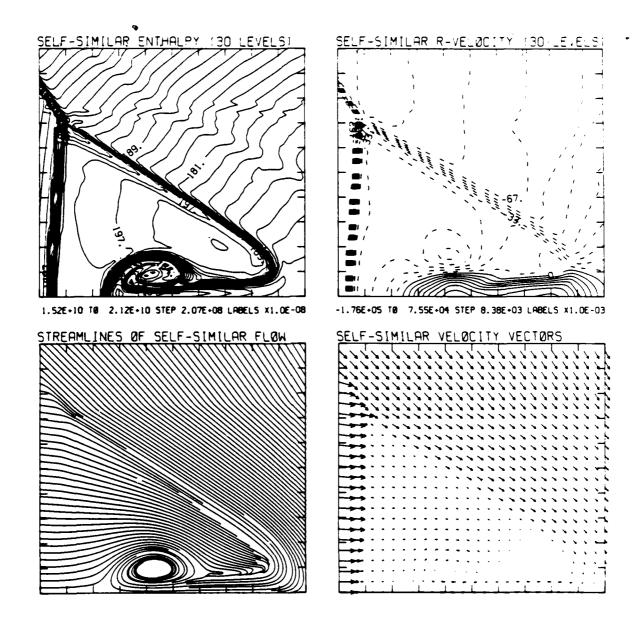
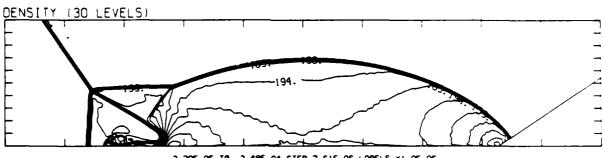


Figure 21.5b. $\theta_{\rm W}$ = 33°, blowup-frame plots - continued.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=34.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 PERFECT



2.70E-05 TO 2.48E-04 STEP 7.61E-06 LABELS X1.0E+06

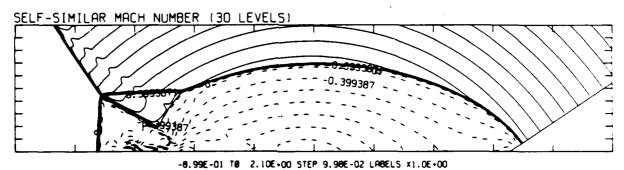


Figure 21.6a. $\theta_w = 34^{\circ}$, whole-flowfield contour-plots.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=34.00 IL=369 IR=443 UT= 71 PO=2.00E+04 PERFECT

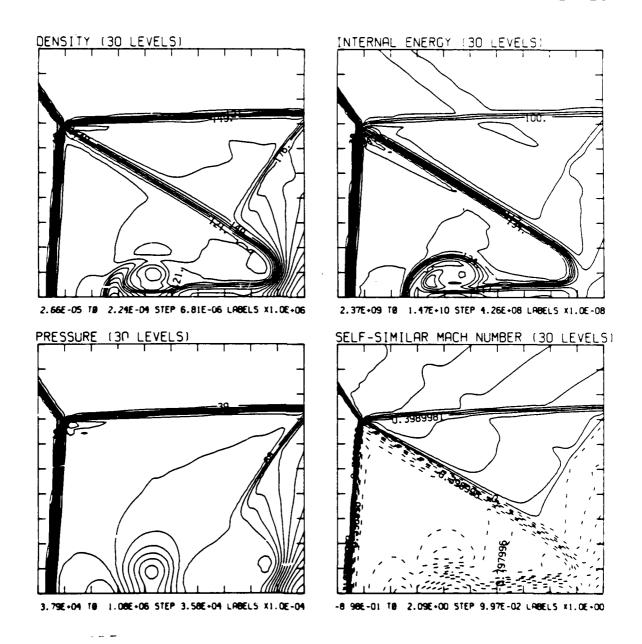


Figure 21.6b. $\theta_w = 34^{\circ}$, blowup-frame plots.

Figure 21. Transition set 2, M_s = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=34.00 IL=369 IR=443 UT= 71 PO=2.00E+34 PERFECT

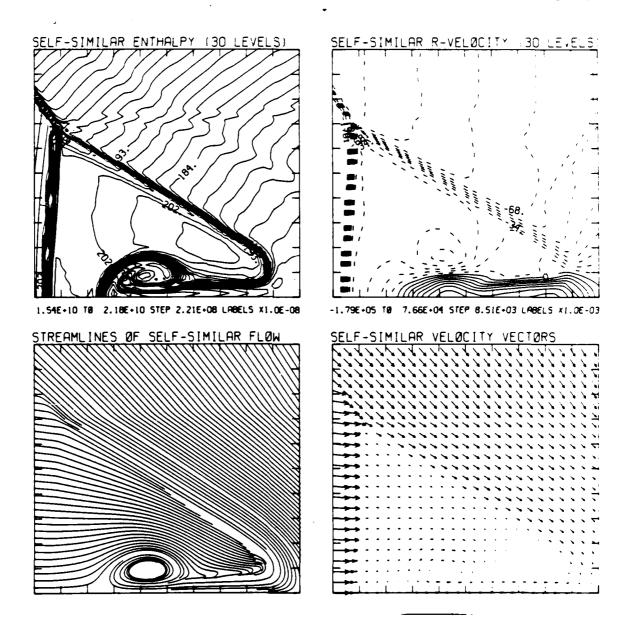
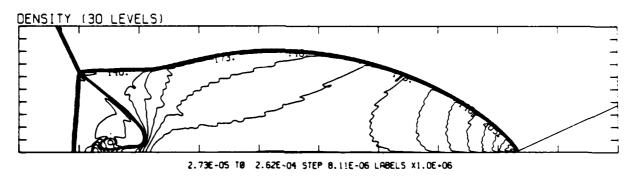


Figure 21.6b. $\theta_{\omega} = 34^{\circ}$, blowup-frame plots - continued.

Figure 21. Transition set 2, M_S = 4.0, γ = 1.4 - continued.

MS= 4.00 ALP=25.00 NR=510 NZ=110 KBEG= 90 P0=2.00E+04 HANSEN



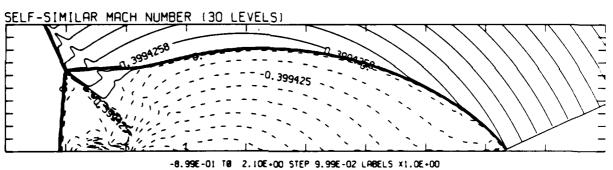


Figure 22.1a. $\theta_{\rm W}$ = 25°, whole-flowfield contour-plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen.

MS= 4.00 ALP=25.00 IL=388 [R=467 UT= 76 PO=2.00E+04 HANSEN

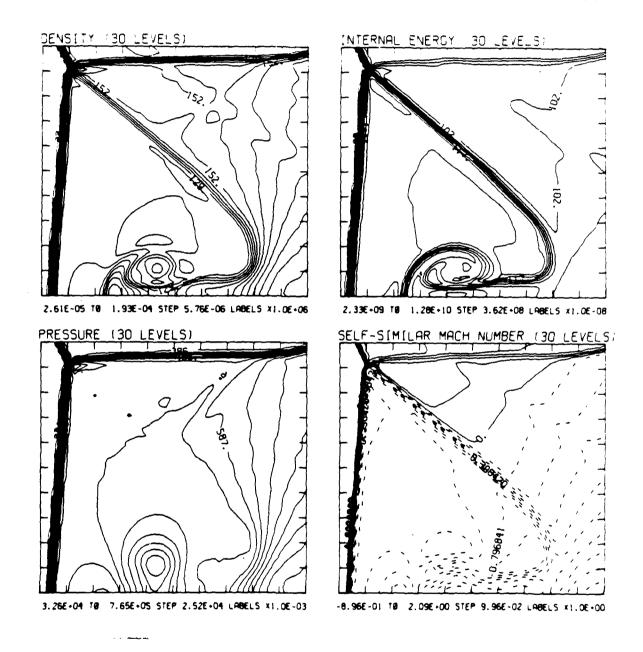


Figure 22.1b. $\theta_{\rm w}$ = 25°, blowup-frame plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=25.00 [L=388 [R=467 LT= 76 PO=2.00E+04 HANSEN

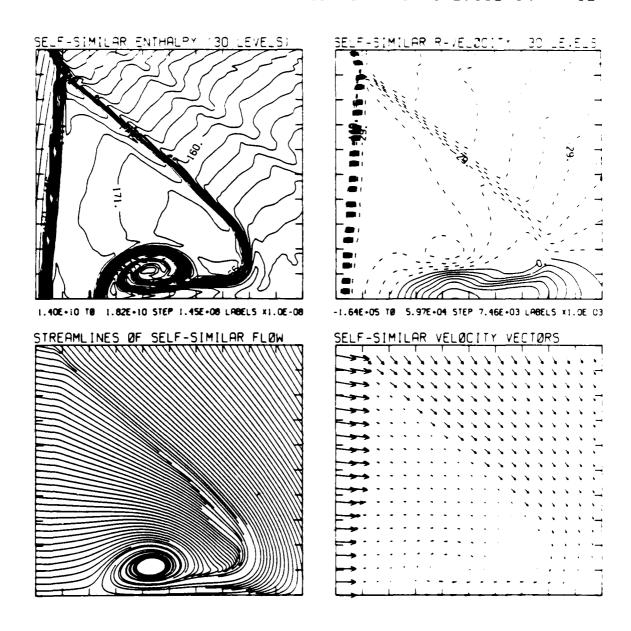
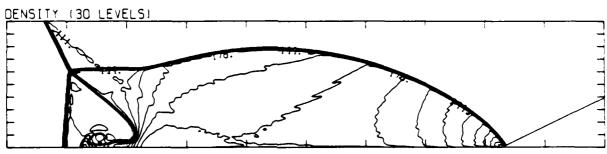


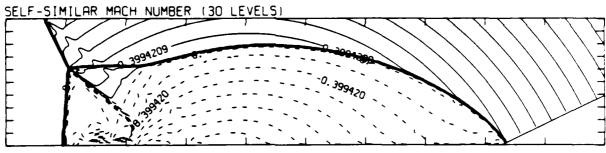
Figure 22.1b. $\theta_{\rm W}$ = 25°, blowup-frame plots - continued.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=26.00 NR=510 NZ=110 KBEG= 90 PG=2.00E+04 HANSEN



2.74E-05 TØ 2.71E-04 STEP 8.39E-06 LRBELS X1.0E+06



-8.99E-01 TB 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 22.2a. q = 26°, whole-flowfield contour-plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=26.00 [L=387 [R=465 JT= 75 PO=2.00E+04 HANSEN

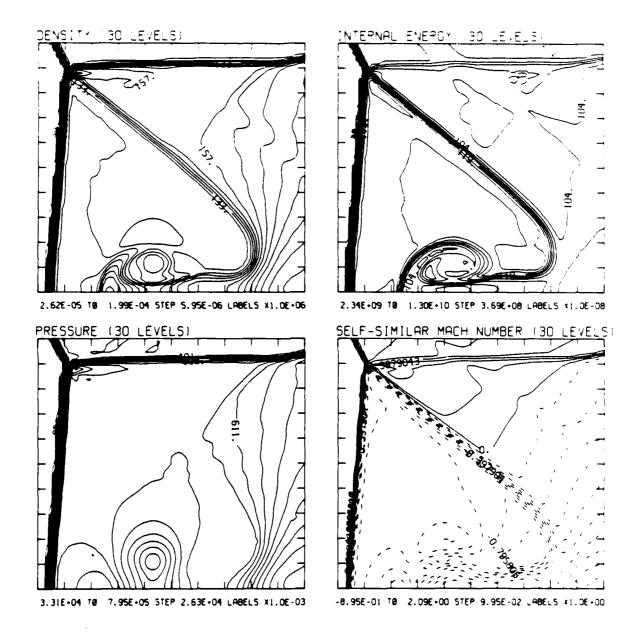


Figure 22.2b. $\theta_w = 26^{\circ}$, blowup-frame plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=26.00 [L=387 [R=465 LT= 75 PO=2.00F+04 HANSEN

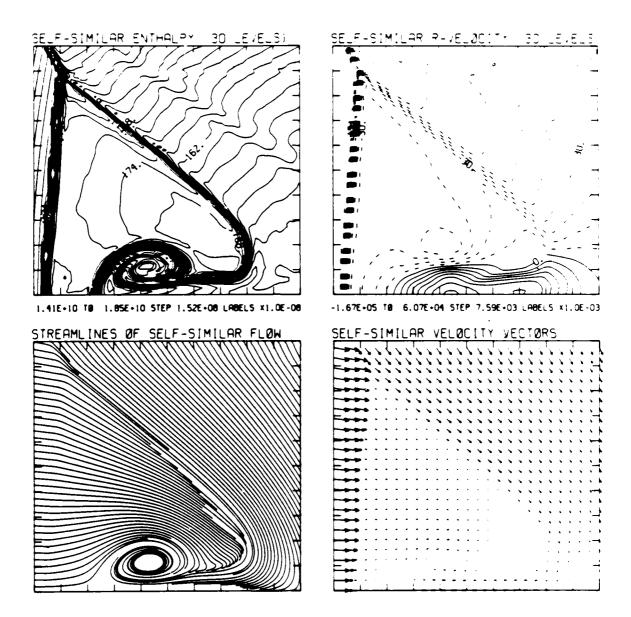
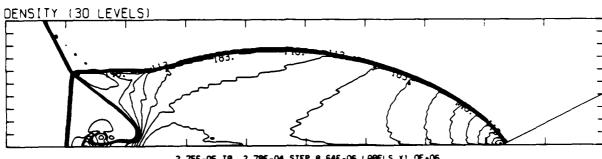


Figure 22.2b. $\theta_w = 26^{\circ}$, blowup-frame plots - continued.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=27.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+34 HANSEN



2.75E-05 TØ 2.78E-04 STEP 8.64E-06 LABELS X1.0E+06

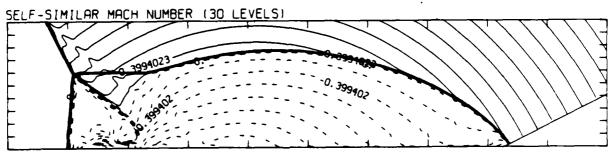


Figure 22.3a. $\theta_{\rm w}$ = 27°, whole-flowfield contour-plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.30 ALP=27.00 [L=384 [R=462 UT= 75 PO=2.30E+34 HANSEN

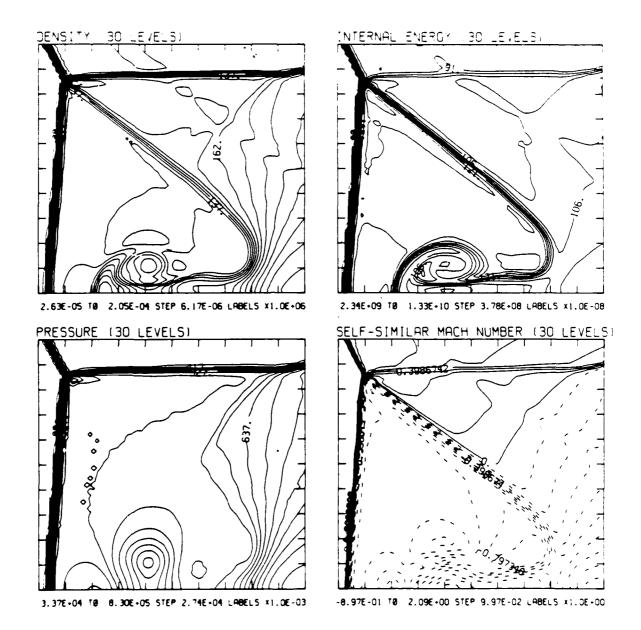


Figure 22.3b. $\theta_{\rm W}$ = 27°, blowup-frame plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=27.00 (L=384 (F=462 UT= 75 P0=2.00E+34 HANSE).

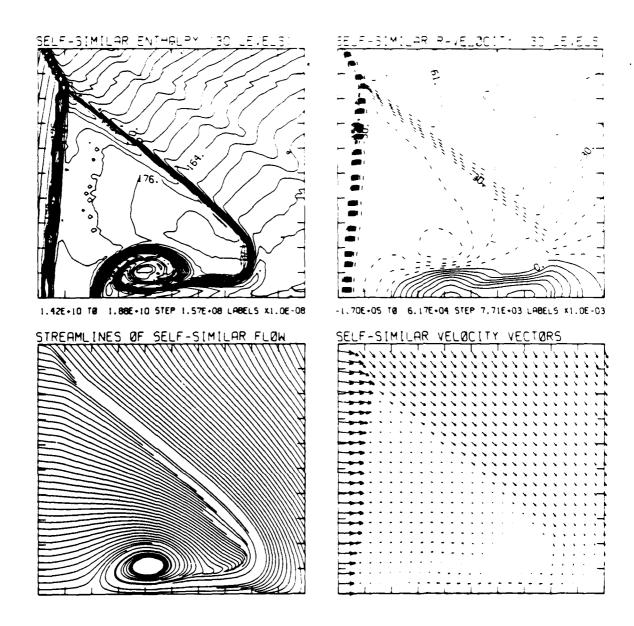
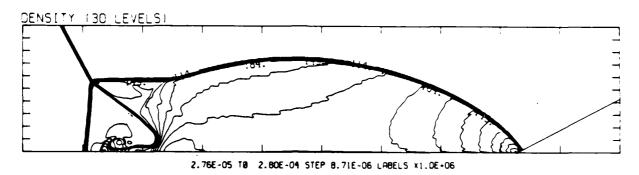


Figure 22.3b. $\theta_{\rm w}$ = 27°, blowup-frame plots - continued.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=28.00 NR=510 NZ=110 KBEG= 90 PC=2.00E+04 HANSEN



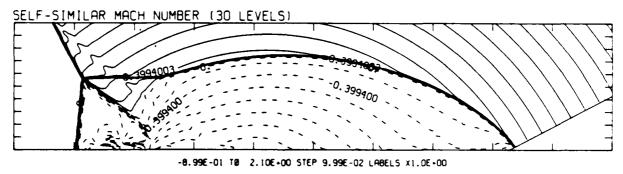


Figure 22.4a. $\theta_{\rm W}$ = 28°, whole-flowfield contour-plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=28.00 [L=383 [R=460 UT= T4 PC=2.00E+04 HANSEN

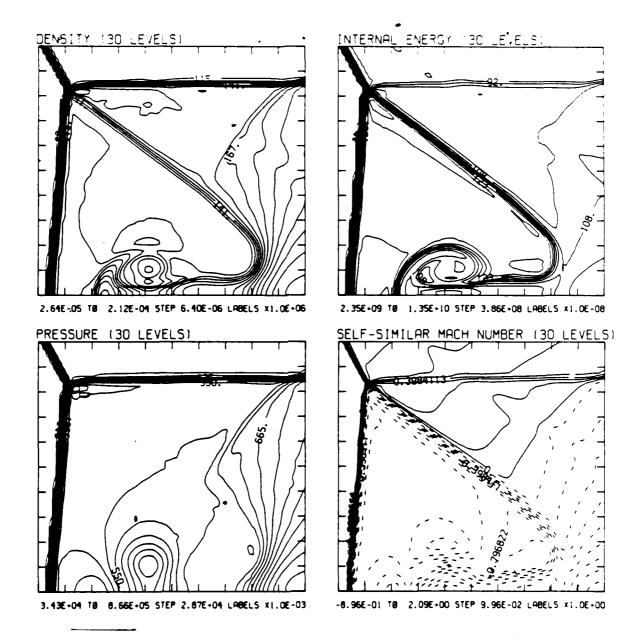


Figure 22.4b. $\theta_{\rm W}$ = 28°, blowup-frame plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=28.00 IL=383 IR=460 UT= T4 PC=2.00E+04 HANSEN

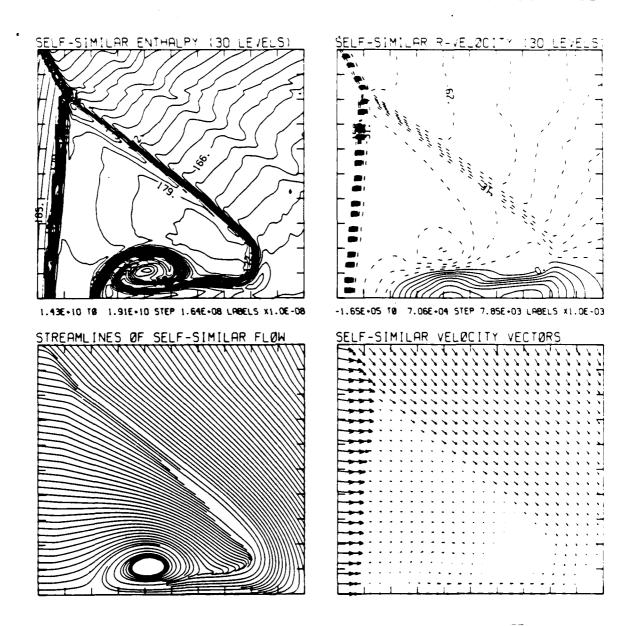
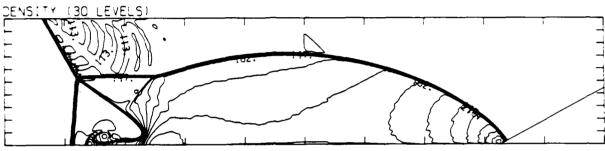


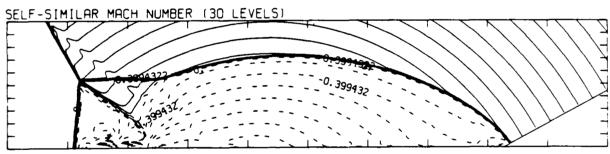
Figure 22.4b. $\theta_{\rm W}$ = 28°, blowup-frame plots - continued.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=29.00 NR=510 NZ=110 KBEG= 90 PC=2.00E+04 HANSEN



2.75E-05 TØ 2.77E-04 STEP 8.59E-06 LABELS X1.0E+06



-8.99E-01 TØ 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 22.5a. $\theta_{\rm w}$ = 29°, whole-flowfield contour-plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=29.00 IL=379 [R=456 UT= T4 P0=2.00E+04 HANSEN

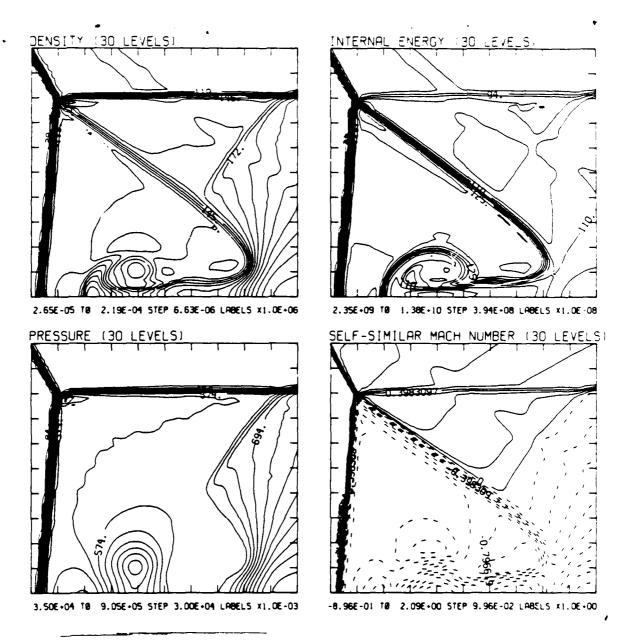


Figure 22.5b. $\theta_{\rm W}$ = 29°, blowup-frame plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=29.00 IL=379 IR=456 UT= 74 PO=2.00E+04 HANSEN

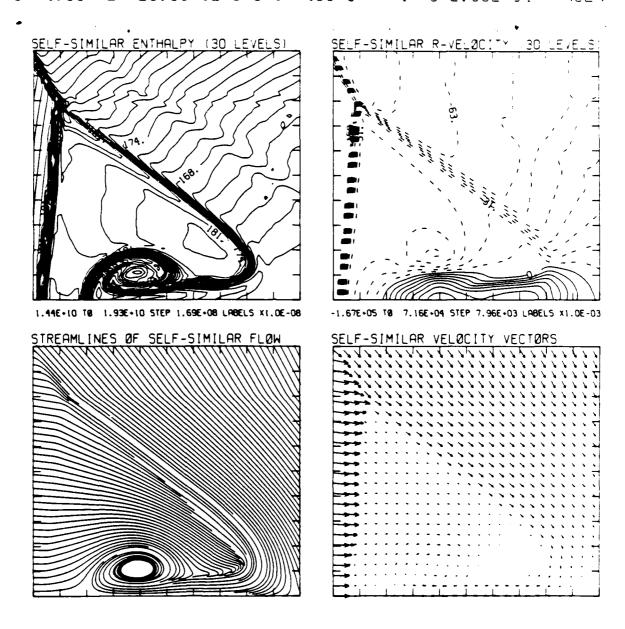
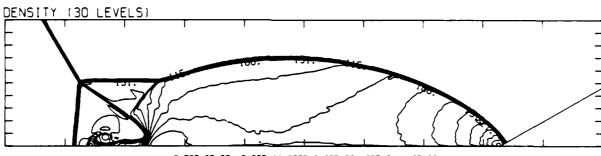


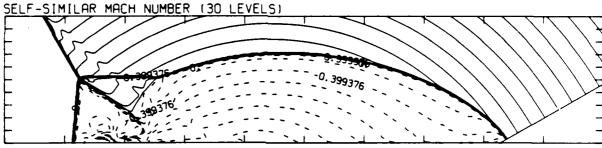
Figure 22.5b. $\theta_{\rm W}$ = 29°, blowup-frame plots - continued.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=30.00 NR=510 NZ=110 KBEG= 90 PO=2.00E+04 HANSEN



2.76E-05 T0 2.83E-04 STEP 8.82E-06 LABELS X1.0E+06



-8.99E-01 TØ 2.10E+00 STEP 9.98E-02 LABELS X1.0E+00

Figure 22.6a. $\theta_{\rm W}$ = 30°, whole-flowfield contour-plots.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 4.00 ALP=30.00 [L=379 [R=455 UT= 73 P0=2.00E+04 H4N35]

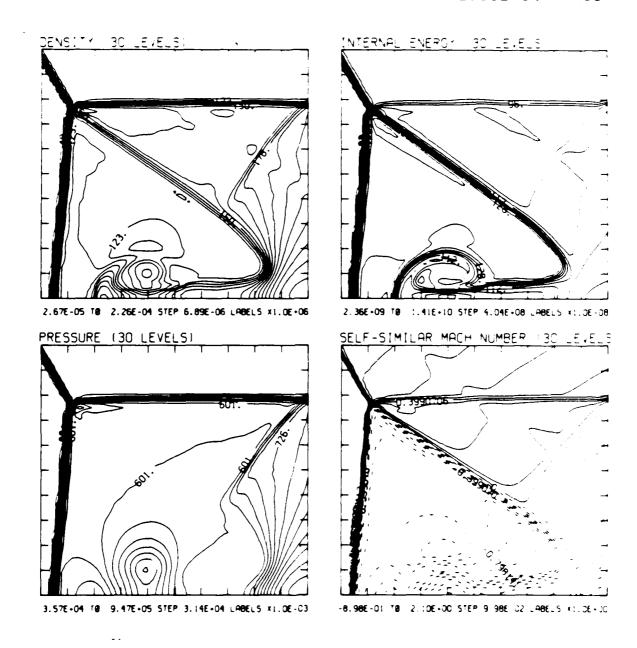


Figure 22.6b. $\theta_{\rm w}$ = 30°, blowup-frame plots.

Figure 22. Transition set 2, $M_S = 4.0$, Hansen - continued.

MS= 4.00 ALP=30.00 (L=379 (P=455 UT= 73 P0=2.00E+04 HANSE).

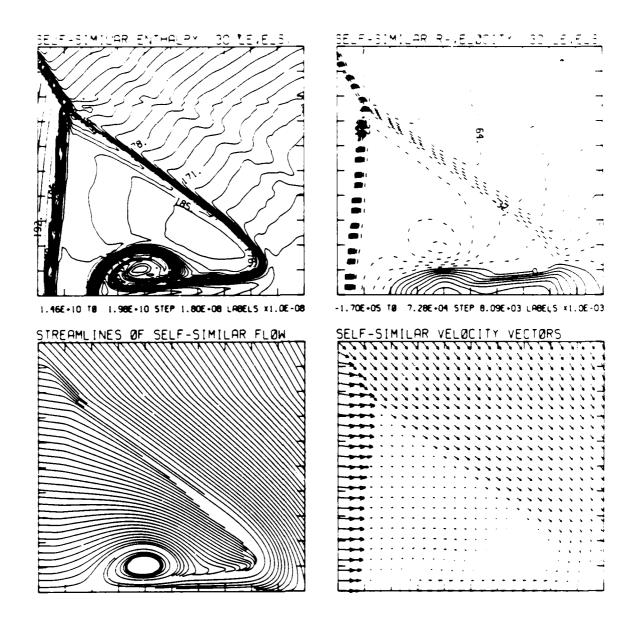
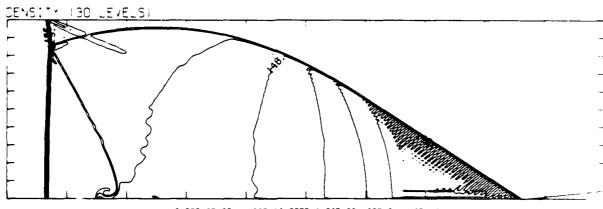


Figure 22.6b. $\theta_{\rm W}$ = 30°, blowup-frame plots - continued.

Figure 22. Transition set 2, M_S = 4.0, Hansen - continued.

MS= 8.75 ALP= 6.00 NR=525 NZ=160 KBEG= 75 PC=2.00E+34 PERFETT



2.56E-05 TØ 1.63E-04 STEP 4.74E-06 LABELS x1.0E+06

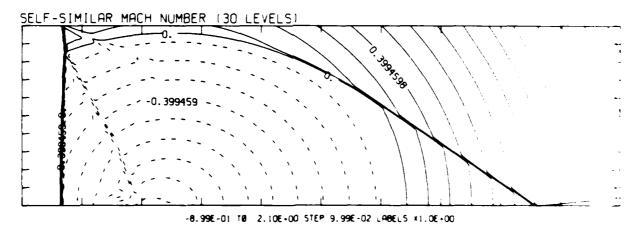


Figure 23.1a. $\theta_w = 6^\circ$, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4.

MS= 8.75 ALP= 8.00 (L=332 (F=494 UT=155 P0=2.008+04 PE4PE):

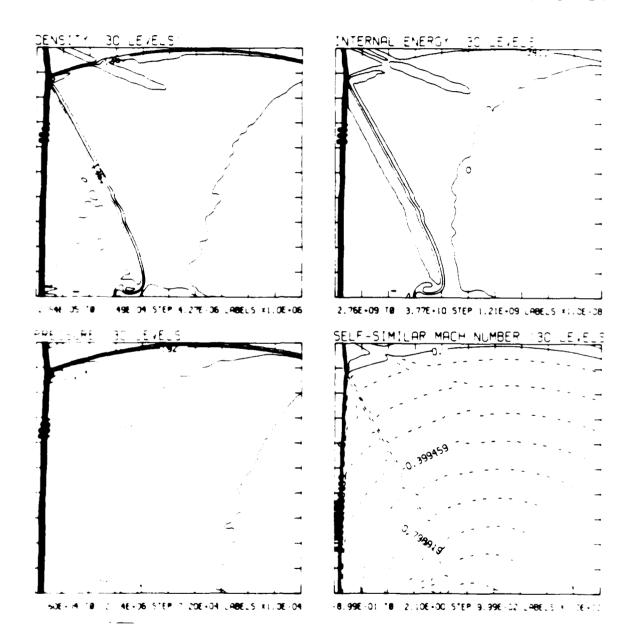


Figure 23.1b. $\theta_a = 6^\circ$, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP= 8.00 (L=330 (P=434)7=155 PD=7.77P.74 PP=4.4

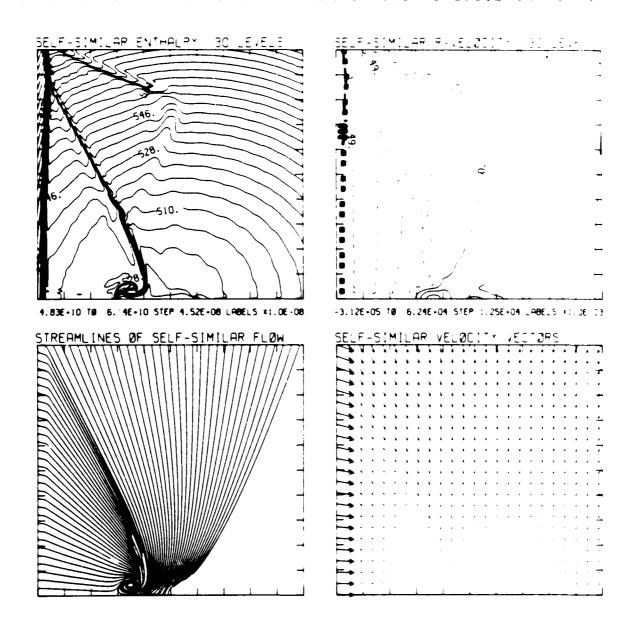
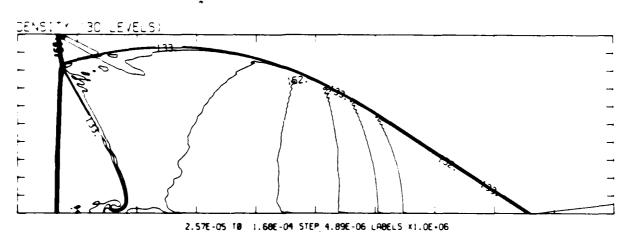


Figure 23.1b. $\theta_u = 6^\circ$, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 3.75 ALP= 7.00 NR=525 NZ=160 KBEG= 75 PC=2.00E+04 PERFECT



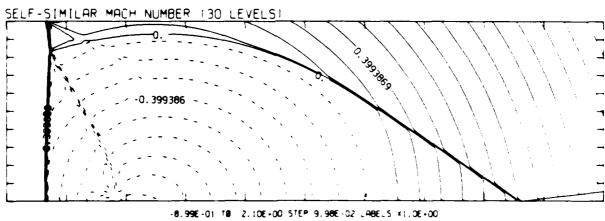


Figure 23.2a. $\theta_{\rm w}$ = 7°, whole-flowfield contour-plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP= 7.00 [L=332 [R=494 UT=155 PO=2.00E+04 PERFECT

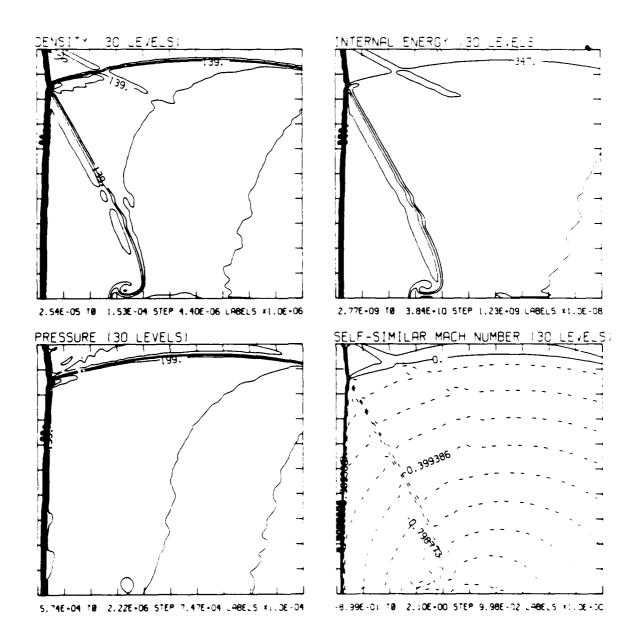


Figure 23.2b. $\theta_w = 7^{\circ}$, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALR= 7.00 (L=330 (R=494 UT=155 PD=0.00E+04 PERFECT

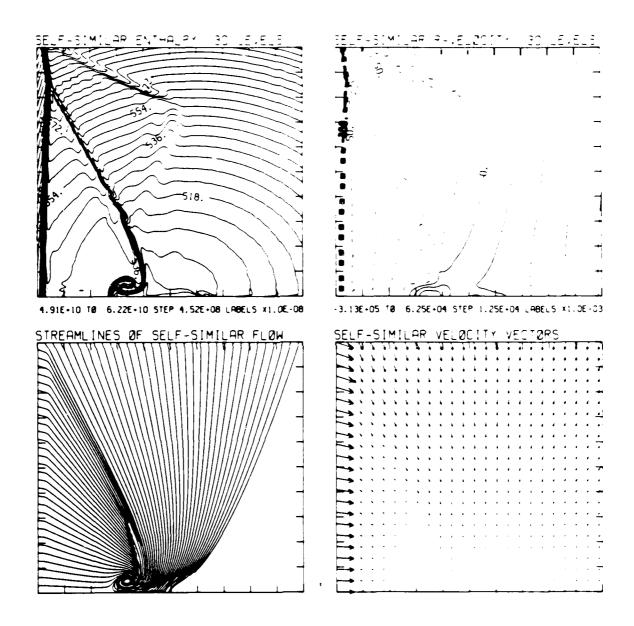
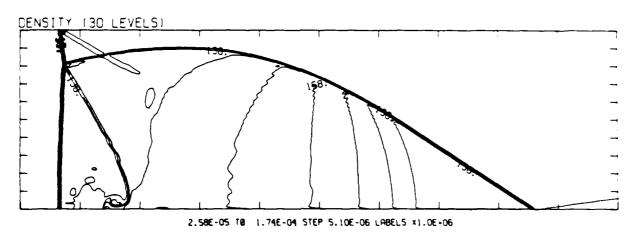


Figure 23.2b. $\theta_{w} = 7^{\circ}$, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 8.30 NR=525 NZ=160 KBEG= 75 PC=2.00E+04 FERFEDT



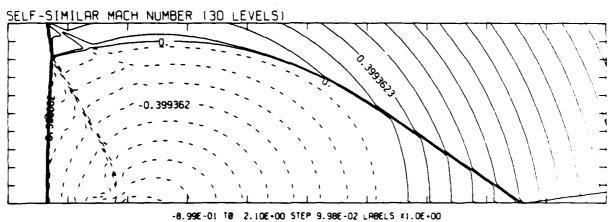


Figure 23.3a. $\theta_{\rm w}$ = 8°, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 3.75 ALP= 3.00 [L=332 [R=494 UT=155 PC=2.00E+04 PEFFE0T

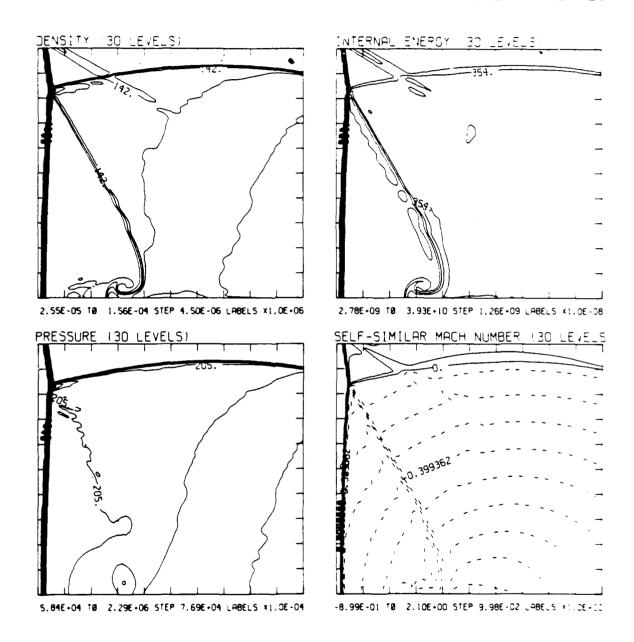


Figure 23.3b. $\theta_{\rm W}$ = 8°, blowup-frame plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 8.30 (L=332 [R=494 UT=155 PO=2.88E+34 PERFE]

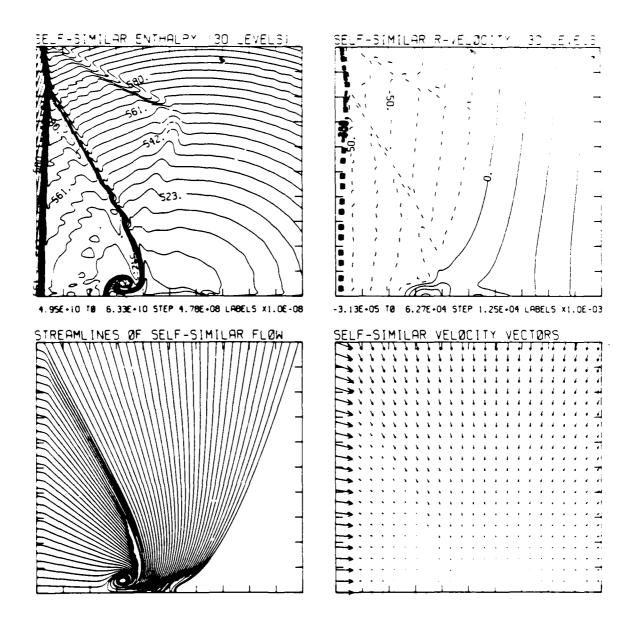
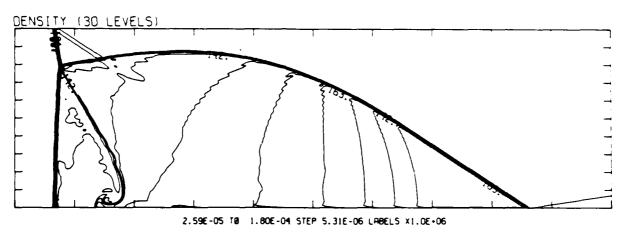


Figure 23.3b. $\theta_{\rm W}$ = 8°, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 9.00 NR=525 NZ=160 KBEG= 75 PG=2.30E+04 PERFECT



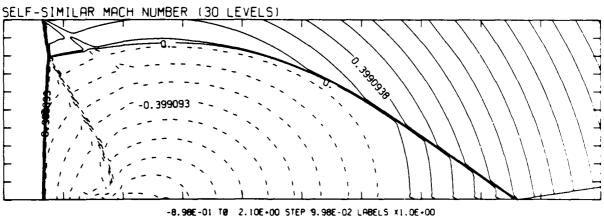


Figure 23.4a. $\theta_{\omega} = 9^{\circ}$, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 9.00 IL=332 IP=494 UT=155 PC=2.00E+04 PERFECT

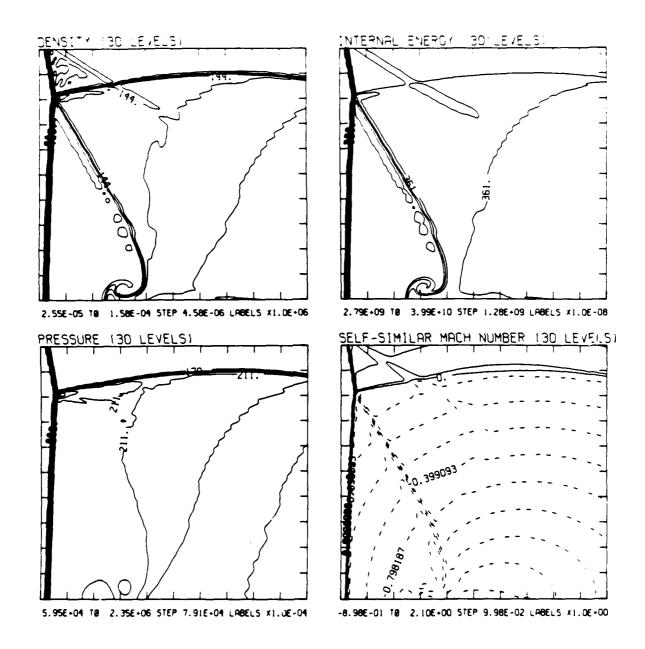


Figure 23.4b. $\theta_{\rm W}$ = 9°, blowup-frame plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 9.00 [L=332 [R=494 UT=155 PC=2.00E+04 PERFECT

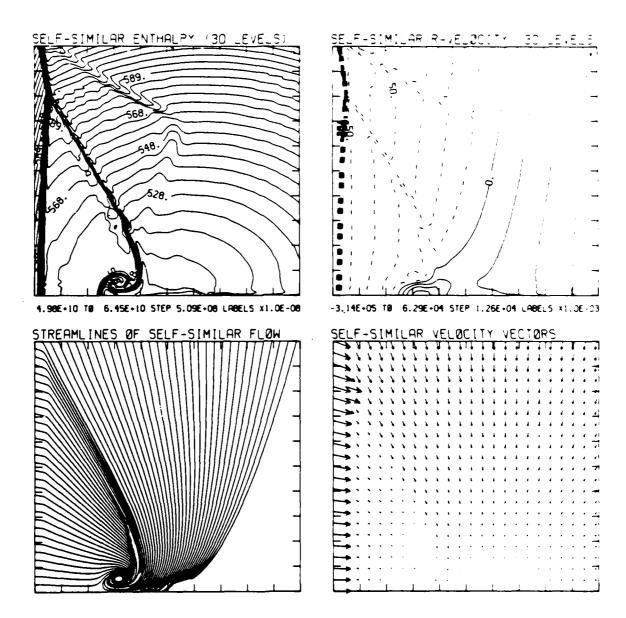
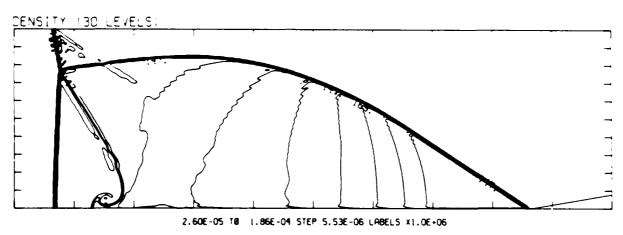


Figure 23.4b. $\theta_w = 9^{\circ}$, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALR=10.00 NR=525 NZ=160 KBE0= 75 P0=0.00E+04 PERFEDT



SELF-SIMILAR MACH NUMBER (30 LEVELS)

-8.97E-01 TO 2.09E+00 STEP 9.97E-02 LABELS X1.0E+00

Figure 23.5a. $\theta_{\rm w}$ = 10°, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=10.00 (L=332 (F=494 LT=155 R0=2.00E+24 REFEEDT

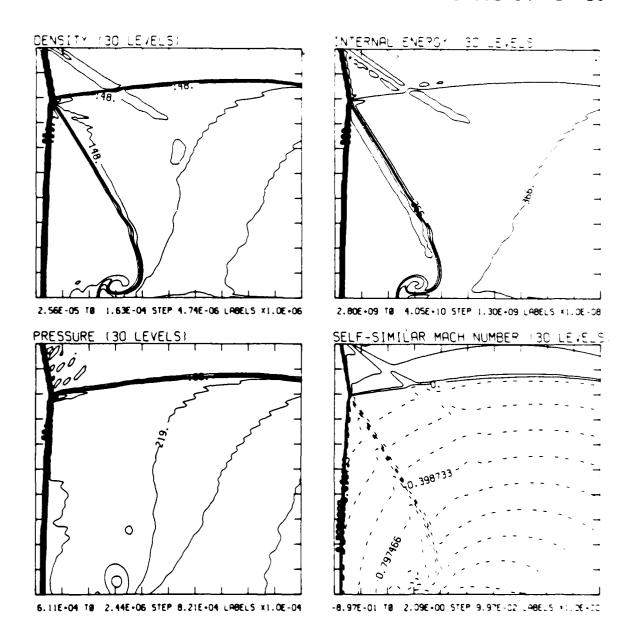


Figure 23.5b. $\theta_{\rm w} = 10^{\circ}$, blowup-frame plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=10.00 [L=332 [R=494 UT=155 PD=2.00E+24 PERFECT

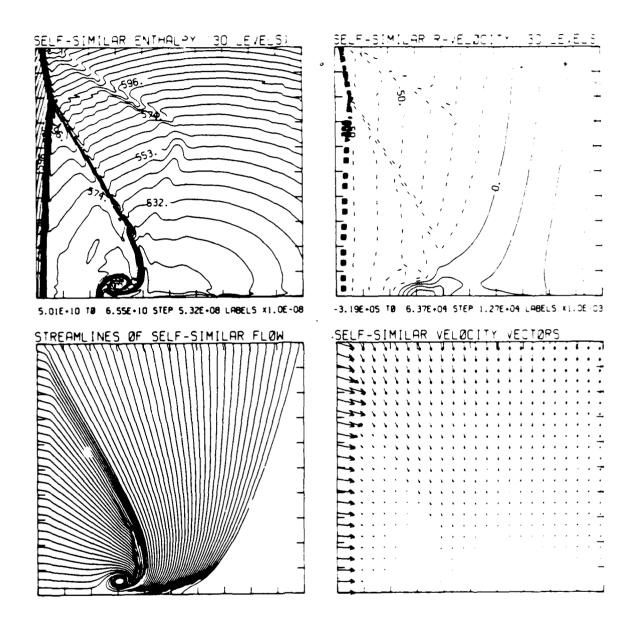
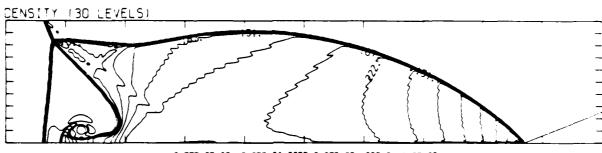


Figure 23.5b. $\theta_{\rm W}$ = 10°, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=22.00 NR=550 NZ=115 KBEG= 75 PC=2.00E+04 FEFFET



2.77E-05 TO 2.85E-04 STEP 8.87E-06 LABELS X1.0E+06

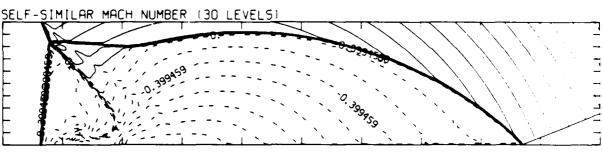


Figure 23.6a. $\theta_{\rm W}$ = 22°, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 ~ continued.

MS= 8.75 ALP=22.30 (L=42) (P=5)9 LT=)(0 P0=0.77F474 AFSF

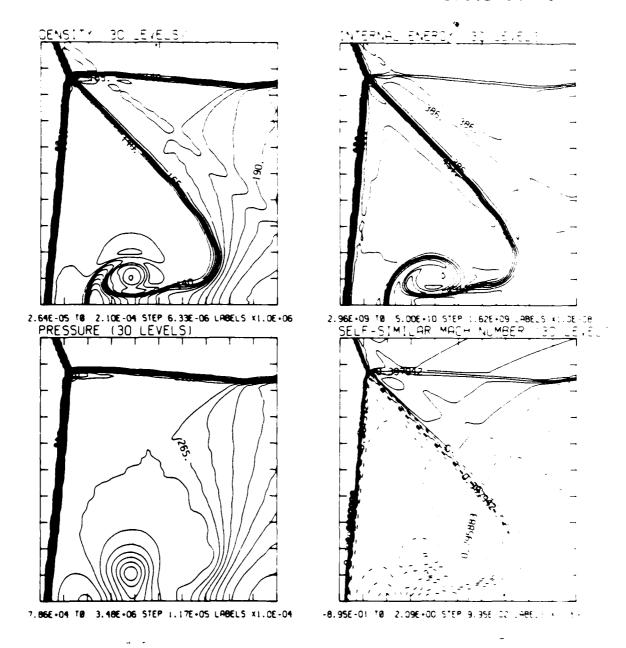


Figure 23.6b. $\theta_w = 22^{\circ}$, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MGE ALTS A PERTLOS I E48, TRESTA TENT PRETITABLOS PARARITA

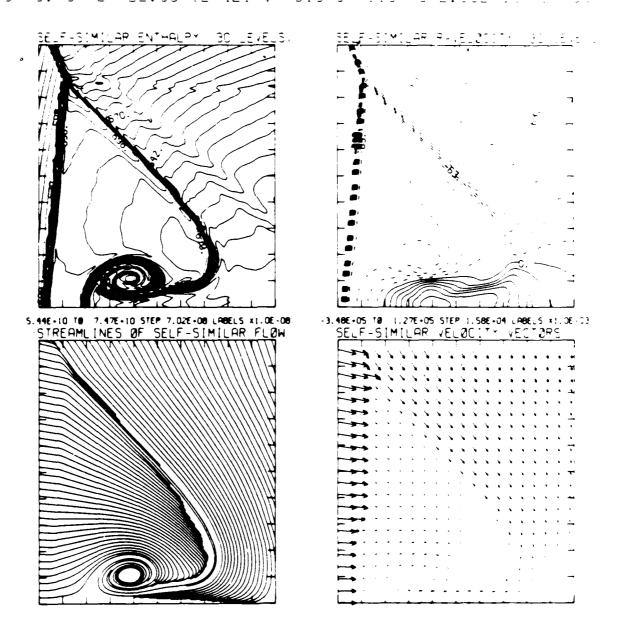


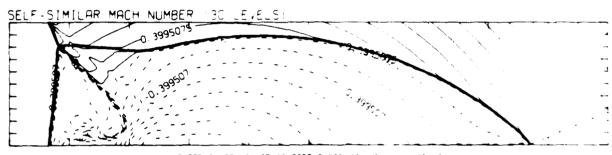
Figure 23.6b. $\theta_{\rm w}$ = 22°, blowup-frame plots - continued.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=23.00 NP=550 NZ=115 MBEG= 75 PG=2.00E+04 PEFFE 17



2.796-05 Ta 2.996-04 STEP 9.346-06 JABELS 41.06+06



-8.99E-01 18 2.10E+00 S1EP 9.99E+02 148Ept #1.0E+00

Figure 23.7a. $9_w = 23^\circ$, whole-flowfield contour-plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=28.00 (L=42: [A=5:8 UT=::0 P0=2.005+04 PEFFE:T

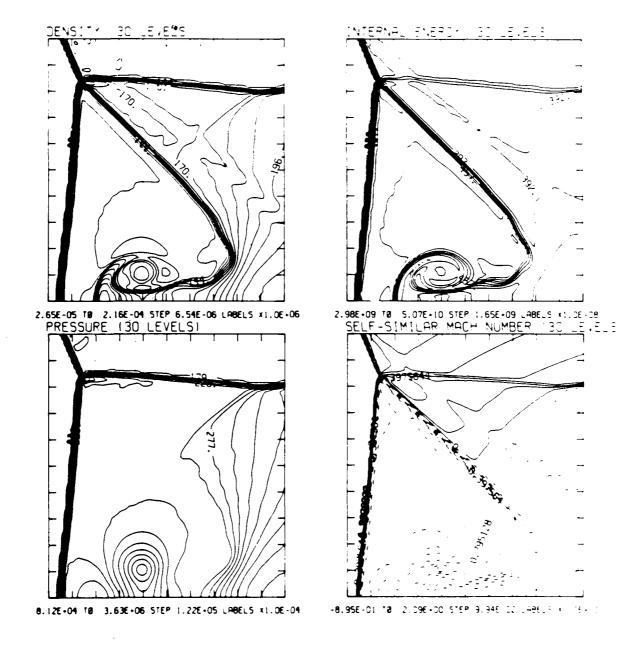


Figure 23.7b. $\theta_{\rm w}$ = 23°, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=23.00 (L=42) (P=518 UT=110 F0=2.00E+04 PEFFE)

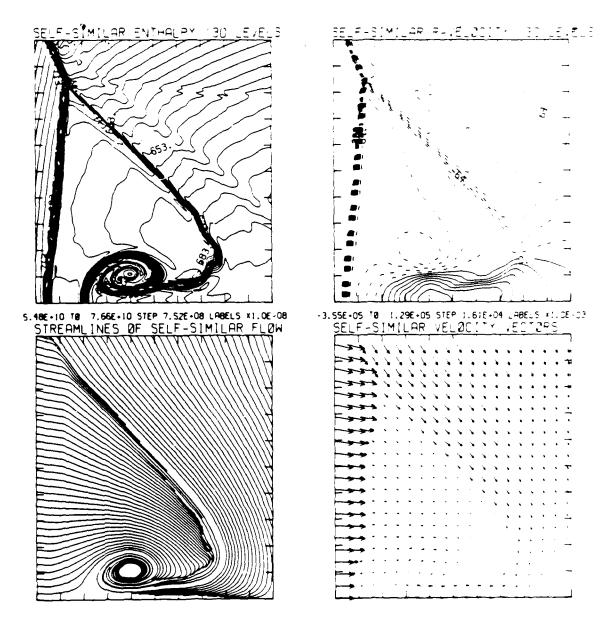
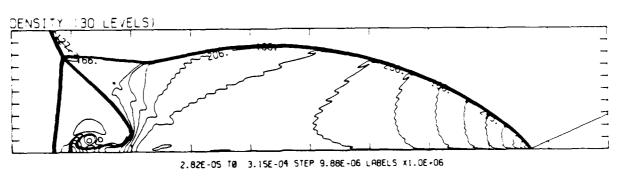
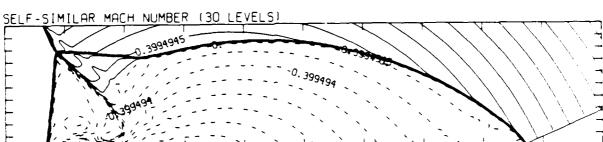


Figure 23.7b. $\theta_w = 23^{\circ}$, blowup-frame plots - continued.

Figure 23. Transition set 3, $M_S = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=24.00 NR=550 NZ=115 KBEG= 75 PG=2.00E+04 PEFFED





-8.99E-01 T0 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 23.8a. $\theta_{\rm w}$ = 24°, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

M5= 3.75 ALP=04.00 (L=40) (P=518 UT=110 P0=0.00E+04 PERFECT

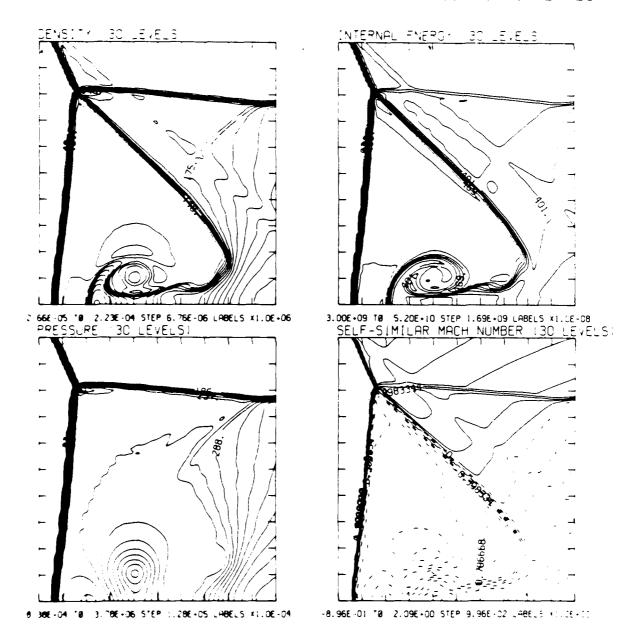


Figure 23.8b. $\theta_{\rm w} = 24^{\circ}$, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=24.00 [L=421 [R=518 UT=110 PC=2.00E+04 PEFFE]]

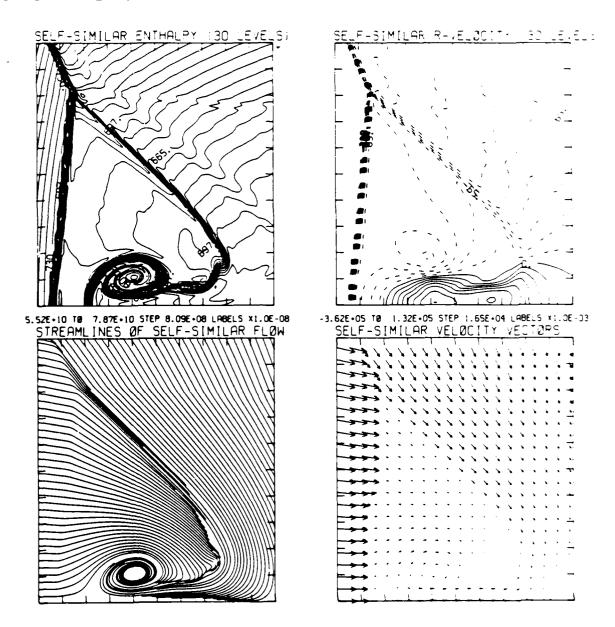
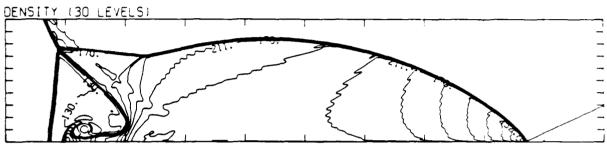


Figure 23.8b. θ_{ω} = 24°, blowup-frame plots - continued.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=25.00 NR=550 NZ=115 KBEG= 75 PG=2.00E+34 PERFECT



2.83E-05 TØ 3.23E-04 STEP 1.02E-05 LABELS X1.0E+06

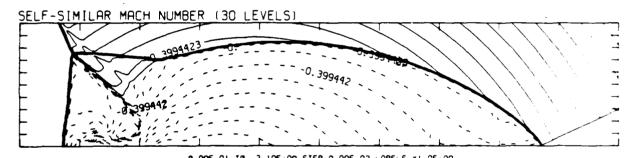


Figure 23.9a. $\theta_w = 25^{\circ}$, whole-flowfield contour-plots.

Figure 23. Transition set 3, M_s = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=25.00 [L=42] [R=518 UT=110 PO=2.00E+34 FEFFE]T

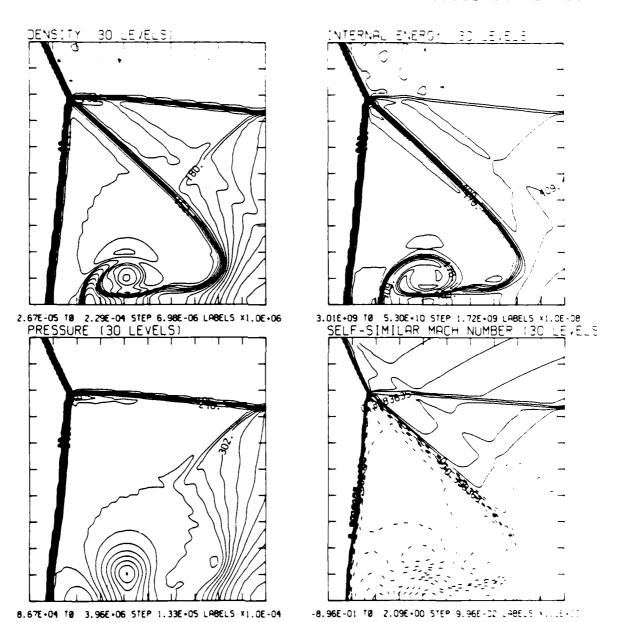


Figure 23.9b. $\theta_{\rm w}$ = 25°, blowup-frame plots.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=25.00 [L=421 [R=518 UT=110 PG=2.00E+04 PERFECT

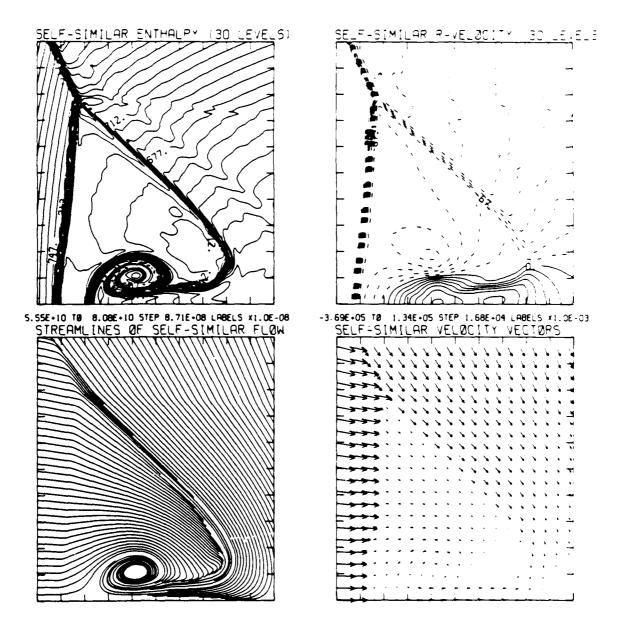
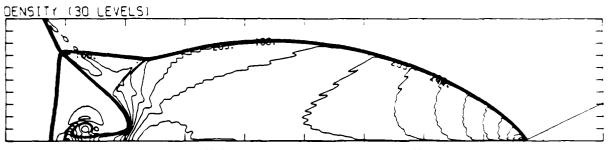


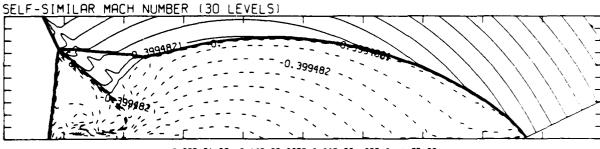
Figure 23.9b. $\theta_a = 25^{\circ}$, blowup-frame plots - continued.

Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP=26.00 NR=550 NZ=115 KBEG= 75 PD=2.00E+04 PERFEDT



2.82E-05 T0 3.14E-04 STEP 9.85E-06 LABELS x1.0E+06



-8.99E-01 TØ 2.10E+00 STEP 9.99E-02 LABELS X1.0E+00

Figure 23.10a. $\theta_{\rm w}$ = 26°, whole-flowfield contour-plots.

Figure 23. Transition set 3, $M_S = 8.75$, $\gamma = 1.4$ - continued.

MS= 3.75 ALP=26.00 [L=42] [R=5:8 UT=::0 PO=2.00E+04 PERFECT

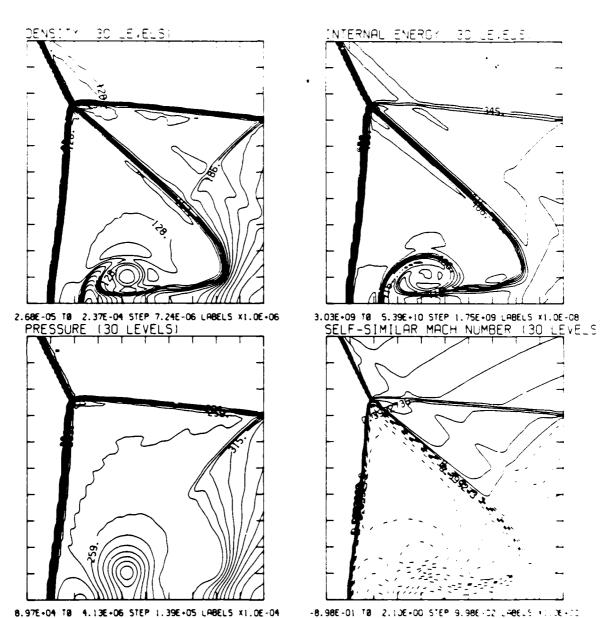


Figure 23.10b. $\theta_{\text{W}} = 26^{\circ}$, blowup-frame plots.

Figure 23. Transition set 3, $M_s = 8.75$, $\gamma = 1.4$ - continued.

MS= 8.75 ALP=26.00 [L=42] [R=518 UT=110 PC=2.00E+04 FERFE07

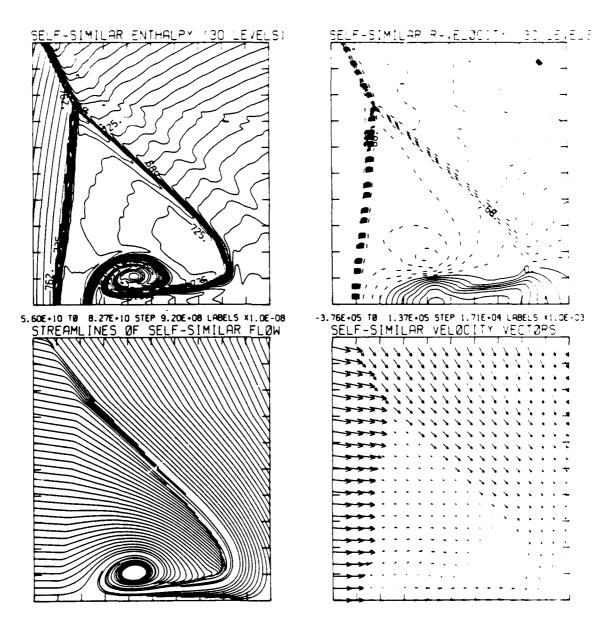
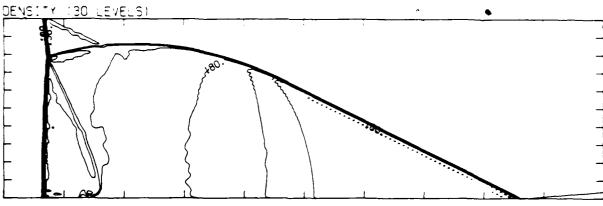


Figure 23.10b. θ_W = 26°, blowup-frame plots - continued. Figure 23. Transition set 3, M_S = 8.75, γ = 1.4 - continued.

MS= 8.75 ALP= 5.00 NR=525 NZ=180 KBEG= 75 PG=2.00E+04 HANSEN



2.62E-05 TØ 1.98E-04 STEP 5.93E-06 LABELS X1.0E+06

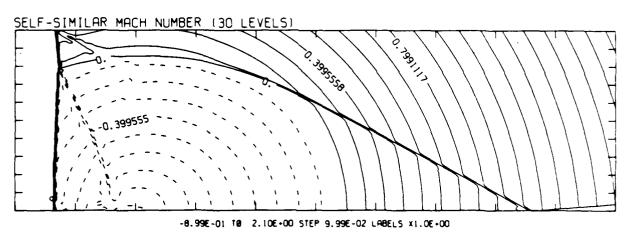


Figure 24.1a. $\theta_{\rm W}$ = 5°, whole-flowfield contour-plots.

Figure 24. Transition set 3, M_S = 8.75, Hansen.

MS= 8.75 ALP= 5.88 (L=345 [R=437 UT=15] Andronalization (A)

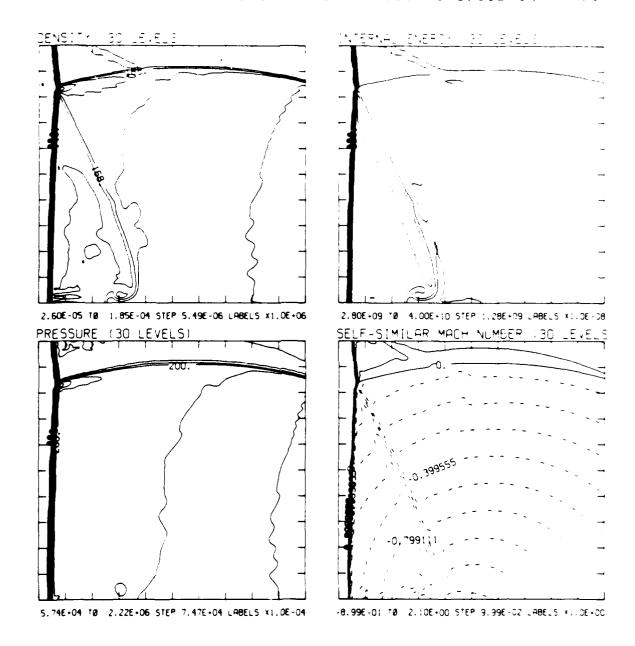


Figure 24.1b. $\theta_{W} = 5^{\circ}$, blowup-frame plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

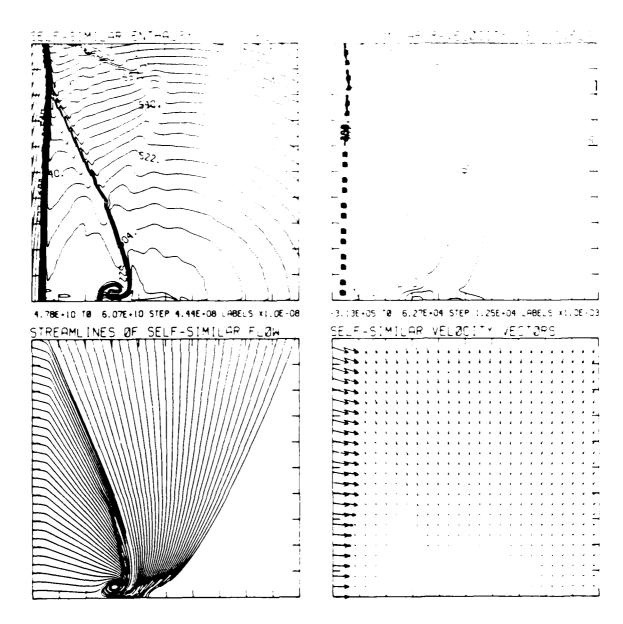
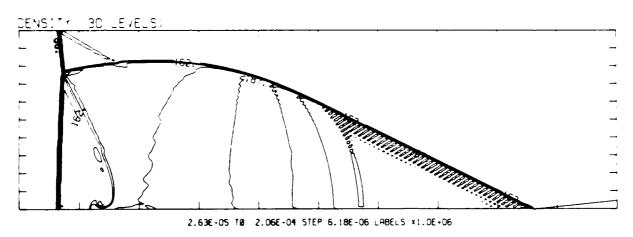


Figure 24.1b. $\theta_{\rm W}$ = 5°, blowup-frame plots - continued. Figure 24. Transition set 3, $\rm M_S$ = 8.75, Hansen - continued.

MS# 8.75 ALP= 8.00 12#=525 NZ#160 KBEG# 75 FG#2.00E+04 HANSEN



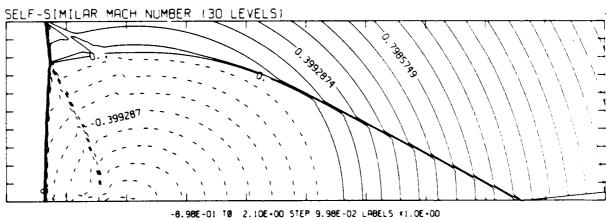


Figure 24.2a. $\theta_w = 6^{\circ}$, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP= 6.00 (1=345 (R=497 UT=150 PO=2.00E+04 HANSEN

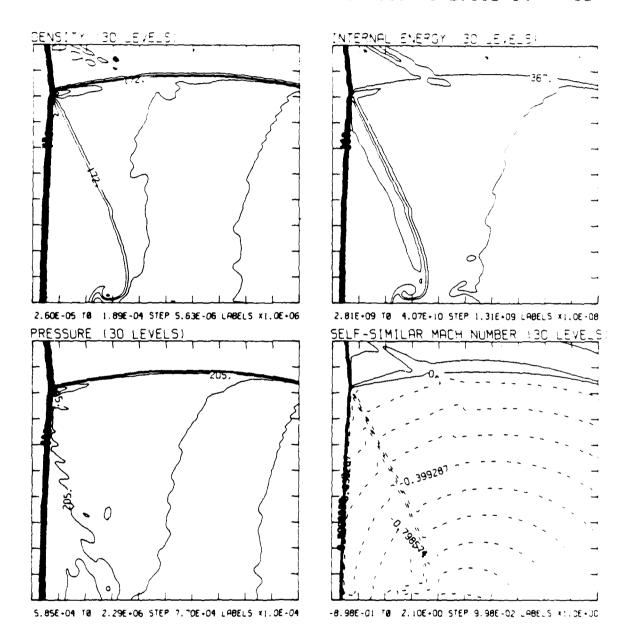


Figure 24.2b. $\theta_{w} = 6^{\circ}$, blowup-frame plots.

Figure 24. Transition set 3, $M_s \approx 8.75$, Hansen - continued.

MS= 8.75 ALP= 6.00 [L=345 [R=497 UT=150 PC=2.00E+04 HANSEN

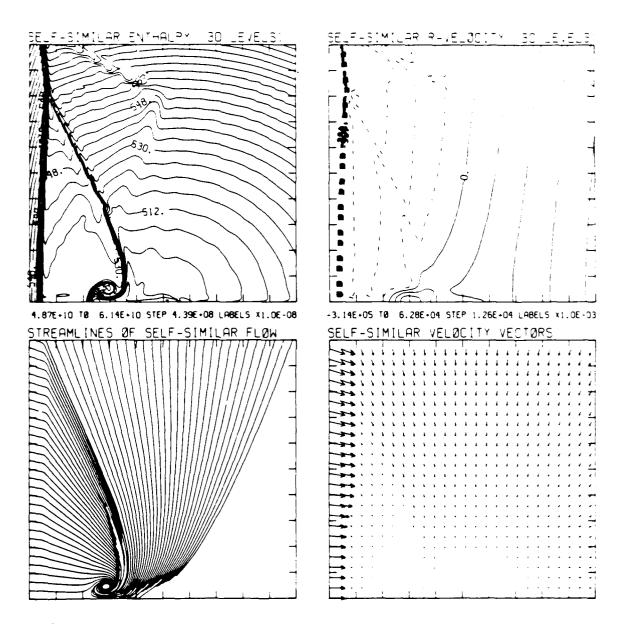
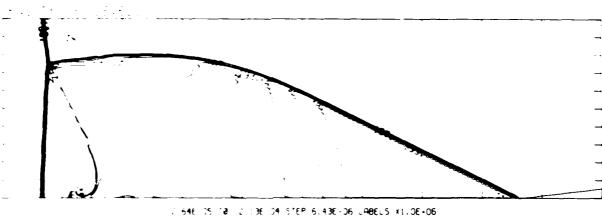
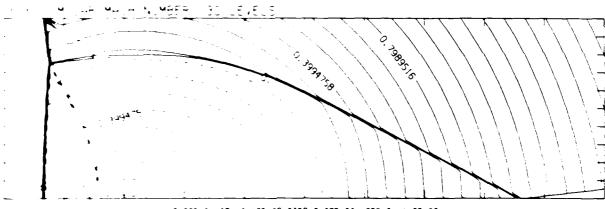


Figure 24.2b. $\theta_{\rm w}$ = 6°, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.





8 996 0. 18 2:106+00 STEP 9:396-02 LABELS X1.06+00

Figure 24.3a. $A_{\rm w}=7^{\circ}$, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_5 = 8.75$, Hansen - continued.

MS= 8.75 ALP= 7.00 [L=345 [R=497 UT=150 PO=2.00E+04 HANSEN

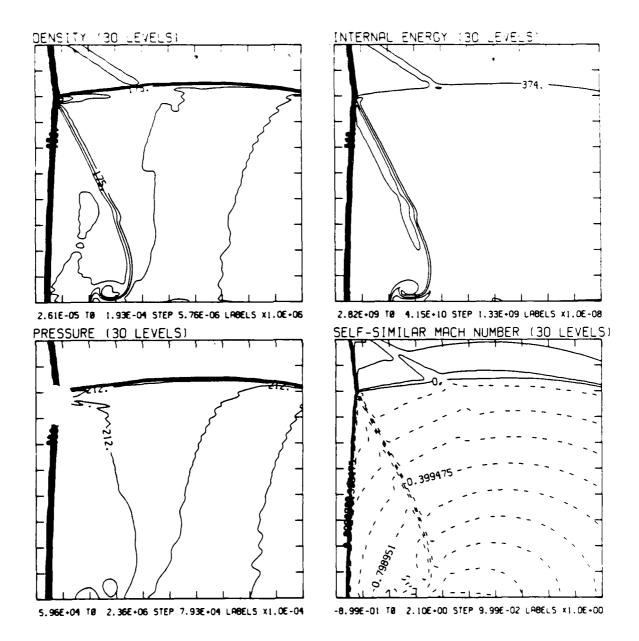


Figure 24.3b. $\theta_w = 7^\circ$, blowup-frame plots.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP= 7.00 [L=345 [R=497 LT=150 PC=2.00E+04 HANSEN

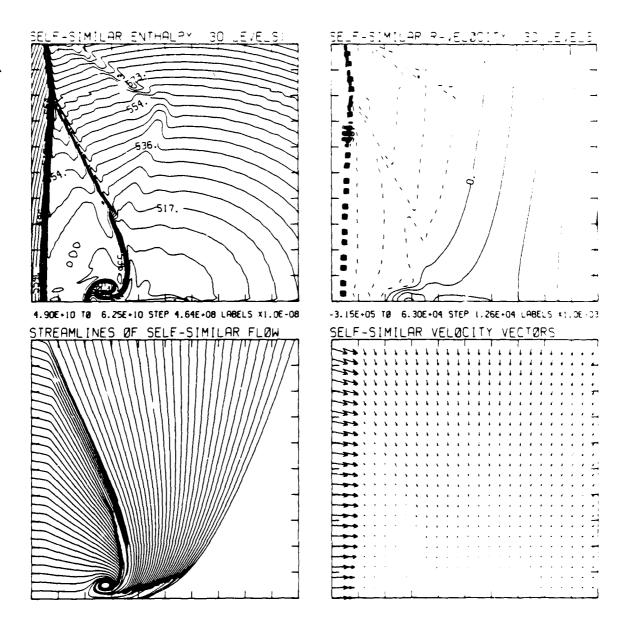
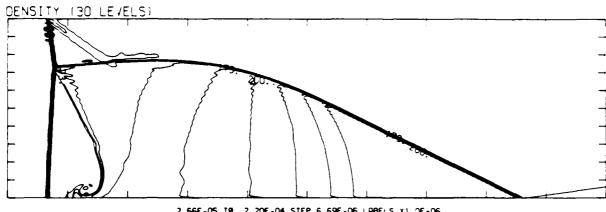


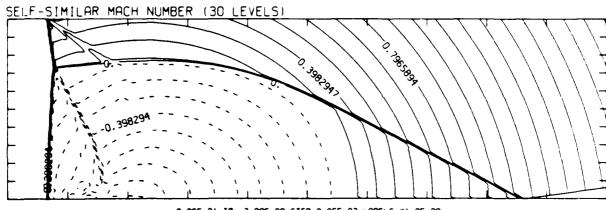
Figure 24.3b. $\theta_w = 7^\circ$, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP= 8.00 NR=525 NZ=160 KBEG= 75 PC=2.005+04 HANSE'.



2.66E-05 T8 2.20E-04 STEP 6.69E-06 LABELS X1.0E+06



-8.96E-01 TØ 2.09E+00 STEP 9.96E-02 LABELS X1.0E+00

Figure 24.4a. $\theta_w = 8^{\circ}$, whole-flowfield contour-plots.

Figure 24. Transition set 3, M_S = 8.75, Hansen - continued.

MS= 8.75 ALP= 8.00 IL=345 [R=497 UT=150 PC=2.00E+04 HANSEN

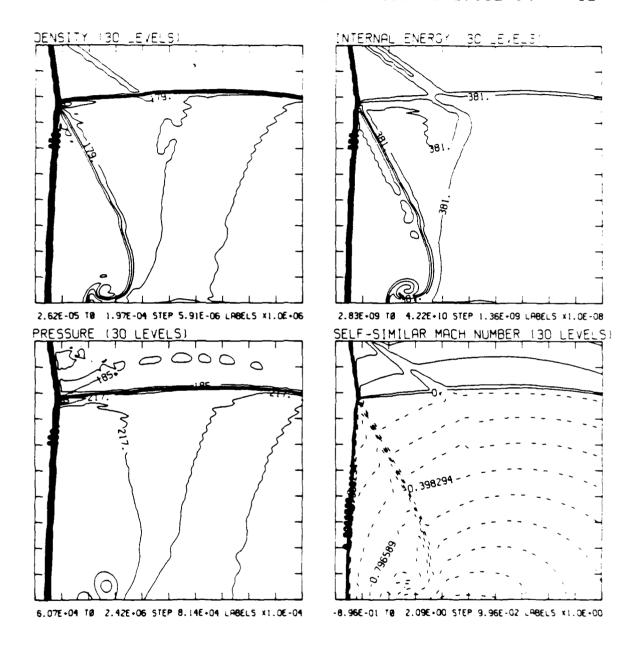


Figure 24.4b. $\theta_{\rm W}$ = 8°, blowup-frame plots.

Figure 24. Transition set 3, M_S = 8.75, Hansen - continued.

MS= 8.75 ALP= 8.00 (L=345 (R=497 UT=150 P0=2.00E+04 HANSE).

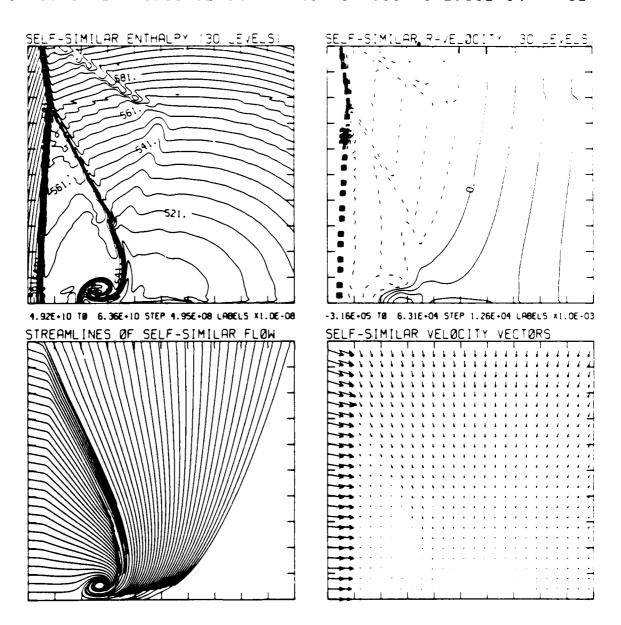
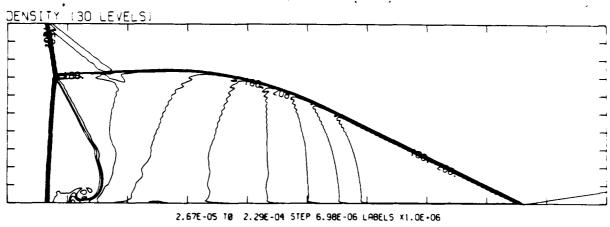
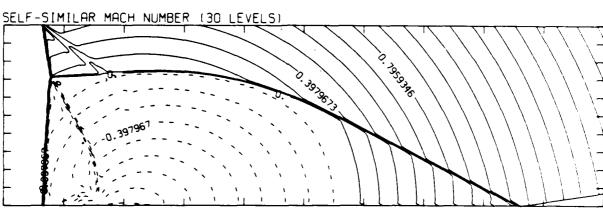


Figure 24.4b. $\theta_{\rm W}$ = 8°, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP= 9.00 NR=525 NZ=160 KBEQ= 75 PO=2.00E+64 HANSEN





-8.95E-01 TØ 2.09E+00 STEP 9.9SE-02 LABELS X1.0E+00

Figure 24.5a. $\theta_{\rm W}$ = 9°, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP= 9.00 [L=345 [R=497 UT=150 PC=2.00E+04 HANSEN

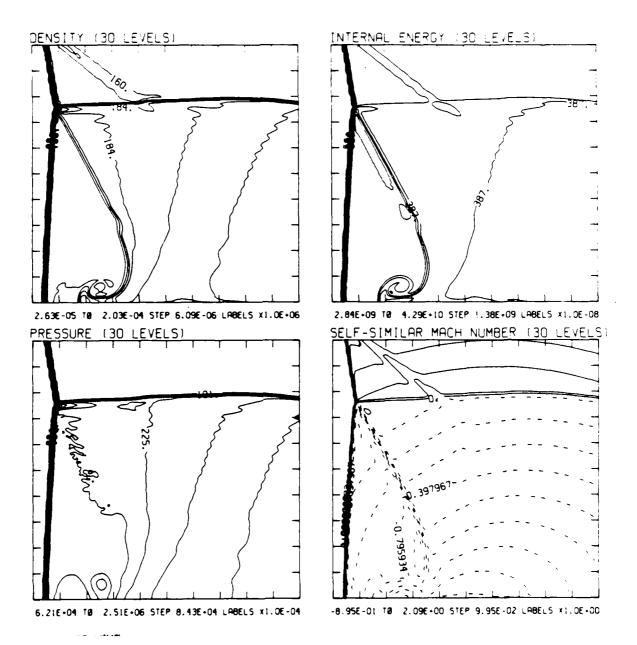


Figure 24.5b. $\theta_{\rm w}$ = 9°, blowup-frame plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP= 9.30 [L=345 [R=497 LT=150 PD=2.00E+04 HANSEN

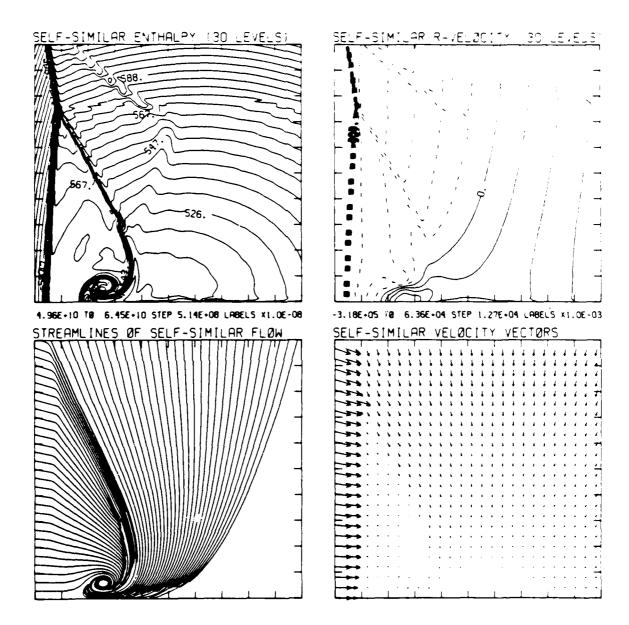
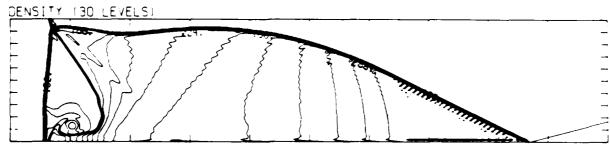


Figure 24.5b. $\theta_w = 9^{\circ}$, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP=15.00 NR=550 NZ=115 KBEG= T5 P0=2.00E+04 Hansen



2.77E-05 TO 2.87E-04 STEP 8.93E-06 LABELS X1.0E+06

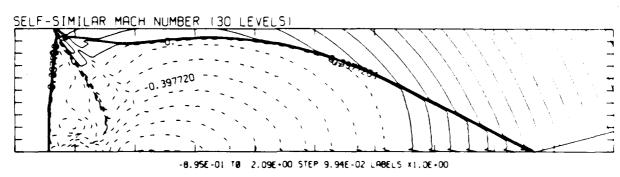


Figure 24.6a. $\theta_{\rm w}$ = 15°, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALF=15.00 (L=406 (P=523 UT=115 P0=2.00E+04 HANSEN

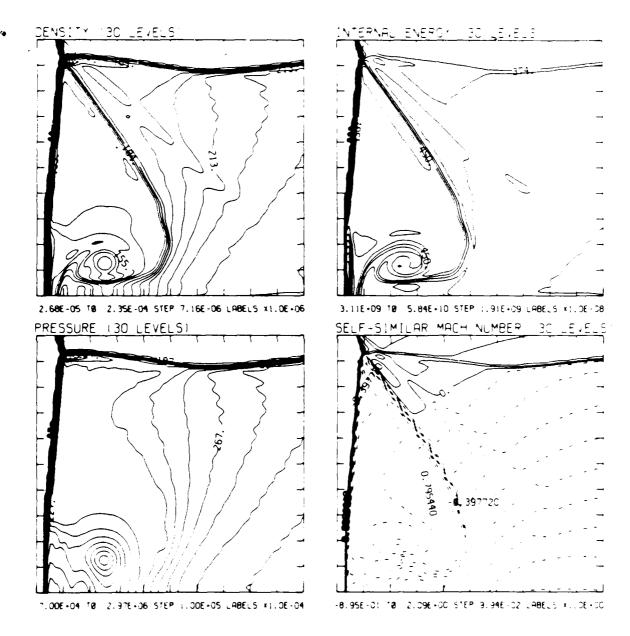


Figure 24.6b. $\theta_{w} = 15^{\circ}$, blowup-frame plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 3.75 ALP=15.00 [L=406 [P=503 LT=115 PD=0.00E+04 H4N.55%

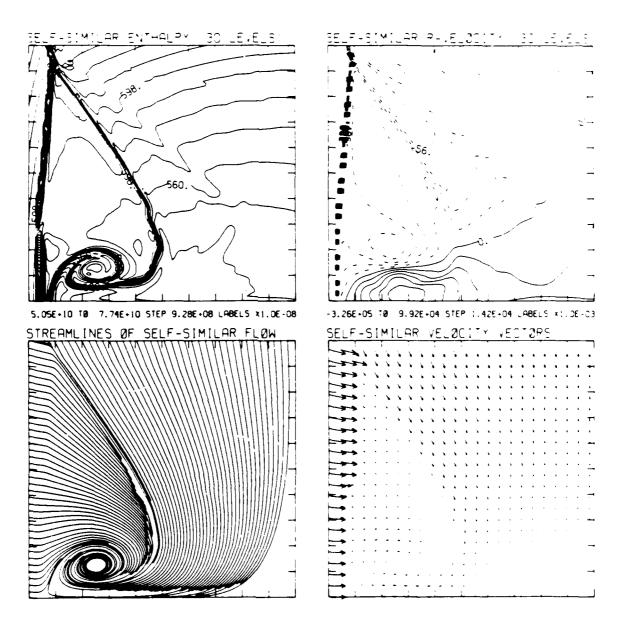
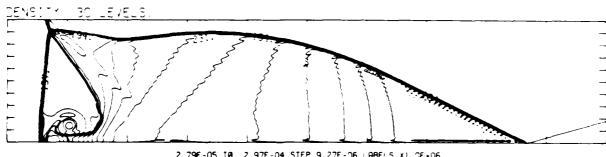


Figure 24.6b. $\theta_{\rm w}$ = 15°, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=16.00 NP=550 NZ=115 ABEG= 75 PG=2.00E+04 HANSEN



2.79E-05 TO 2.97E-04 STEP 9.27E-06 LABELS X1.0E+06

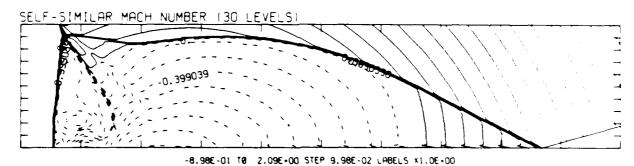


Figure 24.7a. $\theta_{W} = 16^{\circ}$, whole-flowfield contour-plots.

Figure 24. Transition set 3, M_S = 8.75, Hansen - continued.

MS= 8.75 ALP=16.00 IL=406 IR=523 UT=115 PO=2.00E+04 HANSEN

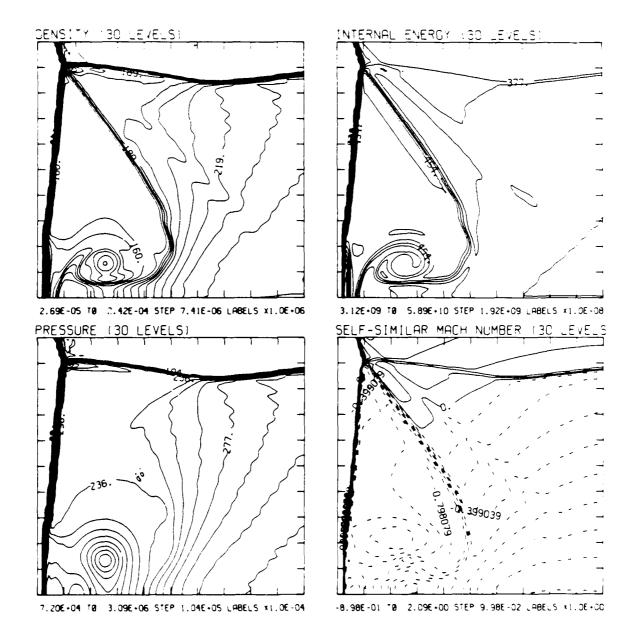


Figure 24.7b. $\theta_{\rm w}$ = 16°, blowup-frame plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

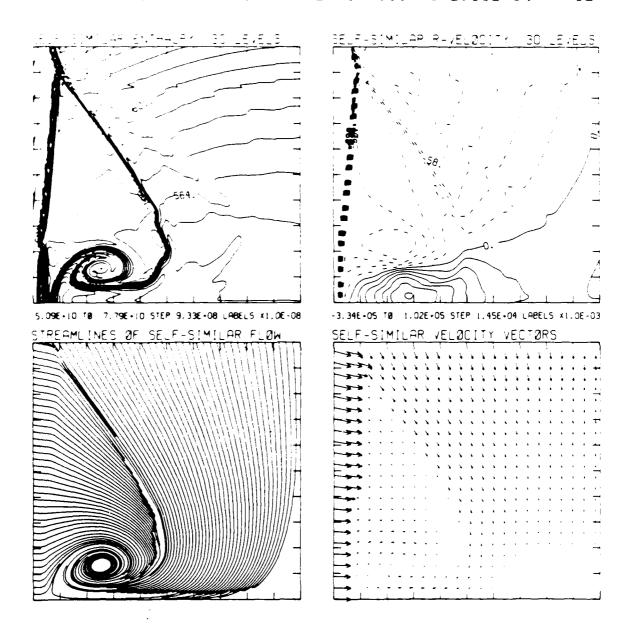
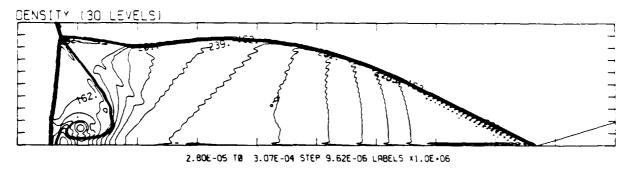


Figure 24.7b. $\theta_{\rm W}$ = 16°, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=17.00 NR=550 NZ=115 KBEG= 75 PC=2.00E+64 HANSEN



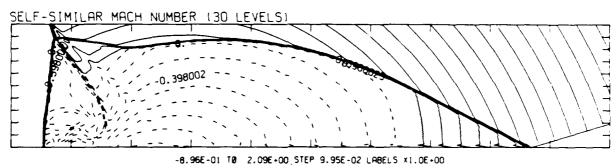


Figure 24.8a. $\theta_{\rm W}$ = 17°, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP=17.00 [L=406 [F=528 UT=115 PD=2.00E+04 HANSEN

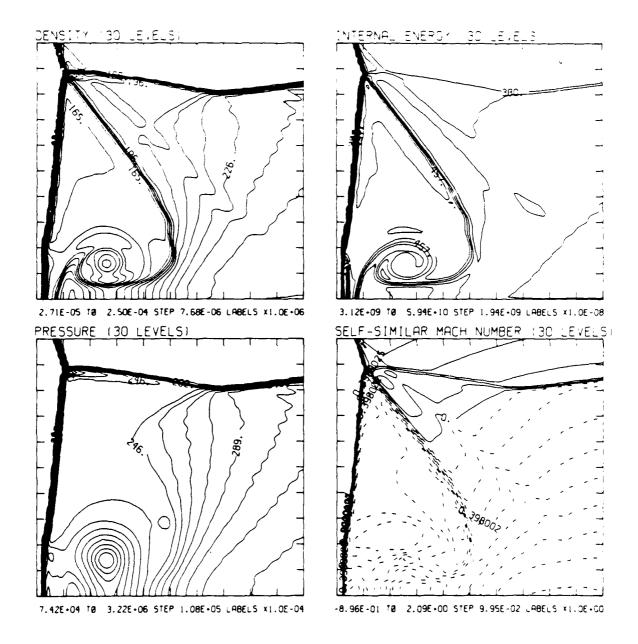


Figure 24.8b. $\theta_{\rm W}$ = 17°, blowup-frame plots.

Figure 24. Transition set 3, $M_s = 8.75$, Hansen - continued.

MS= 8.75 ALP=17.00 IL=406 [R=523 UT=115 PC=2.00E+04 HANSEN

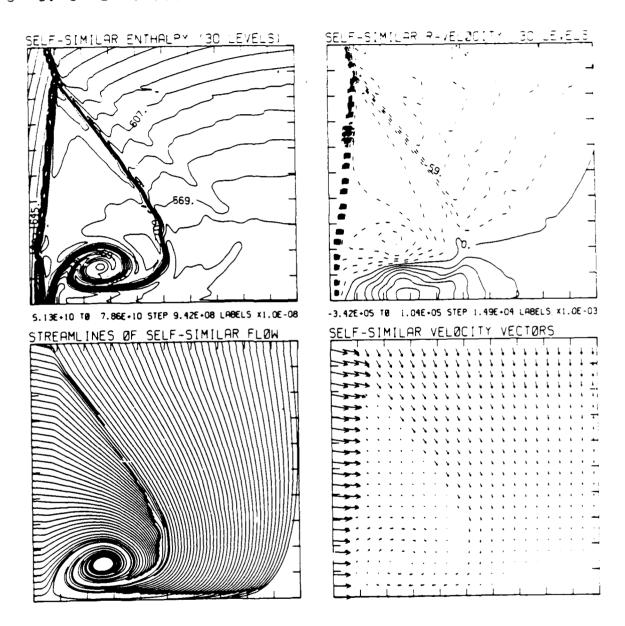
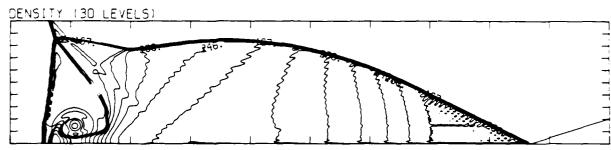


Figure 25.8b. $\theta_{\rm W}$ = 17°, blowup-frame plots - continued.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=18.00 NR=550 NZ=115 KBEG= 75 P0=2.00E+04 HANSEN



2.82E-05 T0 3.16E-04 STEP 9.92E-06 LABELS X1.0E+06

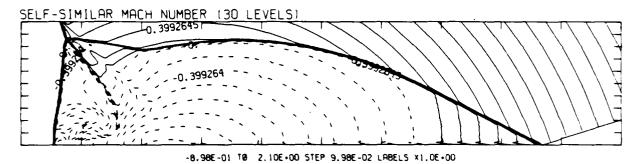


Figure 24.9a. $\theta_{\rm W}$ = 18°, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=18.00 (L=406 (R=523 UT=115 PD=2.00E+04 HANSEY

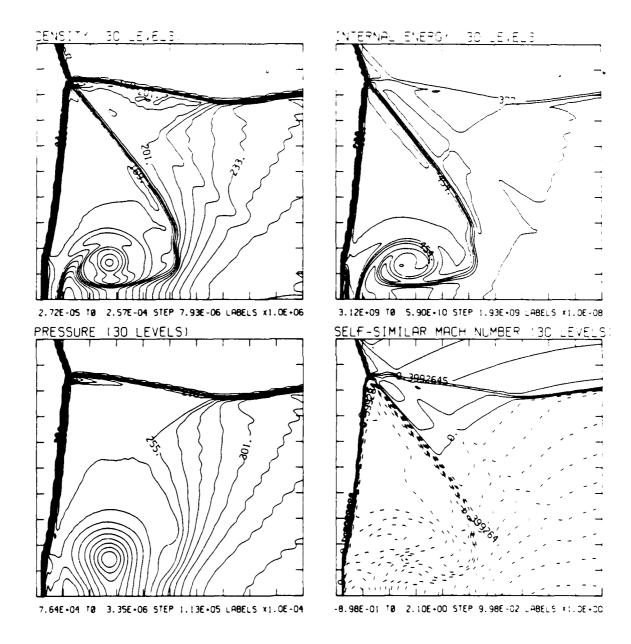


Figure 24.9b. $\theta_{\rm W}$ = 18°, blowup-frame plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=18.00 [L=406 [R=523 UT=115 PO=2.00E+04 H44.5E4.

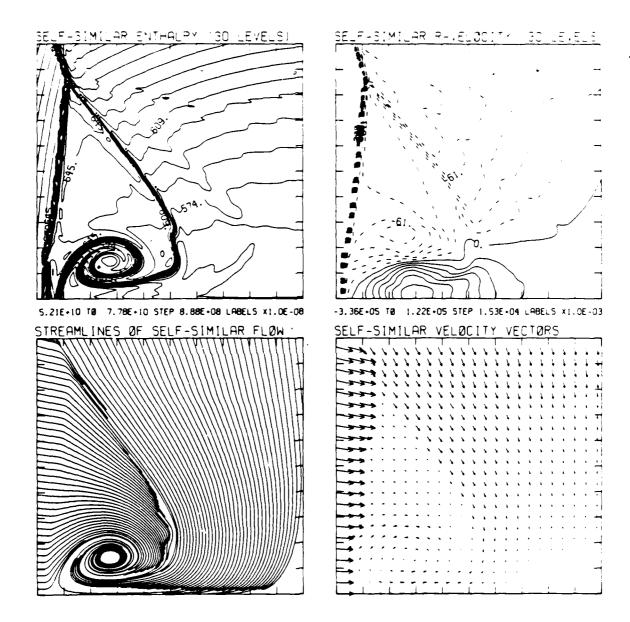
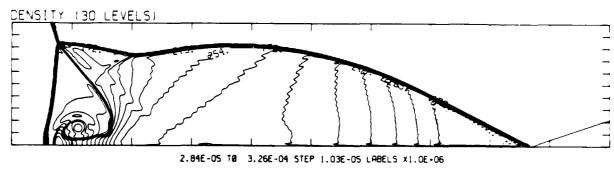


Figure 24.9b. $\theta_{\rm W}$ = 18°, blowup-frame plots - continued.

Figure 24. Transition set 3, M_s = 8.75, Hansen - continued.

MS= 8.75 ALP=19.00 NR=550 NZ=115 KBEG= 75 PC=2.00E+04 HANSEN



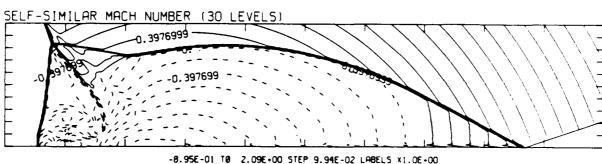


Figure 24.10a. $\theta_{\rm W}$ = 19°, whole-flowfield contour-plots.

Figure 24. Transition set 3, $M_S = 8.75$, Hansen - continued.

MS= 8.75 ALP=19.00 [L=406 [R=523 LT=115 PC=2.00E+04 HaNSEN

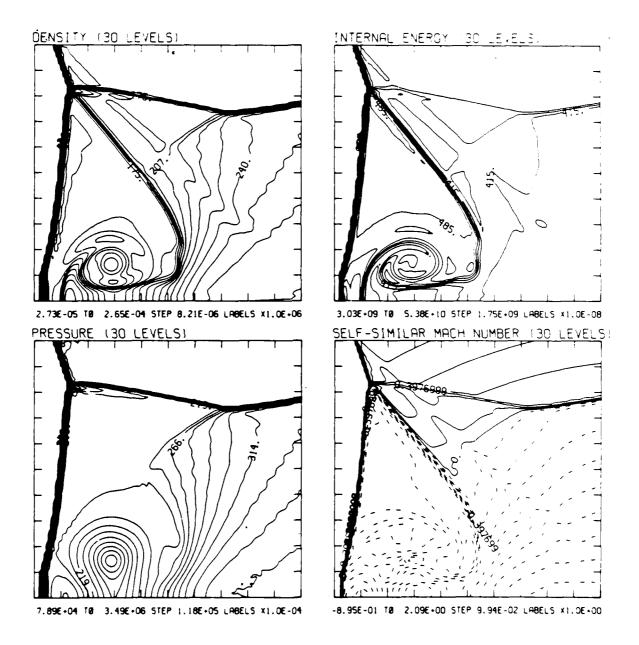


Figure 24.10b. $\theta_{\rm W}$ = 19°, blowup-frame plots.

Figure 24. Transition set 3, M_S = 8.75, Hansen - continued.

MS= 8.75 ALP=19.00 IL=406 IR=523 JT=115 P0=2.00E+04 HANSEN

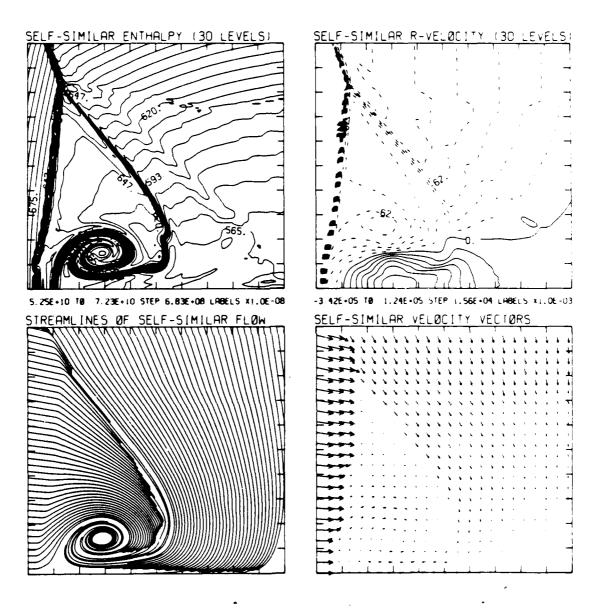


Figure 24.10b. $\theta_{\rm W}$ = 19°, blowup-frame plots - continued. Figure 24. Transition set 3, M_S = 8.75, Hansen - continued.



Region	:1.
()	1.00
1	3.78
2	6.69
3	3.91
а	7.94
1	9.19
L	6.69
d	5.44

Figure 25a. Interferogram, $\theta_{\mathbf{w}} = 49^{\circ}$

18호 지, 18 유 의원교육, 38 NA ERMA 시간 사이 교육의



Figure 25b. $\theta_{\rm w} = 490$

XBB 859-7209

Figure 25. Transition set 4, $M_S = 7.10$, $\gamma = 5/3$, density contour-plots

MS= 7.10 ALP=50.00 NR=575 NZ=115 KBEG= 80 PG=2.005+04 ARODY.

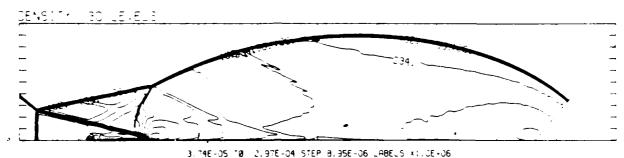


Figure 25c. $\theta_{\rm w} = 500$

MS= 7.10 ALP=51.00 NR=575 NZ=115 KBEG= 80 PC=2.005+04 AF021

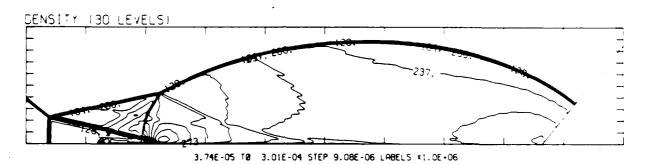


Figure 25d. $\theta_{\omega} = 510$

MS= 7.10 ALP=52.00 NR=575 NZ=115 KBEG= 80 PC=2.00E+34 ARGS.

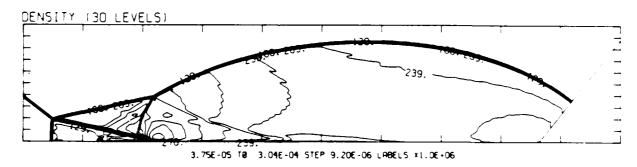


Figure 25e. $\theta_{w} = 52^{\circ}$

Figure 25. Transition set 4, M_S = 7.10, γ = 5/3, density contour plots - continued.

MS= 7.10 ALP=52.75 NR=575 NZ=115 KBEG= 80 PC=2.00E+04 ARDZN

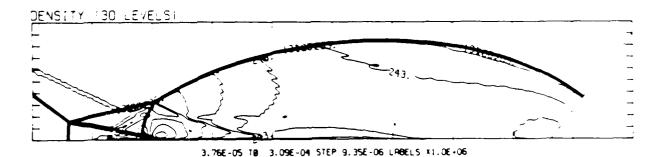


Figure 25f. $\theta_{w} = 52.750$

MS= 7.10 ALP=53.00 NR=575 NZ=115 KBEG= 80 PC=2.00E+14 AF00.

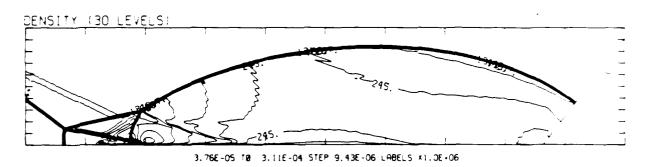


Figure 25g. $\theta_{w} = 53.00$

MS= 7.10 ALP=53.10 NR=575 NZ=115 KBEG= 80 PC=2.00E+34 AFIZ".

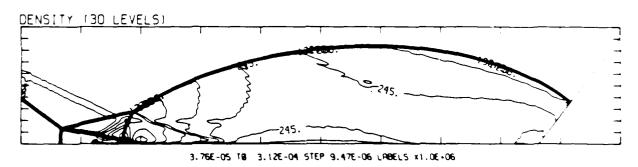
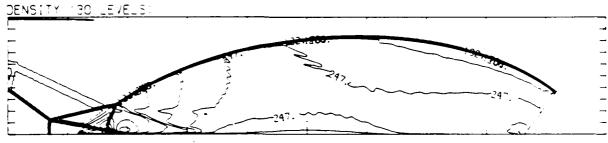


Figure 25h. $\theta_{w} = 53.10^{\circ}$

Figure 25. Transition set 4, M_S = 7.10, γ = 5/3, density contour plots - continued.

MS= 7.10 ALP=53.20 NR=575 NZ=115 KBEG= 80 PG=2.00E+04 APGC1.



3.77E-05 TØ 3.14E-04 STEP 9.52E-06 LABELS X1.0E+06

Figure 25i.
$$\theta_{w} = 53.20^{\circ}$$

MS= 7.10 ALP=53.30 NR=575 NZ=115 KBEG= 80 PO=2.00E+04 APDZ*,

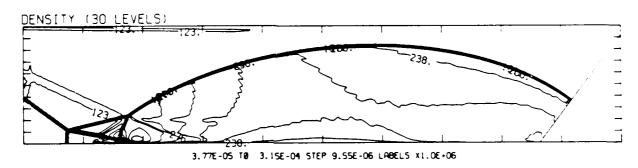


Figure 25j. $\theta_{\rm w} = 53.30^{\circ}$

MS= 7.10 ALP=53.40 NR=575 NZ=115 KBEG= 80 PO=2.00E+34 AF32".

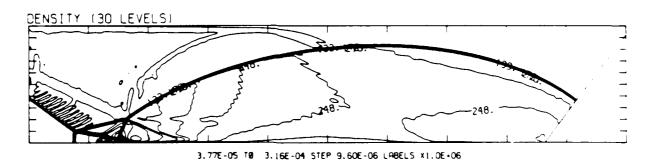
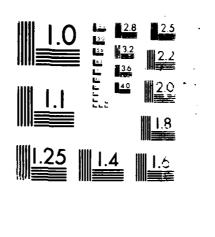


Figure 25k. $\theta_{\rm w} = 53.40^{\circ}$

Figure 25. Transition set 4, M_S = 7.10, γ = 5/3, density contour plots - continued.

A DETAILED NUMERICAL GRAPHICAL AND EXPERIMENTAL STUDY OF DBLIQUE SHOCK MA (U) TORONTO UNITY DOMNSYIEM (ONTRRIO) INST FOR AEROSPACE STUDIES. H M GLAZ ET AL 81 AUG 86 UTIRS-285 DNA-TR-86-365 F/G 20/4 AD-8186 440 5/5 H H GLAZ ET AL F/G 20/4 UNCLASSIFIED NL.



MICROCOPY RESOLUTION TEST CHAINATIONS HUBES, IN STANCARDS (186-

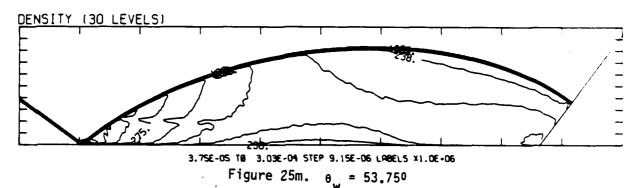
MS= 7.10 ALP=53.50 NR=575 NZ=115 KBEG= 80 PO=2.00E+04 ARGCV



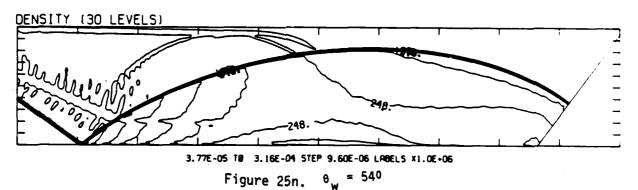
3.77E-05 TO 3.18E-04 STEP 9.65E-06 LABELS X1.0E+06

Figure 251. $\theta_{...} = 53.50^{\circ}$

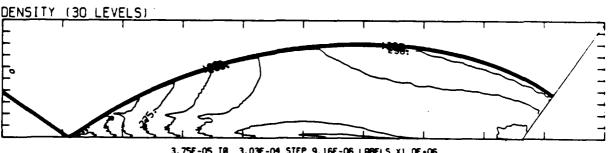
MS= 7.10 ALP=53.75 NR=575 NZ=115 KBEG= 80 PO=2.00E+04 APG21.



MS= 7.10 ALP=54.00 NR=575 NZ=115 KBEG= 80 PO=2.00E+04 APGC.



MS= 7.10 ALP=55.00 NR=575 NZ=115 KBEG= 80 PO=2.00E+04 AFG2*.



3.75E-05 TØ 3.03E-04 STEP 9.16E-06 LRBELS X1.0E+06

Figure 250. $\theta_{\rm W}$ = 550

Figure 25. Transition set 4, M_S = 7.10, γ = 5/3, density contour plots continued.

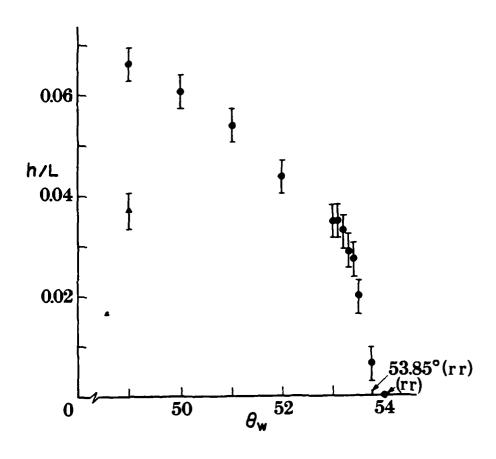


Figure 26. Plot of DMR Mach stem height versus θ_w , extrapolated to zero height for RR(h/L = 0 for θ_w = 53.85°), h/L = 0 for θ_w = 54° is a numerical result

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